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The Algebra of Mirković-Vilonen Cycles in Type A

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This paper is dedicated to Robert MacPherson on the occasion of his 60th birthday.

Abstract: Let Gr be the affine Grassmannian for a connected complex reductive group G. Let \mathcal{C}_G be the complex vector space spanned by (equivalence classes of) Mirković-Vilonen cycles in Gr. The Beilinson-Drinfeld Grassmannian can be used to define a convolution product on MV-cycles, making \mathcal{C}_G into a commutative algebra. We show, in type A, that \mathcal{C}_G is isomorphic to $\mathbb{C}[N]$, the algebra of functions on the unipotent radical N of a Borel subgroup of G; then each MV-cycle defines a polynomial in $\mathbb{C}[N]$, which we call an MV-polynomial. We conjecture that those MV-polynomials which are cluster monomials for a Fomin-Zelevinsky cluster algebra structure on $\mathbb{C}[N]$ are naturally expressible as determinants, and we conjecture a formula for many of them.

1. Introduction

Since Mirković and Vilonen first announced their discovery of a new geometric canonical basis in representation theory [MV1], the study of their algebraic varieties—here called MV-cycles—has grown quickly, with applications to geometry, representation theory, and combinatorics. Let us begin by mentioning a few of these.

In the extended version of their original paper, Mirković and Vilonen [MV2] used MV-cycles to study the relation between the representation theory of a complex reductive group and the equivariant perverse sheaves on the affine Grassmannian **Gr** for the Langlands dual group. The MV-cycles are subvarieties of

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Gr giving a canonical basis for every irreducible representation. Braverman and Gaitsgory [BG] used MV-cycles to give a geometric definition of crystal graphs. Gaussent and Littelmann [GL] gave an interpretation of Littelmann's path model using MV-cycles. Kamnitzer [K] discovered the relation between the combinatorics of the MV-polytopes (moment map images of MV-cycles) and the Berenstein-Zelevinsky combinatorics used to compute Littlewood-Richardson coefficients [BZ].

Essential to the Mirković-Vilonen results was the Beilinson-Drinfeld Grassmannian. This is a space which allows one to define a convolution product of MV-cycles (sometimes called a fusion product), making the vector space spanned by all MV-cycles into a commutative algebra. The purpose of this paper is to discuss this algebra and its combinatorial properties in type A. In particular, we identify it with a polynomial algebra, and conjecture a combinatorial formula for some of the polynomials.

We will continue in the spirit of [AK], in which the authors computed all MV-polytopes in type A; we use mostly elementary tools and combinatorial models, and we avoid abstraction, always preferring to make concrete choices that make it easy to work with examples. There are two reasons for working in this way: (1) Our results and conjectures can be stated in this elementary language. (2) The literature about MV-cycles often gives the inaccurate impression that considerable heavy machinery and abstruse mathematics is required to understand them.

Let us now discuss the main results and conjectures of the paper. As mentioned above, the MV-cycles are certain algebraic varieties lying inside the affine Grassmannian \mathbf{Gr} for a connected complex reductive group \mathbf{G} . The standard short definition of \mathbf{Gr} is that $\mathbf{Gr} = \mathbf{G}(\mathcal{K})/\mathbf{G}(\mathcal{O})$, where $\mathcal{K} = \mathbb{C}((t))$ and $\mathcal{O} = \mathbb{C}[[t]]$. But, for our purposes, this definition is difficult to work with and provides little geometric intuition; we prefer to use a version of Lusztig's lattice model [Lu] for \mathbf{Gr} in type A, which we will discuss in Section 2.1. Although the affine Grassmannian is an infinite-dimensional space, it is an increasing union of finite-dimensional pieces (formally an ind-scheme). In any example we will be interested in, we may work completely within one of these pieces.

One can also give a concise definition of MV-cycles: A maximal torus T of G acts on \mathbf{Gr} and has fixed point set canonically isomorphic to the coweight lattice $\Lambda = \operatorname{Hom}(\mathbb{C}^*, T) \colon \lambda \in \Lambda \leftrightarrow \underline{\lambda} \in \mathbf{Gr}$. Moreover, this action is Hamiltonian and has a naturally defined moment map $\Phi : \mathbf{Gr} \to \mathfrak{t}_{\mathbb{R}} = \operatorname{Lie}(\mathsf{T})$. Consider the Morse flow of a generic component $\xi : \mathbf{Gr} \to \mathbb{R}$ of the moment map. Let S_{λ} be the space of points of \mathbf{Gr} which converge to $\underline{\lambda}$. Similarly, define T_{λ} for the inverse flow. (Actually, later we will pick a concrete component ξ .) Then MV-cycles are defined to be the irreducible components of the closures of $S_{\lambda} \cap T_{\mu}$. Their moment map images are called MV-polytopes. In Section 2.3 we will repeat this definition in the language of the lattice model, and also recall from [AK] how to

describe and parametrize the irreducible components (Theorem 1), a description we will use heavily throughout.

Our primary concern is with the convolution product of MV-cycles: given two MV-cycles M and N, this expresses their product as a linear combination of MV-cycles

$$(1.1) M*N = \sum c_{MN}^K K.$$

This convolution is defined using the Beilinson-Drinfeld Grassmannian. This space is a fibration over \mathbb{C} whose fiber over zero is \mathbf{Gr} and whose fibers over \mathbb{C}^* are isomorphic to $\mathbf{Gr} \times \mathbf{Gr}$. Within this family, it is possible to construct a family $M *_{\mathbb{C}} N$ which degenerates $M \times N$ to the union of K's with multiplicities c_{MN}^K . We will present in Section 3 a lattice model for this family, which will allow us to identify, in type A, the leading term on the right side of (1.1). (Propositions 3.2 and 3.3).

This will result in an algebra \mathcal{C}_G for which the MV-cycles provide a basis, as a complex vector space. (Actually, we will have some simple equivalence relation on MV-cycles, and \mathcal{C}_G will be generated by the equivalence classes.) It was conjectured by the first author [A1] that \mathcal{C}_G is isomorphic as a Hopf algebra to the ring of functions $\mathbb{C}[N]$ on the unipotent radical N of a Borel subgroup of G. In Section 4 we will prove part of this conjecture in type A, where N is the group of upper triangular unipotent matrices. Using the properties we have proved about convolution, we will construct a map $\mathcal{I}:\mathbb{C}[N] \to \mathcal{C}_G$ and prove that it is an algebra isomorphism (Theorem 2). So corresponding to each MV-cycle $M \in \mathcal{C}_G$ we have a polynomial $\mathcal{P} \in \mathbb{C}[N]$; we call these MV-polynomials.

Section 5 addresses the computation of MV-polynomials and their relation to cluster algebras. The simplest MV-cycles are Richardson varieties (intersections of opposite Schubert varieties) and the corresponding MV-polynomials are exactly the minors in $\mathbb{C}[N]$. But understanding MV-polynomials in general seems to be a difficult problem and we will only conjecture how to compute some of them.

Of crucial theoretical and practical importance to our conjectures is the Fomin-Zelevinsky theory of cluster algebras [FZ1, FZ2, BFZ]. Part of this structure is a canonical set of polynomials in $\mathbb{C}[N]$ —the *cluster monomials*. We conjecture (Conjecture 5.1) that these are themselves MV-polynomials, and, moreover, are all naturally expressible as certain determinants.

Although the inductive computation of MV-polynomials from the definition of convolution of MV-cycles is difficult, the computation of the cluster monomials is straightforward and has been programmed onto a computer. (The algorithm is specified in Appendix A and our data for type A_5 may be found at $[AK^*]$.) This provided many examples to work with, based upon which we sharpened Conjecture 5.1 to Conjecture 5.2, which gives an explicit determinantal formula for a large class of MV-polynomials.

Finally, Conjecture 5.3 relates the MV-polynomials that are cluster monomials to the geometry of MV-cycles. The most intriguing thing about all our conjectures is the connection between cluster algebras and MV-cycles, which we hope will eventually be understood on a deeper level.

To summarize, Section 2 recalls what we need from [AK]: the definition of the lattice model for \mathbf{Gr} and the explicit parametrization and description of MV-cycles in type A. Section 3 discusses the convolution product of MV-cycles using a lattice model for the Beilinson-Drinfeld Grassmannian. In Section 4 we prove that the convolution algebra of MV-cycles is isomorphic to the polynomial algebra $\mathbb{C}[\mathsf{N}]$, and use this isomorphism to define MV-polynomials. Section 5 contains conjectural determinantal formulas for some MV-polynomials. Appendix A is a terse summary of what we use from cluster algebras.

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2. MV-Cycles

In this section, we review the type A description of MV-cycles and their parametrization by the Kostant parameter set, as set forth in [AK]. We use throughout the lattice model of the affine Grassmannian, the space in which the MV-cycles live.

- 2.1. Lattice model. Following Lusztig [Lu], we will introduce the lattice model of the affine Grassmannian \mathbf{Gr} for $\mathsf{G} = \mathsf{GL_n}$. Points in \mathbf{Gr} are subspaces of a certain infinite-dimensional complex vector space X, satisfying some extra conditions. The vector space X is defined by specifying a basis: if e_1, e_2, \ldots, e_n denotes the standard basis for \mathbb{C}^n , then X is the span of $t^j e_i$ for $1 \leq i \leq n, -\infty < j < \infty$, where we regard this as a symbol with two indices. We usually picture these basis vectors in an array consisting of n columns, infinite in both directions. Let t be the invertible linear operator on X that sends each $t^j e_i$ to $t^{j+1} e_i$. Then \mathbf{Gr} consists of those subspaces Y of X such that
 - (1) Y is closed under the action of t.
 - (2) $t^N X_0 \subseteq Y \subseteq t^{-N} X_0$ for some N, where X_0 is the span of those $t^j e_i$ with $j \geq 0$.

We call such Y lattices.

Three examples of lattices are illustrated by the pictures in Figure 1. Each dot, or set of dots connected by line segments, represents a vector; these, together with all the dots t^je_i , $j \geq 3$ below the pictures are a \mathbb{C} -basis for the lattice. The first picture represents the span of the basis vectors drawn, namely t^je_1 with $j \geq -2$, t^je_2 , t^je_3 with $j \geq 1$, and t^je_4 , t^je_5 , t^je_6 with $j \geq 2$. The second picture represents the lattice spanned by the vectors $t^{-2}e_1 + 3e_2 + 4e_3$, $3te_2 + 4te_3$,

 $-te_3 + 6te_5$, $t^2e_4 + 2te_5 - te_6$, and all the vectors of the form t^je_i corresponding to the unconnected dots. Similarly for the third picture.

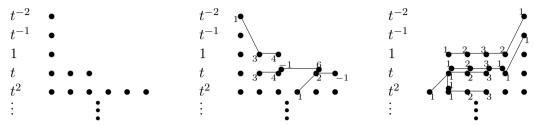


Figure 1. Examples of lattices.

Lattices spanned by vectors $t^j e_i$ are in one-to-one correspondence with the coweights of $\mathsf{GL_n}$ (which we will regard as n-tuples of integers). The first picture in Figure 1 corresponds to the coweight (2,-1,-1,-2,-2,-2) according to our convention. A coweight with nonincreasing entries, like this one, is said to be dominant; one with nondecreasing entries is said to be antidominant. In general, if λ is a coweight then we denote the corresponding lattice by $\underline{\lambda}$. Also, any coweight $\lambda = (\lambda_1, \lambda_2, \dots, \lambda_n)$ acts on X by shifting the i^{th} column by λ_i : $\lambda \cdot t^j e_i = t^{j-\lambda_i} e_i$. This allows us to form the lattice $\lambda \cdot Y$, the shift of Y by λ , by setting $\lambda \cdot Y = \{\lambda \cdot v \mid v \in Y\}$. It is then clear that $\lambda \cdot \mu = \lambda + \mu$.

Note that if N is fixed in condition (2), then the set of all lattices is a finite-dimensional projective variety \mathbf{Gr}^N . So \mathbf{Gr} is an increasing union of complex projective varieties. Moreover, although the definition of \mathbf{Gr} at first sounds very infinite-dimensional, in practice everything is finite-dimensional: all but a finite-dimensional subspace of each lattice is spanned by basis vectors, and the varieties we will study will each be contained in one of the \mathbf{Gr}^N .

The affine Grassmannian Gr for GL_n is not connected. Each connected component contains all lattices Y with fixed dimension $\dim_0 Y = \dim(Y/Y \cap X_0) - \dim(X_0/Y \cap X_0)$ relative to X_0 . All the connected components are isomorphic to each other by shifts. Also, the connected component where $\dim_0 Y = 0$ is the affine Grassmannian for SL_n .

2.2. Torus action and moment map. There is an algebraic action of an (n-1)-dimensional torus T on \mathbf{Gr} : $(\mathbb{C}^*)^n$ acts on X by multiplication on the n columns, and we divide by the trivial action of the diagonal one-dimensional subtorus. (More precisely, $(z_1, z_2, \ldots, z_n) \in (\mathbb{C}^*)^n$ acts on a vector $t^j e_i$ by multiplying it by z_i . This action on X induces an action on \mathbf{Gr} .) Note that the lattices corresponding to coweights are exactly the torus-fixed points.

This torus action gives rise to a moment map Φ from \mathbf{Gr} to a real (n-1)-dimensional Euclidean space $\mathfrak{t}_{\mathbb{R}}$. It maps compact irreducible torus-invariant subvarieties of \mathbf{Gr} onto convex polytopes. It is easy to describe this map explicitly in terms of lattices: We first define a Hermitian inner product on our ambient vector space X by declaring $\{t^j e_i\}$ to be an orthonormal basis. We also take $\mathfrak{t}_{\mathbb{R}}$

to be the hyperspace perpendicular to the vector $(1,1,\ldots,1)$ in \mathbb{R}^n (with respect to the standard inner product). If $Y=\underline{\lambda}$ is a torus-fixed point, we may define $\Phi(Y)$ to be the orthogonal projection of $\lambda\in\mathbb{R}^n$ onto $\mathfrak{t}_{\mathbb{R}}$. On an arbitrary lattice Y, first write it as an orthogonal sum $Y=\underline{\lambda}\oplus V$ where V is a finite-dimensional vector space. (For instance, take $\underline{\lambda}=\operatorname{span}\{t^je_i\mid t^je_i\in Y\}$, and take V to be the orthogonal complement in Y.) Choose an orthonormal basis v_1,v_2,\ldots,v_k for V. Define $x\in\mathbb{R}^n$ by $x_i=\sum_j|\operatorname{Proj}_iv_j|^2$ where Proj_iv_j is the orthogonal projection of v_j onto the i^{th} column. Then $\Phi(Y)$ is the orthogonal projection of $\lambda+x\in\mathbb{R}^n$ onto $\mathfrak{t}_{\mathbb{R}}$.

2.3. **Definition of MV-cycles.** We recall (an essentially equivalent version of) the Mirković-Vilonen definition of MV-cycles. Given any lattice Y, let us define its highest coweight λ and its lowest coweight μ to be such that $\underline{\lambda}$ results from using the torus action to "flow left" and $\underline{\mu}$ results from using it to "flow right"; that is, we require that $\lim_{z\to 0}(1,z,z^2,\ldots,z^{n-1})\cdot Y=\underline{\lambda}$ and $\lim_{z\to \infty}(1,z,z^2,\ldots,z^{n-1})\cdot Y=\underline{\mu}$. For example, in Figure 1, the first lattice corresponds to the highest coweight for the second lattice. Let S_{λ} denote the set of lattices with highest coweight λ and T_{μ} denote the set of lattices with lowest coweight μ . So S_{λ} and T_{μ} are stable and unstable manifolds with respect to a Morse function given by a certain component of the moment map Φ .

We define the MV-cycles with highest coweight λ and lowest coweight μ to be the irreducible components of the closure of $S_{\lambda} \cap T_{\mu}$. This intersection is pure dimensional [MV2].

Notice that a coweight ν shifts S_{λ} and T_{μ} to $S_{\lambda+\nu}$ and $T_{\mu+\nu}$ respectively. Hence, the MV-cycles with highest coweight λ and lowest coweight μ are isomorphic to the MV-cycles with highest coweight $\lambda + \nu$ and lowest coweight $\mu + \nu$. We will say that two MV-cycles are equivalent if one is a shift of another; it will often be convenient to work with the equivalence classes rather than with the cycles themselves.

2.4. **Kostant parameter set.** Here we recall the definition of Kostant partitions in type A, which are just formal sums of positive roots. But it is convenient to develop an unusual notation for these as pictures of loops on a Dynkin diagram, for two reasons: (1) With this notation, it is easy to see that each Kostant partition may naturally be given the structure of a partially ordered set, which will be useful. (2) These pictures make it easy to visualize the relationship between Kostant partitions and lattices in the affine Grassmannian.

Recall that the Dynkin diagram in type A_{n-1} is n-1 dots in a row, one for each simple root α_i (connected by line segments, which we will not draw). We denote a positive root by a loop around a sequence of consecutive dots in the Dynkin diagram, and we call the number of dots it encloses the length of the loop. The simple roots are the loops of length 1. The other positive roots are loops of length ≥ 2 ; each is the sum of the simple roots corresponding to the dots

it encloses. A Kostant picture is a picture of the Dynkin diagram together with a finite number of such loops. We draw the loops so that if the dots contained in one loop are a proper subset of the dots contained in another, then the one is encircled by the other; if two loops contain precisely the same dots, we still draw one encircled by the other. In this way, the loops in a Kostant picture are partially ordered by encirclement. We write $L \subset L'$ if loop L' encircles loop L, and $L \subseteq L'$ if either L = L' or L' encircles L. The number of loops in a Kostant picture \mathbf{p} is denoted $|\mathbf{p}|$.

The Kostant parameter set \mathbb{K} is the collection of all Kostant pictures. Figure 2 shows examples of Kostant pictures for n=6. If α_{ij} is the root that is the sum of simple roots $\alpha_i + \cdots + \alpha_{j-1}$, then these three pictures represent the following Kostant partitions: $\alpha_3 + \alpha_{24} + \alpha_{35} + \alpha_{25} + 3\alpha_{46}$, $\alpha_2 + \alpha_{13} + \alpha_{35} + \alpha_{46}$ and $\alpha_5 + \alpha_{24} + \alpha_{14} + \alpha_{25} + \alpha_{26}$.







Figure 2. Examples of Kostant pictures.

It will be useful to imagine each of the n-1 dots of the Dynkin diagram as lying on the boundary between two of the n columns of basis vectors for X. Then, associated to each loop L, will be the vector subspace V_L spanned by the columns the loop passes through. To be precise, for each $i=1,2,\ldots,n$, let V_i be the span of the basis vectors $t^j e_i$ in the i^{th} column. If we number the dots of the Dynkin diagram $1,2,\ldots,n-1$ and if a loop L contains dots $\ell,\ell+1,\ldots,r-1$ then $V_L=V_{\ell,r}=V_\ell\oplus V_{\ell+1}\oplus\cdots\oplus V_r$. We say that the loop L has its left end at column ℓ , its right end at column r, and passes through columns $\ell,\ell+1,\ldots,r$.

By an extended Kostant picture we mean an ordered pair $\tilde{\mathbf{p}} = (\mathbf{p}, \eta)$ where \mathbf{p} is a Kostant picture and η is a coweight. (It is sometimes convenient to use this terminology: given an extended Kostant picture $\tilde{\mathbf{p}}$ and an integer m that's less than or equal to each component of η , we say that, relative to level m, $\tilde{\mathbf{p}}$ has $\eta_i - m$ loops of length zero, or zero loops, in column i.) Let $\tilde{\mathbb{K}}$ denote the collection of all extended Kostant pictures.

Given an extended Kostant picture $\tilde{\mathbf{p}}$, we define its highest coweight λ and lowest coweight μ : we set λ_i equal to η_i plus the number of loops of \mathbf{p} with left end at column i; we set μ_i equal to η_i plus the number of loops of \mathbf{p} with right end at column i. Observe that \mathbf{p} together with any one of η, λ, μ uniquely determines $\tilde{\mathbf{p}}$.

2.5. Parametrization of MV-cycles. We partition \mathbf{Gr} into pieces parametrized by $\tilde{\mathbb{K}}$, by defining a map which sends $Y \in \mathbf{Gr}$ to $(\mathbf{p}(Y), \eta(Y)) \in \tilde{\mathbb{K}}$.

Let $\mathbf{p}(Y)$ be the Kostant picture where the number of loops encircling dots $i, i+1, \ldots, j-1$ is equal to the dimension of the vector space $(Y \cap V_{i,j})/((Y \cap V_{i+1,j}) + (Y \cap V_{i,j-1}))$. Let $\eta(Y)$ be the coweight associated to the span of all basis vectors $t^j e_i$ contained in Y.

Then using [AK] terminology we say that the lattice Y is weakly compatible to the extended Kostant picture $\tilde{\mathbf{p}} = (\mathbf{p}(Y), \eta(Y))$, and we denote by $M(\tilde{\mathbf{p}})$ the closure of the set of lattices weakly compatible to $\tilde{\mathbf{p}}$.

Theorem 1. [AK] The map $\tilde{\mathbf{p}} \to M(\tilde{\mathbf{p}})$ is a one-to-one correspondence between extended Kostant pictures and MV-cycles.

Under the equivalence relation by shifts, we have that $M((\mathbf{p}, \eta_1))$ is equivalent to $M((\mathbf{p}, \eta_2))$ by a shift of $\eta_2 - \eta_1$. If we let $M(\mathbf{p})$ denote the equivalence class $\{M(\tilde{\mathbf{p}}) \mid \tilde{\mathbf{p}} = (\mathbf{p}, \eta) \text{ for some } \eta\}$, then we also have a one-to-one correspondence $\mathbf{p} \to M(\mathbf{p})$ between Kostant pictures and equivalence classes of MV-cycles.

Remark. Kamnitzer [K] generalized Theorem 1 to other types. In particular, for any G and a choice of a reduced word \mathbf{i} for the longest element of the Weyl group of G, he parametrized the set of equivalence classes of MV-cycles by the Kostant parameter set. The parametrization of Theorem 1 coincides with Kamnitzer's parametrization for $G = GL_n$ and a certain choice of \mathbf{i} .

The relationship between a Kostant picture and the lattices inside the corresponding MV-cycle can perhaps be seen more clearly by choosing a basis for the lattice: Suppose that $\tilde{\mathbf{p}} = (\mathbf{p}, \eta)$ is an extended Kostant picture and Y is a lattice with basis $\{y_L\}_{L \in \mathbf{p}} \cup \{t^j e_i \mid t^j e_i \in \underline{\eta}\}$ satisfying $y_L \in V_L$ and $t \cdot y_L \in \text{span}\{y_{L'} \mid L' \subset L\} + (\underline{\eta} \cap V_L)$. Then it is not hard to see that $Y \in M(\tilde{\mathbf{p}})$. Moreover, if $L \in \mathbf{p}$ has left end at column ℓ then the maximum value of j for which y_L can have a nonzero component at $t^{-j}e_\ell$ is equal to $\eta_\ell + |\{L' \in \mathbf{p} \mid L' \subseteq L \text{ and } L' \text{ has left end at column } \ell\}|$; similarly for the right end r of L. Then Y is weakly compatible to $\tilde{\mathbf{p}}$ provided that each y_L actually does have a nonzero component at this maximum position, for both its left and right end.

For example, let Y_2 and Y_3 denote the second and third lattices in Figure 1 and let \mathbf{p}_2 and \mathbf{p}_3 denote the second and third Kostant pictures in Figure 2. Then Y_2 is weakly compatible to $\tilde{\mathbf{p}}_2 = (\mathbf{p}_2, (1, -2, -2, -3, -2, -2))$ and Y_3 is weakly compatible to $\tilde{\mathbf{p}}_3 = (\mathbf{p}_3, (-3, -3, -3, -3, -2, 0))$.

It's also not hard to see from this description that each $M(\tilde{\mathbf{p}})$ has dimension equal to the sum of the lengths of the loops of \mathbf{p} [AK].

2.6. **Richardson varieties.** As an important special case of MV-cycles one finds all of the *Richardson varieties* in a Grassmannian—i.e. intersections of opposite Schubert varieties. These are exactly the MV-cycles corresponding to Kostant pictures where no loop is encircled by another.

Suppose that **p** is such a Kostant picture; note that its loops are naturally ordered from left to right, say $L_1, L_2, \ldots, L_{|\mathbf{p}|}$. Consider the MV-cycle $M = M(\tilde{\mathbf{p}})$

corresponding to $\tilde{\mathbf{p}} = (\mathbf{p}, \eta)$, where, to be definite, choose $\underline{\eta} = X_0$. Let $X_1 = t^{-1} \cdot X_0$. Then each lattice Y in M satisfies $X_0 \subseteq Y \subseteq X_1$ with $\dim(Y/X_0) = |\mathbf{p}|$. In fact, the MV-cycle M is the intersection of the two Schubert varieties S_ℓ and S_r determined by the left and right ends of the loops of \mathbf{p} . Let $\ell_1 < \ell_2 < \cdots < \ell_{|\mathbf{p}|}$ denote the left ends and $r_1 < r_2 < \cdots < r_{|\mathbf{p}|}$ the right ends. Then S_ℓ is the collection of lattices Y (still with $X_0 \subseteq Y \subseteq X_1$ and $\dim(Y/X_0) = |\mathbf{p}|$) such that $\dim(Y \cap V_{\ell_i,n})/(X_0 \cap V_{\ell_i,n}) \ge |\mathbf{p}| - i + 1$ for $i = 1, 2, \ldots, |\mathbf{p}|$; and S_r is the collection of lattices Y such that $\dim(Y \cap V_{1,r_i})/(X_0 \cap V_{1,r_i}) \ge i$ for $i = 1, 2, \ldots, |\mathbf{p}|$.

2.7. **MV-polytopes.** MV-polytopes are defined as the moment map images of MV-cycles. Included among them are all of the Weyl polytopes from representation theory (the convex hulls of a Weyl group orbit through a coweight). An MV-polytope $\Phi(M(\tilde{\mathbf{p}}))$ is a translation of a Weyl polytope if and only if the number of copies of a particular loop in \mathbf{p} depends only on the length of the loop; for instance if \mathbf{p} contains five loops of length 1 around the second dot of the Dynkin diagram, then it must also contain exactly five loops of length 1 around each of the other dots as well.

One reason to be interested in MV-polytopes is that they are closely related to representation theory. For instance, as shown by the first author [A2], both weight multiplicities and tensor product multiplicities may be expressed as the number of MV-polytopes fitting inside a certain region. Recent work of Kamnitzer [K] shows that this is really an alternative way of looking at Bernstein-Zelevinsky's computation of tensor product multiplicities as the number of lattice points inside a convex polytope [BZ].

In [AK] all MV-polytopes are computed in type A, by constructing an explicit map from the Weyl group onto the vertices of each polytope. More recently Kamnitzer [K] has generalized this, using Bernstein and Zelevinsky's work [BZ] to compute MV-polytopes in any type.

3. Convolution of MV-cycles.

We will explain how to define a product structure on MV-cycles. More specifically, given two MV-cycles M and N we will describe a family $M*_{\mathbb{C}}N$ fibered over \mathbb{C} , whose fibers over \mathbb{C}^* are canonically isomorphic to $M\times N$ and whose fiber over 0 has top-dimensional component equal to the union of other MV-cycles K with multiplicities c_{MN}^K . Then the convolution product of M and N is defined to be

$$M*N = \sum c_{MN}^K K.$$

We will sketch a definition of this product for all types in Section 3.1 and then concentrate on a lattice model version of it in type A. Most of this section can be omitted on first reading, as we will summarize the important results at the beginning of the next section; only their statements as well as the partial order defined in Section 3.4 are needed to understand the rest of the paper.

3.1. Beilinson-Drinfeld Grassmannian and convolution. The key object used in the definition of the convolution product of MV-cycles is the Beilinson-Drinfeld Grassmannian $\mathbf{Gr}*_{\mathbb{C}}\mathbf{Gr}$ over \mathbb{C} of the group \mathbf{G} . It is an infinite-dimensional space which fibers over \mathbb{C} . Just as for the affine Grassmannian, the Beilinson-Drinfeld Grassmannian can be approximated by finite-dimensional algebraic varieties, so that in most cases we will be working in a finite-dimensional space. All fibers of $\mathbf{Gr}*_{\mathbb{C}}\mathbf{Gr}$ over \mathbb{C}^* are isomorphic to $\mathbf{Gr} \times \mathbf{Gr}$ while the fiber over 0 is isomorphic to \mathbf{Gr} . A precise definition of $\mathbf{Gr}*_{\mathbb{C}}\mathbf{Gr}$ along with its many properties is discussed in [BD], [MV2], [Na1]. Instead of repeating this definition we will give a lattice description of $\mathbf{Gr}*_{\mathbb{C}}\mathbf{Gr}$ in type A a little later.

Both the affine Grassmannian and the Beilinson-Drinfeld Grassmannian can be defined for any Borel subgroup B of G. It turns out that the sets of points of the affine Grassmannians for G and B are in bijection. However, their topological and algebraic structures differ. In particular, the connected components for the affine Grassmannian of B are the spaces S_{λ} (similarly T_{μ} for the opposite Borel B⁻). Analogously, the sets of points of the Beilinson-Drinfeld Grassmannians for G and B are identical, but the topological and algebraic structures differ. Each irreducible component of the Beilinson-Drinfeld Grassmannian for B is a family $S_{\lambda}*_{\mathbb{C}}S_{\nu}$ ($T_{\mu}*_{\mathbb{C}}T_{\gamma}$ for B⁻) which fibers over \mathbb{C} such that its fibers over \mathbb{C}^* are naturally isomorphic to $S_{\lambda} \times S_{\nu}$ (respectively $T_{\mu} \times T_{\gamma}$) and its fiber over 0 is isomorphic to $S_{\lambda+\nu}$ (respectively $T_{\mu+\gamma}$). We will treat both $S_{\lambda}*_{\mathbb{C}}S_{\nu}$ and $T_{\mu}*_{\mathbb{C}}T_{\gamma}$ as subspaces of $\mathbf{Gr}*_{\mathbb{C}}\mathbf{Gr}$.

This allows one to define the family $M*_{\mathbb{C}}N$ as the irreducible component of the closure of the intersection $S_{\lambda}*_{\mathbb{C}}S_{\nu}\cap T_{\mu}*_{\mathbb{C}}T_{\gamma}$ whose fibers over \mathbb{C}^* are $M\times N$. Then the fiber of $M*_{\mathbb{C}}N$ over 0 lies inside the intersection $S_{\lambda+\nu}\cap T_{\mu+\gamma}$, and its top dimensional components are MV-cycles for $\overline{S_{\lambda+\nu}\cap T_{\mu+\gamma}}$ which are used to define the product (3.1).

If instead of using M and N here we used their shifts $\alpha \cdot M$ and $\beta \cdot N$, then every fiber of $(\alpha \cdot M) *_{\mathbb{C}} (\beta \cdot N)$ would be a shift of the corresponding fiber of $M *_{\mathbb{C}} N$. Hence (3.1) defines a product on equivalence classes of MV-cycles.

The convolution product is commutative and associative. The commutativity comes from the \mathbb{Z}_2 symmetry of $\mathbf{Gr} *_{\mathbb{C}} \mathbf{Gr}$ which identifies $M *_C N$ with $N *_{\mathbb{C}} M$. The associativity follows from the fact that it is possible to define the product of three MV-cycles using a more general version of the Beilinson-Drinfeld Grassmannian fibered over \mathbb{C}^2 . (See, for example, [MV2] or [Na1]. But note that they define the Beilinson-Drinfeld Grassmannian in more generality, where it is fibered over $C \times C$ for some algebraic curve C. Our $\mathbf{Gr} *_{\mathbb{C}} \mathbf{Gr}$ is a restriction of this to a copy of \mathbb{C} inside $C \times C$ that transversely intersects the diagonal of $C \times C$ at zero.)

Remark. It is possible to define an action of the torus T on $\mathbf{Gr} *_{\mathbb{C}} \mathbf{Gr}$ so that it acts diagonally on every fiber over \mathbb{C}^* and is isomorphic to the usual T action at the fiber over 0. There is a moment map on $\mathbf{Gr} *_{\mathbb{C}} \mathbf{Gr}$ associated to this

action. Then the moment map image (and more generally the Duistermaat-Heckman measure) of all fibers of $M*_{\mathbb{C}}N$ are the same. Since every fiber over \mathbb{C}^* is isomorphic to $M \times N$, its moment map image is the Minkowski sum of the MV-polytopes $\Phi(M)$ and $\Phi(N)$. Hence every MV-polytope $\Phi(K)$ for $c_{MN}^K \neq 0$ lies within this Minkowski sum. This is a very useful tool in proving that certain coefficients from (3.1) are zero. Also if we use the standard definition of convolution of measures on $\mathfrak{t}_{\mathbb{R}}$, the Duistermaat-Heckman measures of MV-cycles satisfy an equation exactly like (3.1).

3.2. Lattice model of the Beilinson-Drinfeld Grassmannian in type A. For $G = GL_n$ several lattice models for the Beilinson-Drinfeld Grassmannian are known. One is due to Nadler [Na2]. We present another model, which only constructs certain finite-dimensional subfamilies inside of $\mathbf{Gr} *_{\mathbb{C}} \mathbf{Gr}$; but since any family $M *_{\mathbb{C}} N$ will embed into one of these, it will be sufficient for the study of convolution of MV-cycles.

Let \mathbf{Gr}_0 be the subspace of \mathbf{Gr} containing those lattices Y for which $X_0 \subset Y$. Clearly, each connected component of \mathbf{Gr}_0 is a finite-dimensional subvariety of \mathbf{Gr} . We will describe the family $\mathbf{Gr}_0 *_{\mathbb{C}} \mathbf{Gr}_0$, which is a subfamily of $\mathbf{Gr} *_{\mathbb{C}} \mathbf{Gr}$. Since any two MV-cycles M and N can be shifted to lie in \mathbf{Gr}_0 , their convolution $M*_{\mathbb{C}}N$ can be shifted to lie in $\mathbf{Gr}_0*_{\mathbb{C}}\mathbf{Gr}_0$. Hence the convolution product can be studied within $\mathbf{Gr}_0*_{\mathbb{C}}\mathbf{Gr}_0$ instead of the whole Beilinson-Drinfeld Grassmannian.

Let X be the vector space of all formal—possibly infinite—linear combinations of vectors $t^j e_i$ with $-\infty < j < \infty$ and $1 \le i \le n$. As before, X_0 is the span of all $t^j e_i$ with $j \ge 0$, when only finite linear combinations are allowed. For a complex number $s \ne 0$, we denote the sum

(3.2)
$$\frac{1}{t}(1 + \frac{s}{t} + \frac{s^2}{t^2} + \dots + \frac{s^k}{t^k} + \dots)$$

by $\frac{1}{t-s}$. Let X^s be the span of X_0 and the vectors $\frac{e_i}{(t-s)^j}$ for j > 0 (where we treat $\frac{1}{(t-s)^j}$ as the j^{th} power of the series (3.2)). A subspace Y of X^s will be called an s-lattice if it is invariant under multiplication by (t-s), it contains X_0 , and for a large N it is contained inside the space spanned by X_0 and the vectors $\frac{e_i}{(t-s)^j}$ for $N \geq j > 0$.

Now consider the variety \mathcal{F} of all the subspaces Y of X of the form $Y = Y_1 + Y_2$, where Y_1 is a 0-lattice (namely an ordinary lattice containing X_0) and Y_2 is an s-lattice for $s \neq 0$. Clearly, \mathcal{F} is a union of finite-dimensional algebraic varieties and an s-fibration over $\mathbb{C} - \{0\}$ whose fibers are isomorphic to $\mathbf{Gr}_0 \times \mathbf{Gr}_0$.

Finally, consider the space \mathcal{X} of subspaces Z of X with $X_0 \subset Z$ and $\dim(Z/X_0) < \infty$. The space $\mathbf{Gr}_0 *_{\mathbb{C}} \mathbf{Gr}_0$ is now defined to be the closure of \mathcal{F} within \mathcal{X} .

3.3. **Example.** This lattice model for the Beilinson-Drinfeld Grassmannian is useful for finding concrete families $M*_{\mathbb{C}}N$ which deform $M\times N$ into unions of MV-cycles. Sometimes the lattice model can be used to guess correctly which

coefficients on the right side of the equation (3.1) are not zero. However, it is not well suited for computing these coefficients. In particular it is very hard to see using this model whether a nonzero coefficient is equal to 1 or to some other integer. Let us present an illustration.

Assume n=3. Let \mathbf{p} be the Kostant partition α_1 and \mathbf{q} be the Kostant partition α_2 . Define $\tilde{\mathbf{p}}=(\mathbf{p},\underline{0})$, $\tilde{\mathbf{q}}=(\mathbf{q},\underline{0})$, where $\underline{0}=(0,0,0)$. Then $M(\tilde{\mathbf{p}})*M(\tilde{\mathbf{q}})=M(\tilde{\mathbf{r}})+M(\tilde{\mathbf{s}})$, where $\tilde{\mathbf{r}}=(\alpha_1+\alpha_2,\underline{0})$ and $\tilde{\mathbf{s}}=(\alpha_{13},(0,1,0))$. In terms of the Kostant pictures we can write this equation as

$$(3.3) M(\bullet \bullet) * M(\bullet \bullet) = M(\bullet \bullet) + M(\bullet \bullet).$$

Both MV-cycles $M(\tilde{\mathbf{p}})$ and $M(\tilde{\mathbf{q}})$ are projective lines \mathbb{CP}^1 , while $M(\tilde{\mathbf{r}})$ and $M(\tilde{\mathbf{s}})$ are projective planes \mathbb{CP}^2 intersecting along a \mathbb{CP}^1 . Hence the family $M*_{\mathbb{C}}N$ degenerates $\mathbb{CP}^1 \times \mathbb{CP}^1$ to two copies of \mathbb{CP}^2 intersecting along \mathbb{CP}^1 . There are many ways of visualizing this family. For example, let W be a four-dimensional vector space with basis w_1, w_2, w_3, w_4 . Then the s-fiber of this family can be identified with the pair of lines, one sitting inside the span of w_1 and w_2 and the other inside the span of $w_2 + sw_4$ and w_3 . The zero fiber contains those two-dimensional planes inside W which either contain the vector w_2 or sit inside the span of w_1, w_2, w_3 .

We will use the lattice model to show that both MV-cycles appearing on the right side of equation (3.3) should indeed be there. However, checking that both multiplicities are 1 and that there are no other terms on the right side requires other techniques (for example, moment map images or, more generally, Duistermaat-Heckman measures of MV-cycles).

Consider the space $\mathcal{F}(\tilde{\mathbf{p}}, \tilde{\mathbf{q}})$ which contains all the vector subspaces $Y \subset \widetilde{X}$ spanned by X_0 and two vectors y_1 and y_2 of the form $y_1 = at^{-1}e_1 + bt^{-1}e_2$ and $y_2 = \frac{ce_2}{t-s} + \frac{de_3}{t-s}$ for some $s \neq 0$. The closure of $\mathcal{F}(\tilde{\mathbf{p}}, \tilde{\mathbf{q}})$ is the family $M(\tilde{\mathbf{p}}) *_{\mathbb{C}} M(\tilde{\mathbf{q}})$.

To show that a certain MV-cycle K enters the product $M(\tilde{\mathbf{p}})*M(\tilde{\mathbf{q}})$ with non-trivial coefficient it is enough to construct a one-parameter family Y(s) inside $M(\tilde{\mathbf{p}})*_{\mathbb{C}}M(\tilde{\mathbf{q}})$ such that Y(s) lies inside the s-fiber of $M(\tilde{\mathbf{p}})*_{\mathbb{C}}M(\tilde{\mathbf{q}})$, and Y(0) is a generic point of K. We will construct two such families in our example, one for each MV-cycle on the right side of equation (3.3).

Given any numbers a, b, c, d, we can define the family Y(s) to be the span of $X_0, y_1(s) = at^{-1}e_1 + bt^{-1}e_2$ and $y_2(s) = \frac{ce_2}{t-s} + \frac{de_3}{t-s}$. Then Y(0) is the span of $X_0, y_1(0) = at^{-1}e_1 + bt^{-1}e_2$, and $y_2(0) = ct^{-1}e_2 + dt^{-1}e_3$, which is a generic point in $M(\tilde{\mathbf{r}})$.

Similarly, given numbers a,d and $b \neq 0$, let Y(s) be spanned by $X_0, y_1(s) = sat^{-1}e_1 + bt^{-1}e_2$, and $y_2(s) = \frac{be_2}{t-s} + s\frac{de_3}{t-s}$ for $s \neq 0$. Then Y(0) is the span of X_0 , $y_1(0) = bt^{-1}e_2$, and $\lim_{s\to 0} \frac{y_1(s)}{s} - \frac{y_2(s)}{s} = at^{-1}e_1 - bt^{-2}e_2 - dt^{-1}e_3$, which is a generic point in $M(\tilde{\mathbf{s}})$.

3.4. A partial order on MV-cycles and its relation to the convolution product. In this section we present a property of the convolution of MV-cycles. Our statement of this property is specific to type A and the proof uses the lattice model for the Beilinson-Drinfeld Grassmannian. However, it should be possible to use the techniques developed by Kamnitzer [K] to generalize this to other types, as will be discussed below.

Let us start by defining a partial order on extended Kostant pictures. For an extended Kostant picture $\tilde{\mathbf{p}} = (\mathbf{p}, \eta)$, fix an m such that each $\eta_i > m$, and think of $\tilde{\mathbf{p}}$ as having $\eta_i - m$ zero loops in column i. Consider loops A, B in $\tilde{\mathbf{p}}$ which overlap, in the sense that neither encircles the other and they pass through at least one common column. Define a new Kostant picture $\tilde{\mathbf{p}}'$ by replacing A and B by their union $A \cup B$ (i.e. the loop passing though the columns which either A or B pass through) and intersection $A \cap B$ (i.e. the loop passing though the columns which both A and B pass through; if there is only one such column, say i, then this is a zero loop, in which case η_i is increased by 1). Call this operation $\tilde{\mathbf{p}} \to \tilde{\mathbf{p}}'$ the fusion of loops A and B. Notice that fusion does not change the highest and lowest coweights of the extended Kostant picture.

For example, consider the extended Kostant picture $\tilde{\mathbf{r}}$ from Section 3.3. It contains two length one loops and their fusion is a length two loop and a length zero loop. So this fusion produces the extended Kostant picture $\tilde{\mathbf{s}}$.

Define a partial order on $\tilde{\mathbb{K}}$ by saying that $\tilde{\mathbf{p}} \geq \tilde{\mathbf{q}}$ if $\tilde{\mathbf{q}}$ is produced from $\tilde{\mathbf{p}}$ by a sequence of fusions. To see that this partial order is well-defined it is enough to show that a sequence of nontrivial fusions of an extended Kostant picture $\tilde{\mathbf{p}}$ cannot produce $\tilde{\mathbf{p}}$ again. Fix an integer m as before. For $1 \leq i \leq j \leq n$ define $d_{ij}(\tilde{\mathbf{p}})$ to be the number of loops (including the zero loops) of $\tilde{\mathbf{p}}$ which lie between columns i and j; in other words $d_{ij}(\tilde{\mathbf{p}})$ is the number of loops of $\tilde{\mathbf{p}}$ which do not pass through any of the columns $1, \ldots, i-1, j+1, \ldots, n$. Let $D(\tilde{\mathbf{p}})$ denote the collection of numbers $d_{ij}(\tilde{\mathbf{p}})$. We say that $D(\tilde{\mathbf{p}}) \leq D(\tilde{\mathbf{q}})$ if $d_{ij}(\tilde{\mathbf{p}}) \leq d_{ij}(\tilde{\mathbf{q}})$ for every i, j; and $D(\tilde{\mathbf{p}}) < D(\tilde{\mathbf{q}})$ if $D(\tilde{\mathbf{p}}) \leq D(\tilde{\mathbf{q}})$ and $d_{ij}(\tilde{\mathbf{p}}) < d_{ij}(\tilde{\mathbf{q}})$ for some i, j.

It is clear that if $\tilde{\mathbf{q}}$ is produced out of $\tilde{\mathbf{p}}$ by a sequence of nontrivial fusions, then $D(\tilde{\mathbf{q}}) > D(\tilde{\mathbf{p}})$, which proves that the partial order on $\tilde{\mathbb{K}}$ is well-defined. Surprisingly the converse is also true.

Lemma 3.1. Assume $\tilde{\mathbf{p}}$ and $\tilde{\mathbf{q}}$ have the same highest and lowest coweights. Fix a common m to define the numbers $d_{ij}(\tilde{\mathbf{p}})$ and $d_{ij}(\tilde{\mathbf{q}})$. Then $D(\tilde{\mathbf{q}}) > D(\tilde{\mathbf{p}})$ if and only if $\tilde{\mathbf{q}} < \tilde{\mathbf{p}}$.

Proof. It remains to show that $D(\tilde{\mathbf{q}}) > D(\tilde{\mathbf{p}})$ implies $\tilde{\mathbf{q}} < \tilde{\mathbf{p}}$. If $\tilde{\mathbf{q}}$ and $\tilde{\mathbf{p}}$ have a common loop, we can remove it and proceed by induction on the number of loops. If all the loops of $\tilde{\mathbf{q}}$ and $\tilde{\mathbf{p}}$ are different, pick a loop C inside $\tilde{\mathbf{q}}$ which does not encircle any other loop. Let the left end of C be ℓ and the right end r. Consider all the loops of $\tilde{\mathbf{p}}$ which encircle C, and among these, consider only those that do not encircle any other loops which encircle C; these are ordered in a sequence from left to right. Moreover there must be at least two of them: in particular,

since the lowest and highest coweights of $\tilde{\mathbf{p}}$ and $\tilde{\mathbf{q}}$ are the same, there is one with right end r and another with left end ℓ . Let A and B be any two consecutive loops in this sequence, and let $\tilde{\mathbf{p}}'$ be produced from $\tilde{\mathbf{p}}$ by the fusion of A and B. Then it is not hard to check that $D(\tilde{\mathbf{q}}) \geq D(\tilde{\mathbf{p}}')$. If $D(\tilde{\mathbf{q}}) = D(\tilde{\mathbf{p}}')$ then $\tilde{\mathbf{q}} = \tilde{\mathbf{p}}'$ and we are done; otherwise $D(\tilde{\mathbf{q}}) > D(\tilde{\mathbf{p}}')$ and we proceed by induction with $\tilde{\mathbf{p}}$ replaced by $\tilde{\mathbf{p}}'$.

Let **p** and **q** be two Kostant pictures. Define $\mathbf{p} + \mathbf{q}$ to be the Kostant picture containing the loops of both **p** and **q**; in other words, we just add the Kostant partitions. If $\tilde{\mathbf{p}} = (\mathbf{p}, \eta)$ and $\tilde{\mathbf{q}} = (\mathbf{q}, \eta')$, then define $\tilde{\mathbf{p}} + \tilde{\mathbf{q}} = (\mathbf{p} + \mathbf{q}, \eta + \eta')$.

Proposition 3.2. If
$$M(\tilde{\mathbf{p}})*M(\tilde{\mathbf{q}}) = \sum c_{\tilde{\mathbf{r}}}M(\tilde{\mathbf{r}})$$
 then $c_{\tilde{\mathbf{r}}} = 0$ for $\tilde{\mathbf{r}} > \tilde{\mathbf{p}} + \tilde{\mathbf{q}}$.

Proof. Without loss of generality we may assume that $\tilde{\mathbf{p}} = (\mathbf{p}, \underline{0})$ and $\tilde{\mathbf{q}} = (\mathbf{q}, \underline{0})$. For the rest of the proof fix m = 0 when defining the d_{ij} 's for extended Kostant pictures.

For a lattice Y containing X_0 define D(Y) to be the set of numbers $d_{ij}(Y) = \dim(Y \cap V_{ij}/X_0 \cap V_{ij})$. (Recall that V_{ij} is the span of columns i through j inside X.) It is easy to see that $D(Y) \geq D(\tilde{\mathbf{p}})$ whenever $Y \in M(\tilde{\mathbf{p}})$ and $D(Y) \geq D(\tilde{\mathbf{q}})$ whenever $Y \in M(\tilde{\mathbf{q}})$. Moreover, generically these inequalities are equalities.

Now let Y be a point in $M(\tilde{\mathbf{p}})*_{\mathbb{C}}M(\tilde{\mathbf{q}})$. We can again define D(Y) to be the set of numbers $d_{ij}(Y) = \dim(Y \cap \widetilde{V}_{ij}/X_0 \cap \widetilde{V}_{ij})$, where \widetilde{V}_{ij} is now the span of columns i though j of \widetilde{X} . If Y is in a fiber of $M*_{\mathbb{C}}N$ over \mathbb{C}^* then $Y = Y_1 + Y_2$ with $Y_1 \in M(\tilde{\mathbf{p}})$ and $Y_2 \in M(\tilde{\mathbf{q}})$. Hence

$$D(Y) = D(Y_1) + D(Y_2) \ge D(\tilde{\mathbf{p}}) + D(\tilde{\mathbf{q}}) = D(\tilde{\mathbf{p}} + \tilde{\mathbf{q}}).$$

Since the functions d_{ij} are lower semi-continuous, we conclude that $D(Y) \geq D(\tilde{\mathbf{p}} + \tilde{\mathbf{q}})$ for any $Y \in M(\tilde{\mathbf{p}}) *_{\mathbb{C}} M(\tilde{\mathbf{q}})$, and in particular for Y in the zero fiber. Applying Lemma 3.1 finishes the proof.

Remark. We believe that Proposition 3.2 holds in a more general situation. Namely as shown by Kamnitzer [K], once a reduced word for the longest element of the Weyl group of G is chosen, it is possible to parametrize MV-cycles by extended Kostant partitions. Moreover, Kamnitzer constructed functions which naturally generalize the numbers d_{ij} in the above proof. So, we expect that Proposition 3.2 holds for any type as soon as the proper partial order on extended Kostant partitions is defined. The proof of the more general statement should rely on the semi-continuity of Kamnitzer's analogues of the numbers d_{ij} .

3.5. The leading term of the convolution product. The converse of Proposition 3.2 is false: one may have $c_{\tilde{\mathbf{r}}} = 0$ even when $\tilde{\mathbf{r}} < \tilde{\mathbf{p}} + \tilde{\mathbf{q}}$. But not for $\tilde{\mathbf{r}} = \tilde{\mathbf{p}} + \tilde{\mathbf{q}}$:

Proposition 3.3. If $M(\tilde{\mathbf{p}})*M(\tilde{\mathbf{q}}) = \sum c_{\tilde{\mathbf{r}}}M(\tilde{\mathbf{r}})$ then $c_{\tilde{\mathbf{p}}+\tilde{\mathbf{q}}} \neq 0$.

Conjecture 3.4. The coefficient $c_{\tilde{\mathbf{p}}+\tilde{\mathbf{q}}}$ is always equal to 1.

The rest of this section is devoted to the proof of Proposition 3.3 using the lattice model.

Without loss of generality $\tilde{\mathbf{p}} = (\mathbf{p}, \underline{0})$ and $\tilde{\mathbf{q}} = (\mathbf{q}, \underline{0})$. Let Z be a lattice weakly compatible to $\tilde{\mathbf{p}} + \tilde{\mathbf{q}}$. It will be enough to prove that for each such Z there exists a one-parameter family Z(s) inside $M(\tilde{\mathbf{p}}) *_{\mathbb{C}} M(\tilde{\mathbf{q}})$ with Z(0) = Z and Z(s) sitting inside the s-fiber of $M(\tilde{\mathbf{p}}) *_{\mathbb{C}} M(\tilde{\mathbf{q}})$.

The lattice Z has a basis $\{y_L\}_{L\in\tilde{\mathbf{p}}+\tilde{\mathbf{q}}}$ indexed by the loops of $\tilde{\mathbf{p}}+\tilde{\mathbf{q}}$ as described in Section 2.5. (Here we treat each vector t^je_i as the basis vector associated to a zero loop.) This basis satisfies the property that each vector ty_L lies in the span of the vectors associated to loops encircled by L. So we can aways uniquely write

(3.4)
$$ty_L = \sum_{L' \subset L} a_L^{L'} y_{L'}.$$

Notice that if we count all the zero loops there are infinitely many loops encircled by L; nevertheless the summation on the right side of (3.4) is finite.

Let us now decompose the Kostant picture $\mathbf{p} + \mathbf{q}$ into \mathbf{p} and \mathbf{q} : color every loop in $\mathbf{p} + \mathbf{q}$ either red or green, so that the red loops give \mathbf{p} and the green loops give \mathbf{q} .

Define vectors $z_L = z_L(s)$ inside \widetilde{X} by

$$z_L = \frac{1}{t-s} \sum_{L' \subset L} a_L^{L'} z_{L'}$$

if L is red and by

$$z_L = \frac{1}{t} \sum_{L' \subset L} a_L^{L'} z_{L'}$$

if L is green.

Let Z(s) be the span of all the vectors $z_L(s)$ and X_0 . Clearly, Z = Z(0). So, to prove Proposition 3.3 it is enough to show that Z(s) is inside the s-fiber of $M(\tilde{\mathbf{p}}) *_{\mathbb{C}} M(\tilde{\mathbf{q}})$ for $s \neq 0$.

To show this we will construct vectors $x_L = x_L(s)$ of the form

(3.5)
$$x_L = z_L + \sum_{L' \subset L} b_L^{L'} z_{L'}$$

such that for every red loop L, the vector $(t-s)x_L$ is a linear combination of vectors $x_{L'}$ for $L' \subset L$ a red loop or a zero loop. Similarly, for every green loop L, the vector tx_L will be a linear combination of vectors $x_{L'}$ for $L' \subset L$ a green loop or a zero loop. Then we define $Y_2(s)$ to be spanned by red vectors x_L and X_0 , and $Y_1(s)$ to be spanned by green vectors x_L and X_0 . Clearly, $Z(s) = Y_1(s) + Y_2(s)$, $Y_2(s)$ is an s-lattice inside $M(\tilde{\mathbf{p}})$, and $Y_1(s)$ is a 0-lattice inside $M(\tilde{\mathbf{q}})$, proving that Z(s) is inside the s-fiber of $M(\tilde{\mathbf{p}})*_{\mathbb{C}}M(\tilde{\mathbf{q}})$. It remains to construct vectors x_L with the above properties.

Construct the vectors x_L by induction on inclusion of loops. Namely, assume we can construct such vectors for every loop $L' \subset L$ and we will prove that x_L

exists for L. Assume that L is red; an analogous argument can be given if L is

Start by setting $x = z_L$ and write

(3.6)
$$(t-s)x = \sum_{L' \subset L} c^{L'} x_{L'}.$$

If in the above equation $c^{L'} = 0$ for all green loops L' we can just set $x_L = x$. Otherwise, let L^g be one of the largest green loops with $c^{L^g} \neq 0$; in other words if L^g is encircled by a green loop L' then $c^{L'} = 0$. We can then modify x to guarantee $c^{L^g} = 0$ and proceed by induction. Namely we replace x by the vector

$$x' = x + \frac{c^{L_g}}{s} x_{L^g}.$$

Since $L^g \subset L$, the vector x' is still of the form (3.5). And we have

$$(t-s)x' = (t-s)x + \frac{(t-s)c^{L^g}}{s}x_{L^g} = (t-s)x + \frac{tc^{L^g}}{s}x_{L^g} - c^{L^g}x_{L^g}$$
$$= \sum_{L' \subset L} c^{L'}x_{L'} - c^{L^g}x_{L^g} + \frac{c^{L^g}}{s}tx_{L^g}$$

Thus the vector x_{L^g} , as well as all green vectors $x_{L'}$ for loops L' encircling L^g , have coefficient zero when we write (t-s)x' as a linear combination of the $x_{L'}$. This allows us to proceed by induction and construct the vector x of the form (3.5) such that in equation (3.6), $c^{L'} = 0$ whenever L' is green. Then we can set $x_L = x$.

4. Convolution algebra and $\mathbb{C}[N]$

4.1. Properties of the convolution product. Let us recall the results of Section 3 as well as restate them in terms of Kostant pictures rather than extended Kostant pictures. We will assume that $G = SL_n$ or $G = GL_n$ for most of this section.

Equation (3.1) is invariant under shifts. Hence we can define the convolution product on equivalence classes of MV-cycles

(4.1)
$$M(\mathbf{p})*M(\mathbf{q}) = \sum c_{\mathbf{p},\mathbf{q}}^{\mathbf{r}} M(\mathbf{r}).$$

From this point on, whenever we mention MV-cycles we will really mean their equivalence classes, and we will only use the convolution defined in equation (4.1). Also, the partial order on extended Kostant pictures immediately induces one on Kostant pictures: we say $\mathbf{p} > \mathbf{q}$ if $\tilde{\mathbf{p}} = (\mathbf{p}, \eta) > \tilde{\mathbf{q}} = (\mathbf{q}, \eta')$ for some η, η' .

Summarizing Section 3 (in particular Propositions 3.2 and 3.3) we have the following properties of convolution:

- (1) The convolution product is commutative and associative,
- (2) $c_{\mathbf{p},\mathbf{q}}^{\mathbf{r}} = 0$ for $\mathbf{r} > \mathbf{p} + \mathbf{q}$, (3) $c_{\mathbf{p},\mathbf{q}}^{\mathbf{p}+\mathbf{q}} \neq 0$.

The commutativity and associativity (which hold for any G) imply that the \mathbb{C} -span of MV-cycles has an algebra structure. We call this the convolution algebra \mathcal{C}_G for the group G.

4.2. **The isomorphism.** For a group G, denote by N the unipotent radical of a Borel subgroup, and let $\mathbb{C}[N]$ be the algebra of functions on N. Assume G is either SL_n or GL_n , so that N can be identified with upper triangular $n \times n$ matrices with 1's along the diagonal; then $\mathbb{C}[N]$ is the polynomial ring $\mathbb{C}[\mathbf{x}]$ in the matrix entries $\mathbf{x} = \{x_{ij}\}_{1 \le i \le j \le n}$ above the diagonal.

Let us define a homomorphism $\mathcal{I}: \mathbb{C}[N] \to \mathcal{C}_{\mathsf{G}}$ by sending x_{ij} to $M(\mathbf{p}_{ij})$, where \mathbf{p}_{ij} is the Kostant picture with a single loop passing through columns i through j.

Theorem 2. The map \mathcal{I} is an isomorphism.

Proof. For a loop L passing through columns i through j let $x_L = x_{ij}$, and for a Kostant partition \mathbf{p} define the monomial $x^{\mathbf{p}} = \prod_{L \in \mathbf{p}} x_L$. Clearly the monomials $x^{\mathbf{p}}$ provide a basis for $\mathbb{C}[\mathsf{N}]$.

Let us show that \mathcal{I} is surjective. We will prove by induction that every $M(\mathbf{p})$ is in the image of \mathcal{I} . Assume $M(\mathbf{q}) \in \text{Im}(\mathcal{I})$ for every $\mathbf{q} < \mathbf{p}$, and we'll prove the same thing for \mathbf{p} . By properties (2) and (3) of the convolution product

$$\mathcal{I}(x^{\mathbf{p}}) = \prod_{L \in \mathbf{p}} \mathcal{I}(x_L) = aM(\mathbf{p}) + \sum_{\mathbf{q} < \mathbf{p}} a_{\mathbf{q}} M(\mathbf{q})$$

with $a \neq 0$. Using the induction assumption, we conclude that $M(\mathbf{p}) \in \text{Im}(\mathcal{I})$, which proves that \mathcal{I} is surjective.

Let us show that \mathcal{I} is injective. For an integer k, let $\mathbb{C}[\mathbb{N}]^k$ be the span of the monomials $x^{\mathbf{p}}$ with $|\mathbf{p}| \leq k$; similarly, let $\mathcal{C}_{\mathsf{G}}^k$ be the span of the $M(\mathbf{p})$ with $|\mathbf{p}| \leq k$. Clearly $\dim(\mathbb{C}[\mathbb{N}]^k) = \dim(\mathcal{C}_{\mathsf{G}}^k)$. Moreover, by property (2) of the convolution product, $\mathcal{I}^{-1}(\mathcal{C}_{\mathsf{G}}^k) \subset \mathbb{C}[\mathbb{N}]^k$. Hence \mathcal{I} is injective. \square

Let us denote $\mathcal{I}^{-1}(M(\mathbf{p}))$ by $\mathcal{P}_{\mathbf{p}}$ and call it an MV-polynomial. It is easy to see that

$$\mathcal{P}_{\mathbf{p}} = \sum_{\mathbf{q} \le \mathbf{p}} c_{\mathbf{q}}^{\mathbf{p}} x^{\mathbf{q}}$$

with $c_{\mathbf{p}}^{\mathbf{p}} \neq 0$. Computation of the the coefficients $c_{\mathbf{q}}^{\mathbf{p}}$ is a challenging unsolved problem, which we hope has a combinatorial solution. Section 5 is devoted to stating some conjectures about this.

The MV-polynomials form a basis of $\mathbb{C}[N]$, which is expected to coincide with one of the known canonical bases of $\mathbb{C}[N]$. We are not experts in this field and will not speculate which one it is.

Remarks.

- The first author in [A1] defined a conjectural Hopf algebra structure on \mathcal{C}_G and conjectured that $\mathbb{C}[N]$ and \mathcal{C}_G are isomorphic as Hopf algebras. Theorem 2 proves part of that conjecture in type A.
- Feigin, Finkelberg, Kuznetzov, and Mirković [FM, FFKM] gave a geometric construction of the universal enveloping algebra dual to $\mathbb{C}[N]$.
- We believe that Theorem 2 can be generalized to other types by using Kamnitzer's results [K]. As mentioned before, we expect that analogues of properties (2) and (3) of the convolution product for other types can be proved using Kamnitzer's results. Then the above proof of Theorem 2 should work for other types with only minor modifications.
- The partial order on Kostant partitions, viewed as a partial order on the monomials of $\mathbb{C}[\mathbf{x}]$, is a diagonal term order.
- Property (2) of convolution is very similar to the multiplicative property of Lusztig's dual canonical basis proved by Caldero [Ca]. Since MV-polynomials are expected to form one of the known canonical basis of $\mathbb{C}[N]$, this similarity is not surprising.

5. Determinantal formulas for MV-polynomials

A large unsolved problem is to find an explicit formula for every MV-polynomial. This is equivalent to the problem of finding a formula for the convolution of any two MV-cycles (at least insofar as a solution to one problem gives an inductive method for solving the other). Likely, these are difficult problems; all we will be able to do here is to give a mysterious conjectural formula for some MV-polynomials as certain determinants. Even more mysterious is the connection to the Fomin-Zelevinsky theory of cluster algebras, for our conjectures apply only to those MV-polynomials which are also cluster monomials.

5.1. MV-polynomials for Richardson varieties. Recall that according to the isomorphism of Section 4, if **p** is a Kostant picture consisting of a single loop L encircling dots $i, i+1, \ldots, j-1$ of the Dynkin diagram, then the corresponding MV-polynomial is just the matrix entry x_{ij} .

Then, by computing convolutions of MV-cycles, one may calculate MV-polynomials for more complicated Kostant pictures. For instance, from the computation of the example in 3.3, we see that

$$\mathcal{P}_{\bullet \bullet} = \mathcal{P}_{\bullet \bullet} \mathcal{P}_{\bullet \bullet} - \mathcal{P}_{\bullet \bullet} = x_{12}x_{23} - x_{13}$$

By doing such computations, one quickly finds that the MV-polynomial for any of the Richardson varieties is given by a minor. More precisely, suppose \mathbf{p} is a Kostant picture without any loops encircled by other loops. Let $\ell_1, \ell_2, \ldots, \ell_{|\mathbf{p}|}$ and $r_1, r_2, \ldots, r_{|\mathbf{p}|}$ denote the left and right ends of the loops of \mathbf{p} , ordered from left to right. Then $\mathcal{P}_{\mathbf{p}}$ is the minor $\det(x_{\ell_i, r_j})$. Here $(x_{ij}) \in \mathbb{N}$ denotes an arbitrary element of the group of upper triangular unipotent matrices, so that $x_{ii} = 1$ and

 $x_{ij} = 0$ if i > j. Moreover, all the nonvanishing minors arise in this way. For instance, in the above example, we have

$$\mathcal{P}_{\bullet\bullet} = \left| \begin{array}{cc} x_{12} & x_{13} \\ 1 & x_{23} \end{array} \right|$$

5.2. Which MV-polynomials are determinants? As one inductively computes more and more complicated MV-polynomials, one observes that often they are naturally expressible as determinants; indeed it is difficult to find one that is not. We begin by saying precisely what we mean by "naturally expressible".

Suppose $\mathcal{P}_{\mathbf{p}}$ is an MV-polynomial, with leading term $x^{\mathbf{p}}$. Label the loops of \mathbf{p} in any order: $L_1, L_2, \ldots, L_{|\mathbf{p}|}$. We will construct a $|\mathbf{p}|$ by $|\mathbf{p}|$ matrix $A_{\mathbf{p}}$ whose rows and columns are labelled by the loops $L_1, L_2, \ldots, L_{|\mathbf{p}|}$ and the product of whose diagonal entries is $x^{\mathbf{p}}$. We set the ij^{th} entry of $A_{\mathbf{p}}$ equal to x_{ℓ_i, r_j} where ℓ_i is the left end of L_i and r_j is the right end of L_j (recalling that $x_{ii} = 1$ and $x_{ij} = 0$ if i > j).

One observes that for many Kostant pictures \mathbf{p} the MV-polynomial $\mathcal{P}_{\mathbf{p}}$ is equal to the determinant of a matrix $\hat{A}_{\mathbf{p}}$, obtained from $A_{\mathbf{p}}$ by changing some of its entries to zeros. In particular, this is true for the Richardson varieties, in which case no entries are changed and $\hat{A}_{\mathbf{p}} = A_{\mathbf{p}}$, the determinant of which is the minor in the previous section.

Example. Let **p** be the Kostant picture



One can compute that the corresponding MV-polynomial is

$$\mathcal{P}_{\mathbf{p}} = x_{13}x_{24}x_{34}x_{15}x_{45}x_{46} - x_{13}x_{24}x_{34}x_{45}^2x_{16} - x_{13}x_{24}x_{15}x_{35}x_{46} \\ + x_{13}x_{24}x_{35}x_{45}x_{16} - x_{13}x_{34}x_{15}x_{45}x_{26} + x_{13}x_{34}x_{25}x_{45}x_{16} \\ + x_{13}x_{15}x_{35}x_{26} - x_{13}x_{25}x_{35}x_{16} - x_{23}x_{14}x_{34}x_{15}x_{45}x_{46} \\ + x_{23}x_{14}x_{34}x_{45}^2x_{16} + x_{23}x_{14}x_{15}x_{35}x_{46} - x_{23}x_{14}x_{35}x_{45}x_{16} \\ - x_{14}x_{15}x_{25}x_{46} + x_{14}x_{15}x_{45}x_{26} + x_{24}x_{15}^2x_{46} \\ - x_{24}x_{15}x_{45}x_{16} - x_{15}^2x_{26} + x_{15}x_{25}x_{16}$$

Let us order the six loops of \mathbf{p} according to the locations of their left ends in the drawing of the Kostant picture, working from left to right (so L_1 is the loop of length 4 and L_6 is the rightmost loop of length 1). Having chosen this order, we may write down the matrix $A_{\mathbf{p}}$, and (for instance by trial and error) we find that $\mathcal{P}_{\mathbf{p}} = \det(\hat{A}_{\mathbf{p}})$ after changing ten entries to zero.

$$A_{\mathbf{p}} = \begin{pmatrix} x_{15} & x_{13} & x_{14} & x_{14} & x_{16} & x_{15} \\ x_{15} & x_{13} & x_{14} & x_{14} & x_{16} & x_{15} \\ x_{25} & x_{23} & x_{24} & x_{24} & x_{26} & x_{25} \\ x_{35} & 1 & x_{34} & x_{34} & x_{36} & x_{35} \\ x_{45} & 0 & 1 & 1 & x_{46} & x_{45} \\ x_{45} & 0 & 1 & 1 & x_{46} & x_{45} \end{pmatrix}$$

$$\hat{A}_{\mathbf{p}} = \left(egin{array}{cccccccc} x_{15} & \mathbf{0} & \mathbf{0} & \mathbf{0} & x_{16} & \mathbf{0} \\ x_{15} & x_{13} & x_{14} & x_{14} & x_{16} & x_{15} \\ x_{25} & x_{23} & x_{24} & x_{24} & x_{26} & x_{25} \\ \mathbf{0} & 1 & \mathbf{0} & x_{34} & \mathbf{0} & x_{35} \\ x_{45} & 0 & 1 & 1 & x_{46} & x_{45} \\ \mathbf{0} & 0 & \mathbf{0} & 1 & \mathbf{0} & x_{45} \end{array}
ight)$$

Notice that there might be several different ways to set entries equal to zero so as to get the same determinant. For instance, in this example, subtracting the third column from the fourth column in $\hat{A}_{\mathbf{p}}$ would result in setting an additional three entries equal to zero without changing the determinant. Nevertheless, as we will discuss in the next section, we believe that in general there should be a canonical best choice of how to do this.

Counterexample. It is not easy to find an example of an MV-polynomial that cannot be expressed as a determinant in the way just described. The first counterexample is in type A_5 , and is actually the only such example we have been able to compute. The Kostant picture \mathbf{p}_1 is



and the MV-polynomial is

$$\mathcal{P}_{\mathbf{p}_{!}} = x_{34}x_{13}x_{46}x_{25} - x_{23}x_{34}x_{46}x_{15} - x_{34}x_{45}x_{13}x_{26} + x_{23}x_{34}x_{45}x_{16} \\ -x_{13}x_{25}x_{36} - x_{46}x_{14}x_{25} + x_{24}x_{46}x_{15} + x_{13}x_{35}x_{26} + x_{23}x_{36}x_{15} \\ +x_{45}x_{14}x_{26} - x_{23}x_{35}x_{16} - x_{45}x_{24}x_{16} - 2x_{15}x_{26} + 2x_{25}x_{16}$$

Then, ordering the loops left to right by their left ends,

$$A_{\mathbf{p}_!} = \begin{pmatrix} x_{13} & x_{15} & x_{14} & x_{16} \\ x_{23} & x_{25} & x_{24} & x_{26} \\ 1 & x_{35} & x_{34} & x_{36} \\ 0 & x_{45} & 1 & x_{46} \end{pmatrix}$$

and one can verify that $\mathcal{P}_{\mathbf{p}_!}$ is not equal to the determinant of any $\hat{A}_{\mathbf{p}_!}$ obtained by changing some entries to zeros.

We believe there to be a relation between the Fomin-Zelevinsky theory of cluster algebras [FZ1, FZ2] and the question of which MV-polynomials are expressible as determinants. The algebra $\mathbb{C}[N]$ carries the structure of a cluster algebra. One of the basic ingredients of this structure is a canonical set of generators for $\mathbb{C}[N]$, called *cluster variables*, and a superset of *cluster monomials*, which are certain monomials in the cluster variables. (Of course a cluster monomial is not necessarily a *monomial* in the variables \mathbf{x} ; it is a polynomial in these variables.) For completeness, more details are provided in Appendix A; in particular, we will give a precise definition of the cluster monomials in $\mathbb{C}[N]$, and describe the algorithm we used to compute them. In this section, we simply view cluster algebra theory as a black box that outputs certain polynomials—the cluster monomials—which we conjecture are a subset of MV-polynomials.

Conjecture 5.1. Every cluster monomial is an MV-polynomial that is naturally expressible as a determinant; that is, if ψ is a cluster monomial, then there is a Kostant picture \mathbf{p} such that $\psi = \mathcal{P}_{\mathbf{p}} = \det(\hat{A}_{\mathbf{p}})$, for some matrix $\hat{A}_{\mathbf{p}}$ obtained from $A_{\mathbf{p}}$ by setting some entries equal to zero.

Remark. If this conjecture is true then the above counterexample $\mathcal{P}_{\mathbf{p}_!}$ is an MV-polynomial which is not a cluster monomial.

This conjecture rests on three pieces of evidence:

- (1) For every cluster monomial that was small enough so as to be able to also compute the MV-polynomial with the same leading term, these two polynomials were identical.
- (2) We have been able to find the matrix $\hat{A}_{\mathbf{p}}$ for approximately a hundred different cluster monomials in type A_5 . A Kostant picture with nine or ten loops is about the maximum size for which it is usually possible to find $\hat{A}_{\mathbf{p}}$ within a few hours using educated guesswork (and letting a computer do the algebra of course). A real problem is that we do not know an efficient algorithm for finding $\hat{A}_{\mathbf{p}}$.
- (3) We strongly believe that the above counterexample of an MV-polynomial that is not expressible as a determinant is not a cluster monomial, because it has not appeared in the first 719 cluster variables for type A₅ output by the computer [AK*], whereas all other Kostant pictures of this size are accounted for. This same polynomial was used by Leclerc [Le] to give a counterexample to a conjecture of Berenstein-Zelevinsky. Leclerc showed that this MV-polynomial is an element of the dual canonical basis, but its square is not.

Remark. Although Conjecture 5.1 only refers to the cluster monomials, one might hope that a more general formula, of the form

$$\mathcal{P}_{\mathbf{p}} = \sum_{w \in \text{Perm}(\mathbf{p})} \text{sgn}(w) c_w x^w$$

holds for any MV-polynomial $\mathcal{P}_{\mathbf{p}}$. Here $\operatorname{Perm}(\mathbf{p})$ is the group of permutations of the loops of \mathbf{p} , $\operatorname{sgn}(w)$ is the sign of the permutation, and $x^w = \prod_{L \in \mathbf{p}} x_{\ell_w(L)}, r_L$ where w(L) has left end at $\ell_{w(L)}$ and L has right end at r_L ; as always $x_{ii} = 1$ and $x_{ij} = 0$ if i < j. (Note that different w can give the same x^w .) The coefficients c_w would be nonnegative integers somehow determined combinatorially from \mathbf{p} and the permutations w. (Unfortunately we are far from understanding the c_w —all we can really say is that c_w should be 0 if $x^w = x^{\mathbf{q}}$ with $\mathbf{q} > \mathbf{p}$.) In the case that $\mathcal{P}_{\mathbf{p}}$ is a cluster monomial, each c_w would be 0 or 1, and the sum would reduce to the definition of a determinant $\det(\hat{A}_{\mathbf{p}})$.

5.3. Which entries should be changed to 0? Here, we will give a conjecture as to how to find $\hat{A}_{\mathbf{p}}$ for a moderately large class of Kostant pictures \mathbf{p} for which we believe $\mathcal{P}_{\mathbf{p}}$ to be a determinant; that is, we will specify exactly which entries of $A_{\mathbf{p}}$ are to be changed to zeros.

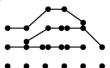
First, let us describe the class of Kostant pictures we will consider. Given any Kostant picture \mathbf{p} , we construct a directed graph $G_{\mathbf{p}}$ whose vertices are the loops of \mathbf{p} . It will be directed in the sense that for each edge, we will say that the loop at one end is on the left and the loop at the other end is on the right.

The first step in constructing $G_{\mathbf{p}}$ is to draw a "line diagram" of \mathbf{p} . This will be somewhat similar to the pictures back in Figure 1. For each loop L of \mathbf{p} and each column i through which it passes, let us say that its height h_i is equal to the length of the longest chain of loops of \mathbf{p} ending in L, say $L_1 \subset L_2 \subset \cdots \subset L_h = L$. Then, starting with a horizontal base of dots at height zero, we draw, for each loop in \mathbf{p} , a dot at height h_i for each column it passes through, and connect these dots by straight line segments.

For instance, for the Kostant picture in the first example of this section,



we draw



(Note that it is easy to draw these line diagrams inductively: (1) start by drawing the line segments for the loops that don't contain other loops; (2) for any loop L that encircles only loops for which the line segments have already been drawn, place the dot in each column at the lowest position that's one unit higher than any dot already drawn in that column for a loop that's encircled by L.)

Now, let us suppose L_1 and L_2 are overlapping loops of \mathbf{p} , meaning that neither loop encircles the other, and they pass through at least one column in common; without loss of generality suppose that their left and right ends satisfy $\ell_1 < \ell_2 \le r_1 < r_2$. We write $L_1 \to L_2$ provided that the line segments we drew for

 L_1 intersect those we drew for L_2 ; more precisely, we require that in column r_1 , the height of L_1 is less than or equal to the height of L_2 , and in column ℓ_2 , the height of L_2 is less than or equal to the height of L_1 . Then we connect L_1 and L_2 by an edge in $G_{\mathbf{p}}$, with L_1 on the left and L_2 on the right, provided that in addition there is no third loop L in \mathbf{p} with $L_1 \to L \to L_2$.

In the above example, the graph looks like this:



(The labelling of the loops is the same as it was before; working left to right, the loops of length 1 are L_4, L_6 , the loops of length 2 are L_2, L_3, L_5 , and the loop of length 4 is L_1 .) The picture is drawn such that for any edge the name of the loop at the left end is situated on the page to the left of the name of the loop at the right end.

Now we may say which Kostant pictures **p** we will consider: those for which the graph $G_{\mathbf{p}}$ is acyclic, i.e., is a disjoint union of trees. (Notice that the graph for the counterexample \mathbf{p}_{l} is not one of these; it is a 4-cycle.) For any such \mathbf{p} we will conjecture a matrix $\hat{A}_{\mathbf{p}}$ for which $\mathcal{P}_{\mathbf{p}} = \det(\hat{A}_{\mathbf{p}})$. So we must specify which entries of $A_{\mathbf{p}}$ to change to zeros. To do this, we need to consider paths in the graph $G_{\mathbf{p}}$.

Suppose that $(L_{a_1}, L_{a_2}, \dots, L_{a_k})$ is a path of distinct vertices in $G_{\mathbf{p}}$ (where we are allowed to travel in either direction along an edge). We say that it is an allowable path provided that for any three consecutive vertices $L_{a_i}, L_{a_{i+1}}, L_{a_{i+2}}$ on it:

- (1) If $L_{a_i} \to L_{a_{i+1}}$ and $L_{a_{i+2}} \to L_{a_{i+1}}$ then L_{a_i} encircles $L_{a_{i+2}}$. (2) If $L_{a_{i+1}} \to L_{a_i}$ and $L_{a_{i+1}} \to L_{a_{i+2}}$ then $L_{a_{i+2}}$ encircles L_{a_i} .

Conjecture 5.2. Suppose **p** is a Kostant picture for which the graph $G_{\mathbf{p}}$ is acyclic. Order the loops of \mathbf{p} and define the matrix $A_{\mathbf{p}}$ as before. Define $\hat{A}_{\mathbf{p}}$ to be the matrix obtained from $A_{\mathbf{p}}$ by changing the ij^{th} entry to 0 if there is no allowable path in $G_{\mathbf{p}}$ from the jth to the ith loop of \mathbf{p} . Then $\mathcal{P}_{\mathbf{p}} = \det(\hat{A}_{\mathbf{p}})$.

One can check in the above example that this results in changing to zeros the ten entries we saw earlier. For instance, the entry of $\hat{A}_{\mathbf{p}}$ in row 1 and column 2 is set to 0 since there is no allowable path from L_2 to L_1 ; indeed the path (L_2, L_3, L_5, L_1) is not allowable since condition (1) fails for the final three vertices $(L_3, L_5, L_1).$

Remarks.

• This conjecture is true for many other Kostant pictures as well, but for a different graph, which we don't know how to define in general. For instance, if **p** is the Kostant picture



then the conjecture does not apply since $G_{\mathbf{p}}$ is a 4-cycle. Nevertheless, we still have that $\mathcal{P}_{\mathbf{p}} = \det(\hat{A}_{\mathbf{p}})$, if instead of using $G_{\mathbf{p}}$ to construct $\hat{A}_{\mathbf{p}}$, we use a different graph: the one obtained from $G_{\mathbf{p}}$ by removing the edge joining the left loop of length 2 to the loop of length 4.

- We believe that the allowable path condition, although correct, is not the natural thing to say; that is, there should be a different, more general, condition which happens to give exactly this answer for this particular class of Kostant pictures.
- 5.4. Connecting MV-polynomials to geometry. We will state a conjecture that directly relates each MV-polynomial to the corresponding MV-cycle, in the case where it is a cluster monomial. Our motivation was to provide a non-inductive geometric definition of each MV-polynomial, but we still cannot do this.

When told this conjecture, a geometer typically expresses some combination of perplexity and dismay, perhaps asking "Where does this come from?" and "What does it mean?". The answers are: It is a purely empirical conjecture based on examples, and, unfortunately, we have no idea what it really means; we are not hiding any geometric intuition from the reader.

To state the conjecture, it is convenient to use the natural group action of $\mathsf{GL}_{\mathsf{n}}(\mathbb{C}[t,t^{-1}])$ on the affine Grassmanian $\mathbf{Gr}\colon\mathsf{GL}_{\mathsf{n}}$ acts by the standard representation on e_1,e_2,\ldots,e_n and t acts by sending each basis vector t^je_i to $t^{j+1}e_i$; extend this by linearity. Let N and N^- denote the subgroups of GL_{n} of upper and lower triangular unipotent matrices. Then it is easy to see that in the definition of MV -cycles as irreducible components of $\overline{S_\lambda \cap T_\mu}$ in Section 2.3, we have $S_\lambda = \mathsf{N}^-(\mathbb{C}[t,t^{-1}]) \cdot \underline{\lambda}$ and $T_\mu = \mathsf{N}(\mathbb{C}[t,t^{-1}]) \cdot \underline{\mu}$. So, given an MV -cycle M, with lowest coweight μ , there is a dense subset M° of points $Y \in M$ for which we can write $Y = \mathbf{n} \cdot \underline{\mu}$, where $\mathbf{n} \in \mathsf{N}(\mathbb{C}[t,t^{-1}])$; of course \mathbf{n} is not uniquely determined by Y.

Let us suppose $M = M(\tilde{\mathbf{p}})$ where $\tilde{\mathbf{p}} = (\mathbf{p}, \eta)$; as usual, λ and μ denote the highest and lowest coweights. We will define two functions $\chi_1, \chi_2 : M^{\circ} \to \mathbb{C}$ and conjecture that they are the same function. (Actually, the first function χ_1 will only be well-defined when the MV-polynomial $\mathcal{P}_{\mathbf{p}}$ is a cluster monomial, and even this well-definedness will be conjectural.) The definition of χ_1 will depend on $\mathcal{P}_{\mathbf{p}}$ but the definition of χ_2 will not.

5.4.1. The first function. Let Y be an arbitrary point of M° and write $Y = \mathbf{n} \cdot \underline{\mu}$. In short, we will define $\chi_1(Y)$ to be the lowest non-zero term of $\mathcal{P}_{\mathbf{p}}(\mathbf{n}) \in \mathbb{C}[t, t^{-1}]$, the evaluation of $\mathcal{P}_{\mathbf{p}}$ on $\mathbf{n} \in \mathsf{N}(\mathbb{C}[t, t^{-1}])$. Of course $\mathcal{P}_{\mathbf{p}}(\mathbf{n})$ depends not just on $Y \in M^{\circ}$, but on \mathbf{n} as well. Let d be the integer that is the net

difference in heights j of the basis vectors $t^{-j}e_i$ of $\underline{\lambda}$ and those of $\underline{\mu}$, that is, $d = (\sum_{t^{-j}e_i \in \underline{\mu} \text{ but } t^{-j}e_i \notin \underline{\lambda}} j) - (\sum_{t^{-j}e_i \in \underline{\lambda} \text{ but } t^{-j}e_i \notin \underline{\mu}} j)$. Then in the case where $\mathcal{P}_{\mathbf{p}}$ is a cluster monomial we conjecture that the coefficient of t^d in $\mathcal{P}_{\mathbf{p}}(\mathbf{n})$ does not depend on the choice of \mathbf{n} , and we denote this coefficient by $\chi_1(Y)$. It is important to assume that $\mathcal{P}_{\mathbf{p}}$ is a cluster monomial; in particular, for the counterexample Kostant picture $\mathbf{p}_!$, this coefficient does depend on the choice of \mathbf{n} . As for the other coefficients of $\mathcal{P}_{\mathbf{p}}(\mathbf{n})$, for k < d we conjecture that the coefficient of t^k is 0; and for k > d, any example will show that the coefficient of t^k depends on the choice of \mathbf{n} .

5.4.2. The second function. Here \mathbf{p} is allowed to be any Kostant picture, and we again let Y be an arbitrary element of M° . Let $\pi_{\lambda}: Y \to \underline{\lambda}$ and $\pi_{\mu}: Y \to \underline{\mu}$ denote orthogonal projection. It is not hard to see that $Y = \mathbf{n} \cdot \underline{\mu}$ implies that π_{μ} is invertible, so that we can define $\pi = \pi_{\lambda} \circ \pi_{\mu}^{-1} : \underline{\mu} \to \underline{\lambda}$. Although formally this is a linear map between infinite-dimensional vector spaces, it is really a map between finite-dimensional vector spaces of dimension $|\mathbf{p}|$: From the discussion in Section 2.5 of bases for lattices inside an MV-cycle, we can write $\underline{\lambda} = \underline{\eta} \oplus \underline{\hat{\lambda}}$ and $\underline{\mu} = \underline{\eta} \oplus \underline{\hat{\mu}}$, where $\underline{\hat{\lambda}}$ and $\underline{\hat{\mu}}$ are $|\mathbf{p}|$ -dimensional subspaces of $\underline{\lambda}$ and $\underline{\mu}$ orthogonal to $\underline{\eta}$, and π restricts to a map $\hat{\pi}: \underline{\hat{\mu}} \to \underline{\hat{\lambda}}$. We define $\chi_2(Y)$ to be the Jacobian of this linear map. Ordinarily, the Jacobian of a map between different vector spaces is only defined up to sign; but here the sign is determined since the Kostant picture \mathbf{p} gives a canonical identification between the basis vectors of $\underline{\hat{\lambda}}$ and of $\underline{\hat{\mu}}$: those for $\underline{\hat{\lambda}}$ correspond to the left ends of the loops and those for $\underline{\hat{\mu}}$ correspond to the right ends of the loops.

Conjecture 5.3. Suppose \mathbf{p} is a Kostant picture for which the MV-polynomial $\mathcal{P}_{\mathbf{p}}$ is a cluster monomial. Let $M = M(\tilde{\mathbf{p}})$ be the corresponding MV-cycle with lowest coweight μ , and let M° be its dense subset $M \cap (\mathbb{N}(\mathbb{C}[t, t^{-1}]) \cdot \underline{\mu})$. Then χ_1 is well-defined and the functions $\chi_1, \chi_2 : M^{\circ} \to \mathbb{C}$ are identical.

Remarks.

- It should be possible to extend $\chi_1 = \chi_2$ from M° to a function $\chi: M \to \mathbb{CP}^1$.
- There should be a direct relation between Conjecture 5.3 and Conjecture 5.2. In particular, one might be able to see from the definitions of χ_1 and χ_2 exactly which entries of the matrix $A_{\mathbf{p}}$ must be changed to zeros, because χ_1 is related to $\det(\hat{A}_{\mathbf{p}})$ and χ_2 is related (via lattices and linear algebra) to the combinatorics of \mathbf{p} . But we have not been able to see the connection, and the conjectures were arrived at independently.
- \bullet The evidence for Conjecture 5.3 is weaker than that for Conjecture 5.2; it has not been tested on Kostant pictures with more than four or five loops. \Box

Example. We will verify the conjecture when the Kostant picture \mathbf{p} is



In this case, the MV-polynomial is

$$\mathcal{P}_{\mathbf{p}} = \begin{vmatrix} x_{14} & \mathbf{0} & x_{15} \\ x_{24} & x_{23} & x_{25} \\ x_{34} & 1 & x_{35} \end{vmatrix} = x_{23}x_{35}x_{14} - x_{23}x_{34}x_{15} - x_{14}x_{25} + x_{24}x_{15}$$

Let us take $M=M(\mathbf{p},\eta_0)$ where $\underline{\eta_0}=X_0$ is spanned by all t^je_i for $j\geq 0$. Then $\mu=\eta_0\oplus \mathrm{span}\{t^{-1}e_3,t^{-1}e_4,t^{-1}e_5\}$. Write

$$\mathbf{n} = \begin{pmatrix} 1 & a & e & h & j \\ 0 & 1 & b & f & i \\ 0 & 0 & 1 & c & g \\ 0 & 0 & 0 & 1 & d \\ 0 & 0 & 0 & 0 & 1 \end{pmatrix}$$

for $\mathbf{n} \in \mathsf{N}(\mathbb{C}[t,t^{-1}])$, where $a = \sum_{k=-\infty}^{\infty} a_k t^k, b = \sum_{k=-\infty}^{\infty} b_k t^k, \dots \in \mathbb{C}[t,t^{-1}]$. Using the description in Section 2.3 of lattices in an MV-cycle, it is not hard to verify that necessary and sufficient conditions for $Y = \mathbf{n} \cdot \underline{\mu}$ to be in M° are: (1) e has degree ≥ 1 ; a,b,d,h,j have degrees ≥ 0 ; c,f,g,i have degrees ≥ -1 (where the degree is the smallest k for which the coefficient of t^k is nonzero). (2) The following minors vanish:

To compute $\chi_1(Y)$ we calculate

$$\mathcal{P}_{\mathbf{p}}(\mathbf{n}) = bgh - bcj - hi + fj$$

$$= \left(b_0 \begin{vmatrix} h_0 & j_0 \\ c_{-1} & g_{-1} \end{vmatrix} - \begin{vmatrix} h_0 & j_0 \\ f_{-1} & i_{-1} \end{vmatrix} \right) t^{-1}$$

$$+ \left(b_0 g_0 h_0 - b_0 c_0 j_0 - h_0 i_0 + f_0 j_0 + b_1 \begin{vmatrix} h_0 & j_0 \\ c_{-1} & g_{-1} \end{vmatrix} \right)$$

$$- j_1 \begin{vmatrix} b_0 & f_{-1} \\ 1 & c_{-1} \end{vmatrix} + h_1 \begin{vmatrix} b_0 & i_{-1} \\ 1 & g_{-1} \end{vmatrix}$$

$$+ terms of degree 1 and higher$$

where we have grouped together certain pairs of terms into vanishing minors; therefore

$$\chi_1(Y) = b_0 g_0 h_0 - b_0 c_0 j_0 - h_0 i_0 + f_0 j_0$$

Now we compute $\chi_2(Y)$. We have

$$\hat{\lambda} = \operatorname{span}\{t^{-1}e_1, t^{-1}e_2, t^{-1}e_3\}, \hat{\mu} = \operatorname{span}\{t^{-1}e_4, t^{-1}e_3, t^{-1}e_5\},$$

where we view these as ordered bases in the order written (so that the i^{th} vector corresponds to the i^{th} loop when the loops are ordered left to right by their left ends). With respect to these ordered bases, it is easy to check that we have

$$\hat{\pi} = \begin{pmatrix} h_0 & 0 & j_0 \\ f_0 & b_0 & i_0 \\ c_0 & 1 & g_0 \end{pmatrix} \begin{pmatrix} 1 & 0 & d_0 \\ c_0 & 1 & g_0 \\ 0 & 0 & 1 \end{pmatrix}^{-1}$$

The Jacobian of $\hat{\pi}$ is then the quotient of the determinants of these two matrices, which is

$$\chi_2(Y) = b_0 g_0 h_0 - b_0 c_0 j_0 - h_0 i_0 + f_0 j_0$$

and we have verified that $\chi_1(Y) = \chi_2(Y)$ for any $Y \in M^{\circ}$.

APPENDIX A. CLUSTER ALGEBRA

Here we define a very interesting set of generators for $\mathbb{C}[N]$ giving it the structure of a cluster algebra. Cluster algebras were discovered by Fomin and Zelevinsky [FZ1, FZ2, BFZ], and have proved to be related to many different topics. By now the theory of cluster algebras is so large, that it would be pointless to even try to summarize their motivation and many applications here. All we will do is recall the definitions and results needed to define a cluster algebra inside $\mathbb{C}[N]$ for $G = GL_n$ or SL_n . For the more general treatment of these results see [BFZ].

For us a cluster \mathbf{c} will be a collection $(c_1,\ldots,c_{\frac{n(n-1)}{2}},B)$ of $\frac{n(n-1)}{2}$ elements c_k of $\mathbb{C}[\mathbb{N}]$ and an $\frac{n(n-1)}{2}\times\frac{(n-1)(n-2)}{2}$ matrix of integers B. The clusters will be defined inductively: we will define an original cluster \mathbf{c}_0 and specify a mutation rule designed to produce more clusters from existing clusters. More specifically, for a cluster \mathbf{c} and an element c_k of this cluster with $k\leq \frac{(n-1)(n-2)}{2}$ we will define another cluster $\mu_k(\mathbf{c})=(c_1,\ldots,c_k',\ldots,c_{\frac{n(n-1)}{2}},\mu_k(B))$ by replacing c_k by a new variable c_k' and replacing c_k by another matrix c_k' and replacing c_k' by another matrix c_k' and c_k' and c_k' by another matrix c_k' and c_k'

The original cluster is defined as follows. Since we can identify N with the space of upper triangular unipotent matrices, every minor of such an $n \times n$ matrix can be though of as an element of $\mathbb{C}[N]$. For $1 \leq a < b \leq n$ define $\Delta_{ab} \in \mathbb{C}[N]$ to be the minor with columns $a, \ldots, b-1$ and the top b-a rows of the $n \times n$ matrix. The functions of the original cluster are given by Δ_{ab} . More specifically, let $k(a,b) = \frac{b(b-1)}{2} - a$, and set $c_{k(a,b)} = \Delta_{ab}$. The matrix B is given as follows. The entry of B in the intersection of row k(a,b) and column k(a',b') is

- 1, if a' = a 1, b' = b 1, or a' = a + 1, b' = b, or a' = a, b' = b + 1;
- -1, if a' = a + 1, b' = b + 1, or a' = a 1, b' = b, or a' = a, b' = b 1;
- 0, otherwise.

Let $B = (b_{ij})$. The mutation rule is given by:

$$c'_{k} = \frac{\prod_{b_{ik} > 0} c_{i}^{b_{ik}} + \prod_{b_{ik} < 0} c_{i}^{-b_{ik}}}{c_{l}}$$

and if $\mu_k(B) = (b'_{ij})$ by

$$b'_{ij} = \begin{cases} -b_{ij} & \text{if } i = k \text{ or } j = k; \\ b_{ij} + \frac{|b_{ik}|b_{kj} + b_{ik}|b_{kj}|}{2} & \text{otherwise.} \end{cases}$$

It is not at all obvious from the above definition that mutation is well-defined, namely why each c'_k is in $\mathbb{C}[N]$. However Berenstein-Fomin-Zelevinsky [BFZ] showed that this is true. In particular, any sequence of mutations produces a cluster. Let us identify two clusters if one can be produced from the other by simultaneous reindexing of the elements c_k and the rows and columns of the matrix B. It is then natural to ask whether the number of clusters is finite or infinite. It turns out that it is finite for $n \leq 5$, but there are infinitely many clusters for $n \geq 6$.

The elements of the clusters are called *cluster variables*, while products of the form $\prod_{c_k \in \mathbf{c}} c_k^{a_k}$ are called *cluster monomials*. These generate the algebra $\mathbb{C}[\mathsf{N}]$. For $n \leq 5$, the cluster monomials span $\mathbb{C}[\mathsf{N}]$ as a vector space, but for $n \geq 6$ they are expected to span a proper subspace.

We have programmed this inductive definition into the computer, and used it to compute all cluster variables in types A_2 , A_3 , A_4 (where there are 4, 12, 40 of them, respectively) and a list of 719 cluster variables in type A_5 [AK*].

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