FORM–TYPE CALABI–YAU EQUATIONS

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Abstract. Motivated from mathematical aspects of the superstring theory, we introduce a new equation on a balanced, hermitian manifold, with zero first Chern class. By solving the equation, one will obtain, in each Bott–Chern cohomology class, a balanced metric which is hermitian Ricci–flat. This can be viewed as a differential form level generalization of the classical Calabi–Yau equation. We establish the existence and uniqueness of the equation on complex tori, and prove certain uniqueness and openness on a general Kähler manifold.

1. Setting and Equations

In the superstring theory, the internal space $X³$ is a complex three-dimensional manifold with a non-vanishing holomorphic three-form Ω [15] (cf. [1]). The $N = 1$ supersymmetry requires [15, 10]

$$
d(\parallel\Omega\parallel_{\omega}\omega^2)=0,
$$

for some hermitian metric (form) ω . The above equation in mathematics says that ω is a conformally balanced metric. (We recall that [14] a hermitian metric ω on an n-dimensional complex manifold X^n is called *balanced* if ω satisfies that

$$
d(\omega^{n-1}) = 0 \qquad \text{on } X^n.)
$$

Note that $[8, 6]$ the torus bundles over K3 surfaces and over complex abelian surfaces twisted by two anti-self dual $(1, 1)$ -forms admit a non-vanishing holomorphic threeform Ω and a natural balanced metric ω_0 such that

$$
(1.1) \t\t\t\t \t\t\t \t\t\t \t\t\t \t\t\t \t\t\t \mathbb{R} \t\t\t \t\t\t \t\t\t \mathbb{R} \t\t\t \math
$$

As important examples in the superstring theory and non-Kähler complex geometry, the complex manifolds $\#_k(S^3 \times S^3)$ for any $k \geq 2$ [4, 12] also admit a nonvanishing holomorphic three-form [4] and a balanced metric [5]. Moreover, we know that $\#_k(S^3 \times S^3)$ satisfies the $\partial\bar{\partial}$ -lemma [4]. A natural question to ask is, whether $\#_k(S^3 \times S^3)$ admits a balanced metric ω_0 such that (1.1) holds. Such a metric ω_0 , if exists, will play an important role in the superstring theory and hermitian geometry.

More generally, let X^n $(n \geq 3)$ be a complex *n*-dimensional manifold with a nonvanishing holomorphic n-form Ω and with a balanced metric ω_0 . We want to look for a balanced metric ω such that √

(1.2)
$$
\omega^{n-1} = \omega_0^{n-1} + \frac{\sqrt{-1}}{2} \partial \bar{\partial} \varphi,
$$

for some real $(n-2, n-2)$ –form φ , and such that

(1.3) $\|\Omega\|_{\omega} =$ some positive constant C_0 .

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In other words, we would like to find solutions of (1.3) in the cohomology class $[\omega_0^{n-1}] \in H_{BC}^{n-1,n-1}(X)$. Here $H_{BC}^{p,q}(X)$ stands for the Bott–Chern cohomology:

$$
H^{p,q}_{BC}(X)=\frac{(\ker\partial\cap\ker\bar{\partial})\cap\Omega^{p,q}(X)}{\operatorname{im}\, \partial\bar{\partial}\cap\Omega^{p,q}(X)}.
$$

One can certainly normalize the constant C_0 in (1.3) to be 1, as in (1.1) . However, it may be more convenient to set

$$
C_0 = \left(\int_X \omega^n\right)^{-\frac{1}{2}},
$$

from the equation point of view. As in the Kähler case, equation (1.3) is equivalent to the equation

(1.4)
$$
\frac{\det \omega}{\det \omega_0} = \frac{\| \Omega \|_{\omega_0}^2}{\| \Omega \|_{\omega}^2} = e^f \frac{\int_X \omega^n}{\int_X \omega_0^n}.
$$

Here we denote

$$
e^f = \|\Omega\|_{\omega_0}^2 \int_X \omega_0^n,
$$

and denote

$$
\det \omega = \det(g_{i\bar{j}}), \quad \text{if} \quad \omega = \frac{\sqrt{-1}}{2} \sum_{i,j=1}^n g_{i\bar{j}} dz_i \wedge d\bar{z}_j.
$$

At the moment, we write

$$
\omega_0^{n-1} + \frac{\sqrt{-1}}{2} \partial \overline{\partial} \varphi = \left(\frac{\sqrt{-1}}{2}\right)^{n-1} (n-1)!
$$

$$
\cdot \sum_{i,j=1}^n (\Psi_{\varphi})_{i\overline{j}} s(i,j) dz_1 \wedge d\overline{z}_1 \wedge \cdots \wedge \widehat{dz}_i \wedge \cdots \widehat{dz}_j \wedge \cdots \wedge dz_n \wedge d\overline{z}_n.
$$

Here the sign function $s(i, j)$ is equal to 1 if $i \leq j$, and is equal to -1 if $i > j$. By (1.2) and

$$
\omega^{n-1} = \left(\frac{\sqrt{-1}}{2}\right)^{n-1} (n-1)!
$$

$$
\cdot (\det \omega) \sum_{i,j=1}^{n} g^{i\bar{j}} s(i,j) dz_1 \wedge d\bar{z}_1 \wedge \cdots \wedge \widehat{dz}_i \wedge \cdots \widehat{dz}_j \wedge \cdots \wedge dz_n \wedge d\bar{z}_n,
$$

we have

$$
(\det \omega)g^{i\bar{j}} = (\Psi_{\varphi})_{i\bar{j}}, \quad \text{for all } 1 \le i, j \le n.
$$

Hence,

$$
\det \omega = \left\{ \det \left[(\Psi_{\varphi})_{i\bar{j}} \right] \right\}^{\frac{1}{n-1}} = \left\{ \det \left[\omega_0^{n-1} + (\sqrt{-1}/2) \partial \bar{\partial} \varphi \right] \right\}^{\frac{1}{n-1}}.
$$

Here det $\left[\omega_0^{n-1} + (\sqrt{-1}/2)\partial \bar{\partial}\varphi\right]$ stands for the determinant of $n \times n$ matrix of its coefficients. Thus, equation (1.4) is equivalent to

(1.5)
$$
\frac{\det[\omega_0^{n-1} + (\sqrt{-1}/2)\partial\bar{\partial}\varphi]}{\det \omega_0^{n-1}} = e^{(n-1)f} \left(\frac{\int_X \omega^n}{\int_X \omega_0^n}\right)^{n-1}.
$$

We call the above equation the form-type Calabi–Yau equation. Clearly, by integrating (1.4), we obtain a compatibility condition

(1.6)
$$
\int_X e^f \omega_0^n = \int_X \omega_0^n.
$$

Let us denote by $\mathcal{P}(\omega_0)$ the set of all smooth real $(n-2, n-2)$ –forms ψ such that

(1.7)
$$
\omega_0^{n-1} + \frac{\sqrt{-1}}{2} \partial \bar{\partial} \psi > 0 \quad \text{on } X.
$$

The question is therefore reduced to find, for a given $f \in C^{\infty}(X)$ with (1.6), a smooth real $(n-2, n-2)$ –form $\varphi \in \mathcal{P}(\omega_0)$ satisfying (1.5).

Here is the geometric interpretation of our equation. Let us briefly recall some definitions related to the hermitian connection. We follow $[9]$. Let R be the curvature of hermitian connection with respect to metric ω . Then,

$$
R_{i\bar{j}k\bar{l}} = -\frac{\partial^2 g_{i\bar{j}}}{\partial z_k \partial \bar{z}_l} + \sum_{p,q=1}^n g^{p\bar{q}} \frac{\partial g_{i\bar{q}}}{\partial z_k} \frac{\partial g_{p\bar{j}}}{\partial \bar{z}_l}.
$$

We set

$$
R_{k\bar{l}} = \sum_{i,j=1}^{n} g^{i\bar{j}} R_{i\bar{j}k\bar{l}}
$$

,

and associate with it a real $(1, 1)$ -form given by

$$
Ric^{h} = \sqrt{-1} \sum_{k,l=1}^{n} R_{k\bar{l}} dz_{k} \wedge d\bar{z}_{l}.
$$

We call Ric^h the $Ricci curvature$ of the hermitian connection. Clearly,

$$
Ric^{h} = \sqrt{-1}\bar{\partial}\partial \log(\det \omega).
$$

So $\| \Omega \|_{\omega} = C_0$ is equivalent to the Ricci curvature $Ric^h = 0$.

On the other hand, we can also define the Ricci form Ric^s of the spin connection (i.e. Bismut connection) on a hermitian manifold. The relation between the two Ricci forms is given by [11]

$$
Ric^s = Ric^h + dd^*\omega.
$$

Here d^* is the adjoint operator of d with respect to the metric ω . So when ω is balanced, $Ric^s = Ric^h$, and hence, $|| \Omega ||_{\omega} = C_0$ is also equivalent to the Ricci curvature of the spin connection is zero.

In particular, if ω_0 is Kähler and let φ to be

either
$$
u\sum_{i=0}^{n-2} {n-1 \choose i} \left(\frac{\sqrt{-1}}{2} \partial \overline{\partial} u\right)^{n-i-2} \wedge \omega_0^i, \text{ or}
$$

$$
-\frac{\sqrt{-1}}{2} \partial u \wedge \overline{\partial} u \wedge \sum_{i=0}^{n-3} {n-1 \choose i} \left(\frac{\sqrt{-1}}{2} \partial \overline{\partial} u\right)^{n-i-3} \wedge \omega_0^i,
$$

then (1.5) is reduced to

$$
\frac{\det(\omega_0 + \frac{\sqrt{-1}}{2}\partial\bar{\partial}u)}{\det \omega_0} = e^f.
$$

This is the classic equation in the Calabi Conjecture on $c_1(X) = 0$, which was settled by Yau [16].

It seems to us that a form-type equation such as (1.5) has not yet been studied. To begin with, we consider the form-type Calabi–Yau equation on $Tⁿ$, the complex *n*-torus. Let (z_1, \ldots, z_n) be the complex coordinates on T^n induced from \mathbb{C}^n . Then, any non-vanishing holomorphic n-form Ω on T^n is equal to

$$
dz_1 \wedge \cdots \wedge dz_n
$$

up to multiplying a nonzero constant. We fix such an n-form Ω . By a constant form or a constant metric on T^n we mean a differential form or a metric on T^n with constant coefficients. Let ω_0 be a balanced metric on T^n . As far as the Bott–Chern cohomology class of ω_0^{n-1} is concerned, we can assume, without loss of generality, that ω_0 is a constant metric on T^n . This is due to the fact that any closed differential form on T^n is cohomologous to a constant form, and the $\partial\bar{\partial}$ -Lemma. Our result is as follows:

Theorem 1. Let Ω be a non-vanishing holomorphic n-form on T^n $(n \geq 3)$, and ω_0 be a constant metric on T^n such that $\|\Omega\|_{\omega_0} = 1$. We denote by C_0 a positive constant.

(1) If $C_0 \leq 1$, then for any metric ω on T^n such that $[\omega^{n-1}] = [\omega_0^{n-1}] \in$ $H_{BC}^{n-1,n-1}(T^n)$ and that $\|\Omega\|_{\omega} = C_0$, we must have $C_0 = 1$ and

 $\omega = \omega_0$.

(2) For each $C_0 > 1$, there exists a metric ω on T^n such that $[\omega^{n-1}] = [\omega_0^{n-1}]$ and that

$$
\|\Omega\|_{\omega} = C_0.
$$

One can see from Theorem 1 that the normalization constant C_0 plays a role here. When $C_0 \leq 1$, the theorem tells us that the Calabi–Yau metric is the unique canonical balanced metric. It is the second case, $C_0 > 1$, that marks the difference between a form-type equation and a usual function-type equation. In this case, we establish the existence of a desired balanced metric which is not Calabi–Yau. We further generalize the uniqueness part, Theorem 1 (1), to an arbitrary Calabi–Yau manifold:

Theorem 2. Let X be a compact Kähler manifold with a non-vanishing holomorphic n-form Ω . Let ω_0 be a Calabi–Yau metric such that $\|\Omega\|_{\omega_0} = 1$. Then, for any balanced metric ω on X such that ω^{n-1} represents the Bott–Chern cohomology class of ω_0^{n-1} and such that $\|\Omega\|_{\omega} = C_0 \le 1$, we have

$$
\omega = \omega_0.
$$

For a general case that ω_0 is non-Kähler, one can use the continuity method to solve (1.5) . As an initial step we consider the openness. Here we have to assume X to be a Kähler manifold, endowed with a Kähler metric η . For nonnegative integers k and m, and a real number $0 < \alpha < 1$, we denote by $C^{k,\alpha}(\Lambda^{m,m}(X))$ the Hölder space of real (m, m) –forms on X, and in particular, $C^{k,\alpha}(\Lambda^{0,0}(X)) \equiv C^{k,\alpha}(X)$. Let

$$
\mathcal{F}^{k,\alpha}(X) = \left\{ g \in C^{k,\alpha}(X); \int_X e^g \,\omega_0^n = \int_X \omega_0^n \right\}.
$$

Then $\mathcal{F}^{k,\alpha}(X)$ is a hypersurface in the Banach space $C^{k,\alpha}(X)$. Let ω_0 be a Hermitian metric on X, and $\mathcal{P}(\omega_0)$ be the set given by (1.7). We define a map $M : \mathcal{P}(\omega_0) \cap$ $C^{k+2,\alpha}(\Lambda^{n-2,n-2}(X)) \to \mathcal{F}^{k,\alpha}(X)$ by

$$
M(\psi) = \log\left(\frac{\omega_{\psi}^{n}}{\omega_{0}^{n}}\right) - \log\left(\frac{\int_{X} \omega_{\psi}^{n}}{\int_{X} \omega_{0}^{n}}\right),\,
$$

where by abuse of notation, $\mathcal{P}(\omega_0) \cap C^{k+2,\alpha}(\Lambda^{n-2,n-2}(X))$ stands for

$$
\{\psi \in C^{k+2,\alpha}(\Lambda^{n-2,n-2}(X)); \omega_0^{n-1} + (\sqrt{-1}/2)\partial\bar{\partial}\psi > 0\},\
$$

and for each $\psi \in \mathcal{P}(\omega_0) \cap C^{k+2,\alpha}(\Lambda^{n-2,n-2}(X))$, we denote by ω_{ψ} the positive $(1,1)$ form on X such that

$$
\omega_{\psi}^{n-1} = \omega_0^{n-1} + (\sqrt{-1}/2)\partial\bar{\partial}\psi.
$$

Note that equation (1.5) can be written as

$$
M(\varphi) = f.
$$

Theorem 3. Let X be an n-dimensional Kähler manifold $(n \geq 3)$, ω_0 be a Hermitian metric on X, $k \geq n+4$ be an integer, and $0 < \alpha < 1$ be a real number. Given $f \in \mathcal{F}^{k,\alpha}(X)$, suppose that $\varphi \in \mathcal{P}(\omega_0) \cap C^{k+2,\alpha}(\Lambda^{n-2,n-2}(X))$ satisfies

 $M(\varphi) = f$.

Then, there is a positive number δ , such that for any $g \in \mathcal{F}^{k,\alpha}(X)$ with $||g-f||_{C^{k,\alpha}(X)} \leq$ δ, there exists a function $ψ ∈ P(ω₀) ∩ C^{k+2,α}(Λ^{n-2,n-2}(X))$ such that

$$
M(\psi) = g.
$$

The rest of the paper is organized as follows: In Section 2, we first show Theorem 1 (1). Next, we prove Theorem 1 (2) by explicitly constructing a smooth solution $\varphi \in \mathcal{P}(\omega_0)$ for the form-type equation. These arguments make use of special properties such as the flat structure of T^n . We prove Theorem 2 at the end of Section 2. In this respect, we essentially present two proofs for the uniqueness on $Tⁿ$, as they may have interests of their own. In Section 3, we prove Theorem 3 in full details, where one can see the compatibility condition is crucial. Moreover, the approach differs from the standard one in that, the special $(n-2, n-2)$ –forms $(u\eta^{n-2})$ are taken, and also in the argument of Proposition 15 and Proposition 16.

2. Uniqueness and Existence

In this section, we adopt the following index convention, unless otherwise indicated. For an $(n-1, n-1)$ –form Θ , we denote

$$
\Theta = \left(\frac{\sqrt{-1}}{2}\right)^{n-1} (n-1)!
$$

$$
\sum_{p,q} s(p,q) \Theta_{pq} dz^1 \wedge d\bar{z}^1 \cdots \wedge \widehat{dz}^p \wedge d\bar{z}^p \wedge \cdots \wedge d\bar{z}^q \wedge \widehat{d\bar{z}}^q \wedge \cdots \wedge dz^n \wedge d\bar{z}^n,
$$

in which

(2.1)
$$
s(p,q) = \begin{cases} -1, & \text{if } p > q; \\ 1, & \text{if } p \leq q. \end{cases}
$$

Here we introduce the sign function s so that,

$$
dz^p \wedge d\bar{z}^q \wedge s(p,q)dz^1 \wedge d\bar{z}^1 \cdots \wedge \widehat{dz}^p \wedge d\bar{z}^p \wedge \cdots \wedge d\bar{z}^q \wedge \widehat{dz}^q \wedge \cdots \wedge dz^n \wedge d\bar{z}^n
$$

= $dz^1 \wedge d\bar{z}^1 \wedge \cdots \wedge dz^n \wedge d\bar{z}^n$, for all $1 \leq p, q \leq n$.

And, if the matrix $(\Theta_{p\bar{q}})$ is invertible, we denote by $(\Theta^{p\bar{q}})$ the transposed inverse of $(\Theta_{p\bar{q}})$, i.e.,

$$
\sum_l \Theta_{i\bar{l}} \Theta^{j\bar{l}} = \delta_{ij}.
$$

In the following, we may also use the summation convention on repeating indices.

2.1. Torus case. Throughout this subsection, we consider $X = Tⁿ$, the complex n-torus with $n \geq 3$. We shall prove Theorem 1. Note that the first part of Theorem 1 follows immediately from Lemma 4 below. We shall prove the second part in Lemma 7.

Lemma 4. Let ω_0 be a constant metric on T^n . Suppose that there exists an $(n 2, n-2$)–form $\varphi \in \mathcal{P}(\omega_0)$ and a constant $0 < C_0 \leq 1$ such that

(2.2)
$$
C_0 \det \left(\omega_0^{n-1} + \frac{\sqrt{-1}}{2} \partial \bar{\partial} \varphi \right) = \det \omega_0^{n-1}.
$$

Then, we have $C_0 = 1$ and

$$
\sqrt{-1}\partial\bar{\partial}\varphi=0.
$$

We need two propositions to derive Lemma 4. Let (z_1, \ldots, z_n) be the complex coordinates on T^n induced from \mathbb{C}^n . The corresponding real coordinates are (x_1, \ldots, x_{2n}) . Here we denote

(2.3)
$$
z_i = x_{2i-1} + \sqrt{-1}x_{2i}, \text{ and hence, } \frac{\partial}{\partial z_i} = \frac{1}{2} \left(\frac{\partial}{\partial x_{2i-1}} - \sqrt{-1} \frac{\partial}{\partial x_{2i}} \right),
$$

for all $1 \leq i \leq n$. We choose the following volume form on T^n :

$$
dV = (\sqrt{-1}/2)^n dz_1 \wedge d\bar{z}_1 \wedge \cdots \wedge dz_n \wedge d\bar{z}_n.
$$

Here are two elementary facts:

Proposition 5. For any smooth complex function f defined on T^n , we have

$$
\int_{T^n} \frac{\partial^2 f}{\partial z_i \partial z_j} dV = 0, \quad \text{for all } i, j = 1 \cdots, n.
$$

Proof. We write

$$
f = f_1 + \sqrt{-1}f_2,
$$

where f_1, f_2 are real functions on T^n . Then,

$$
4\frac{\partial^2 f_1}{\partial z_i \partial z_j} = \frac{\partial^2 f_1}{\partial x_{2i-1} \partial x_{2j-1}} - \frac{\partial^2 f_1}{\partial x_{2i} \partial x_{2j}} - \sqrt{-1} \left(\frac{\partial^2 f_1}{\partial x_{2i} \partial x_{2j-1}} + \frac{\partial^2 f_1}{\partial x_{2i-1} \partial x_{2j}} \right).
$$

We have a similar equation for f_2 . And note that

$$
dV = dx_1 \wedge \cdots \wedge dx_{2n}.
$$

The result then obviously follows from the fundamental theorem of calculus. \Box

Proposition 6. Let $B = (b_{i\bar{j}})$ be a hermitian matrix on T^n , in which each entry $b_{i\bar{j}}$ is a complex smooth function defined on T^n such that

$$
\int_{T^n} b_{i\overline{j}} \, dV = 0.
$$

Assume that $I + B$ is everywhere positive definite, and there is a constant $c \geq 1$ such that

$$
\det(I + B) = c \quad \text{on } T^n, \quad \text{where } I \equiv (\delta_{i\bar{j}}).
$$

Then, $c = 1$ and $B = 0$.

Proof. Since $I + B$ is positive definite, we have

(2.4)
$$
\frac{\operatorname{tr}(I+B)}{n} \geq \sqrt[n]{\det(I+B)} = \sqrt[n]{c} \quad \text{on } T^n.
$$

Integrating (2.4) over $Tⁿ$, we obtain

$$
\int_{T^n} dV = \int_{T^n} \frac{\text{tr}(I + B)}{n} dV \ge \sqrt[n]{c} \int_{T^n} dV.
$$

Thus, $c = 1$, and the inequality of (2.4) is in fact an equality. That is,

(2.5)
$$
\frac{\text{tr}(I + B)}{n} = \sqrt[n]{\text{det}(I + B)} = 1, \quad \text{on } T^n.
$$

Now at an arbitrary point x in $Tⁿ$, we choose a unitary matrix U such that

$$
UB\bar{U}^T = \text{dial}\{\lambda_1, \cdots, \lambda_n\}.
$$

Then (2.5) is equivalent to that

$$
1 + \lambda_1 = 1 + \lambda_2 = \dots = 1 + \lambda_n = 1.
$$

This implies that

$$
\lambda_i = 0, \qquad \text{for all } i = 1, \dots, n.
$$

Therefore, $B = 0$ at x. Since x is arbitrary, this finishes the proof.

Let us now proceed to prove Lemma 4:

Proof of Lemma 4. Let
\n
$$
\omega_0^{n-1} = \left(\frac{\sqrt{-1}}{2}\right)^{n-1} (n-1)!
$$
\n
$$
\sum_{p,q} \Psi_{p\bar{q}} s(p,q) dz^1 \wedge d\bar{z}^1 \wedge \cdots \wedge \widehat{dz}^p \wedge \cdots \wedge \cdots \wedge \widehat{d\bar{z}^q} \wedge \cdots \wedge dz^n \wedge d\bar{z}^n.
$$

Here $(\Psi_{i\bar{j}})$ is a constant, positive definite, hermitian matrix, and $s(p,q)$ is given by (2.1). We can then take a non-degenerate constant matrix A such that

$$
(2.6) \t\t A(\Psi_{i\bar{j}})\bar{A}^T = I.
$$

We define a hermitian matrix $F_{\varphi} = ((F_{\varphi})_{i\bar{j}})$ on T^n by

$$
\frac{\sqrt{-1}}{2} \partial \bar{\partial} \varphi = \left(\frac{\sqrt{-1}}{2}\right)^{n-1} (n-1)!
$$

$$
\sum_{p,q} (F_{\varphi})_{p\bar{q}} s(p,q) dz^1 \wedge d\bar{z}^1 \wedge \cdots \wedge \widehat{dz}^p \wedge \cdots \wedge \cdots \wedge \widehat{dz}^q \wedge \cdots \wedge dz^n \wedge d\bar{z}^n.
$$

It follows from Proposition 5 that

$$
\int_{T^n} (F_{\varphi})_{i\overline{j}} dV = 0.
$$

Then, by (2.2) and (2.6) ,

$$
\det(I + AF_{\varphi}\bar{A}^T) = C_0^{-1}.
$$

Since $\varphi \in \mathcal{P}(\omega_0)$, we obtain

$$
I + AF_{\varphi} \bar{A}^T > 0 \qquad \text{on } T^n.
$$

Applying Proposition 6 yields that $C_0 = 1$, and

$$
AF_{\varphi}\bar{A}^T = 0
$$

and therefore,

$$
F_{\varphi}=0.
$$

 \Box

The following lemma establishes the second part of Theorem 1. By a linear transformation, if necessary, we can assume the constant metric ω_0 on T^n to be the standard metric:

$$
\omega_0 = \frac{\sqrt{-1}}{2} \big(dz_1 \wedge d\overline{z}_1 + \cdots + dz_n \wedge d\overline{z}_n \big).
$$

Lemma 7. For any $0 < \delta < 1$, there exists a smooth $(n-2, n-2)$ –form $\varphi \in \mathcal{P}(\omega_0)$ such that √

(2.7)
$$
\det \left(\omega_0^{n-1} + \frac{\sqrt{-1}}{2} \partial \bar{\partial} \varphi \right) = \delta \det \omega_0^{n-1}.
$$

Proof of Lemma 7. We set

(2.8)
$$
\varphi = (n-1)! \left(\frac{\sqrt{-1}}{2} \right)^{n-2} \left[u(z_1, \bar{z}_1) dz_3 \wedge d\bar{z}_3 \wedge \cdots \wedge dz_n \wedge d\bar{z}_n + v(z_1, \bar{z}_1) dz_2 \wedge d\bar{z}_2 \wedge \widehat{dz}_3 \wedge d\bar{z}_3 \wedge dz_4 \wedge d\bar{z}_4 \wedge \cdots \wedge dz_n \wedge d\bar{z}_n \right].
$$

Here u, v are two real, smooth, periodic functions to be determined, with $1 + \Delta u > 0$ and $1 + \Delta v > 0$. Since u and v depend only on the first variable, the equation (2.7) becomes that

(2.9)
$$
\left(1+\frac{\partial^2 u}{\partial z_1 \partial \bar{z}_1}\right)\left(1+\frac{\partial^2 v}{\partial z_1 \partial \bar{z}_1}\right)=\delta.
$$

This reduces to an equation on T^1 . Note that

$$
\frac{\partial^2 u}{\partial z_1 \partial \bar{z}_1} = \Delta u, \quad \frac{\partial^2 v}{\partial z_1 \partial \bar{z}_1} = \Delta v,
$$

where Δ is the standard Laplacian on T^1 , i.e., the Laplacian associated with $\omega_0|_{T^1}$. We can rewrite (2.9) as

(2.10)
$$
1 + \Delta u = \frac{\delta}{1 + \Delta v}.
$$

Our strategy is to fix a function v and then solve (2.10) for a function u . Note that for a fixed v , the necessary and sufficient condition to solve (2.10) is that

(2.11)
$$
\int_{T^1} \omega_0|_{T^1} = \delta \int_{T^1} \frac{\omega_0|_{T^1}}{1 + \Delta v}.
$$

Now let

(2.12)
$$
v = -4k \sin\left(\frac{z_1 + \bar{z}_1}{2}\right) = -4k \sin x_1,
$$

where $0 \leq k \leq 1$ is a constant to be determined, and the change of coordinates is given by (2.3) . Then, (2.11) becomes that

$$
\int_{T^1} dx_1 \wedge dx_2 = \int_{T^1} \frac{\delta}{1 + k \sin x_1} dx_1 \wedge dx_2,
$$

that is,

(2.13)
$$
\int_0^{2\pi} \frac{\delta}{1 + k \sin x_1} dx_1 = 2\pi.
$$

It follows from the proposition below that, for each $0 < \delta < 1$, there exists a real number $0 < k < 1$, depending only on δ , such that (2.13) holds. Therefore, for v given by (2.12) , there is a smooth function u , unique up to a constant, satisfies (2.10) . Also, by the construction,

$$
1 + \Delta v > 0, \qquad 1 + \Delta u > 0.
$$

Thus, by (2.8) we obtain an $(n-2, n-2)$ –form $\varphi \in \mathcal{P}(\omega_0)$ which solves (2.7). \square

Proposition 8. Let

(2.14)
$$
Z(k) = \frac{1}{2\pi} \int_0^{2\pi} \frac{1}{1 + k \sin x} dx, \quad \text{for all } 0 \le k < 1.
$$

Then, for any $0 < \delta < 1$, there exists a unique number $0 < k_{\delta} < 1$ such that

$$
Z(k_{\delta}) = \delta^{-1}.
$$

Proof. Clearly, the function Z is smooth on $0 \le k < 1$. Note that $Z(0) = 1$, and that

$$
Z(k) \ge \frac{1}{2\pi} \int_{3\pi/2}^{2\pi} \frac{dx}{1 + k \sin x}
$$

=
$$
\frac{1}{\pi\sqrt{1 - k^2}} \arctan \sqrt{\frac{1 + k}{1 - k}} \to +\infty, \text{ as } k \to 1^-.
$$

The existence then follows from the intermediate value theorem in calculus. The uniqueness is due to the monotonicity of Z on $[0, 1)$, which is readily seen by verifying $Z'(0) = 0$ and $Z''(k) > 0$ on [0, 1].

Lemma 7 can be easily generalized to the case of the product of a compact hermitian manifold with T^k , $k \geq 3$. See the corollary below:

Corollary 9. Let $M^n = N^{n-k} \times T^k$, $k \geq 3$, where (N^{n-k}, ω_N) is an $(n-k)$ dimensional compact hermitian manifold. We denote $\omega = \omega_N + \omega_0$, where ω_0 is a constant metric on T^k . Then, for any $1 > \delta > 0$, there exists a smooth $(n-2, n-2)$ form $\psi \in \mathcal{P}(\omega)$ on M such that

$$
\det\left(\omega^{n-1} + \frac{\sqrt{-1}}{2}\partial\bar{\partial}\psi\right) = \delta \det(\omega^{n-1}).
$$

Proof. Let

$$
\psi = \omega_N^{n-k} \wedge \varphi,
$$

where φ is the $(k-2, k-2)$ –form on T^k obtained by Lemma 7, i.e., $\varphi \in \mathcal{P}(\omega_0)$ satisfies that such that √

$$
\det \left(\omega_0^{k-1} + \frac{\sqrt{-1}}{2} \partial \bar{\partial} \varphi \right) = \delta \det(\omega_0^{k-1}).
$$

Then, obviously ψ satisfies the requirement.

2.2. Kähler case. In this subsection, we shall prove Theorem 2. Observe that it is sufficient to prove the following lemma.

Lemma 10. Let (X, ω_0) be a compact Kähler manifold. Consider

$$
\det\left(\omega_0^{n-1} + \frac{\sqrt{-1}}{2}\partial\bar{\partial}\varphi\right) = C_1 \det \omega_0^{n-1},
$$

where $\varphi \in \mathcal{P}(\omega_0)$, and $C_1 > 0$ is a constant. If $C_1 \geq 1$, then

$$
\sqrt{-1}\partial\bar{\partial}\varphi=0.
$$

Proof. By a direct calculation, since ω_0 is Kähler, we have

$$
\int_X (\omega_0^{n-1})^{i\bar{j}} \left(\frac{\sqrt{-1}}{2} \partial\bar{\partial}\varphi\right)_{i\bar{j}} \omega_0^n = n \int_X \omega_0 \wedge \left(\frac{\sqrt{-1}}{2} \partial\bar{\partial}\varphi\right) = 0.
$$

Similar to the torus case, we apply the arithmetic–geometric mean inequality to obtain

(2.15)
\n
$$
C_1^{1/n} = \left[\frac{\det(\omega_{\varphi}^{n-1})}{\det(\omega_0^{n-1})} \right]^{1/n}
$$
\n
$$
\leq 1 + (\omega_0^{n-1})^{i\overline{j}} \left(\frac{\sqrt{-1}}{2} \partial \overline{\partial} \varphi \right)
$$

Integrating over X with respect to ω_0 and using first equality yields that

$$
C_1^{1/n} \int_X \omega_0^n \le \int_X \omega_0^n.
$$

This shows that $C_1 = 1$ and we must have a pointwise equality in (2.15). This forces that √

$$
\frac{\sqrt{-1}}{2}\partial\bar{\partial}\varphi=0.
$$

 $i\bar{j}$.

3. Openness

Let (X, η) be a Kähler manifold, and ω_0 be a Hermitian metric on X. Given $f \in C^{\infty}(X)$, we would like to study the solution $\varphi \in \mathcal{P}(\omega_0)$ of the following equation

(3.1)
$$
\frac{\omega_{\varphi}^n}{\omega_0^n} = \frac{e^f}{V} \int_X \omega_{\varphi}^n.
$$

Here ω_{φ} is a positive $(1, 1)$ –form on X such that

$$
\omega_{\varphi}^{n-1} = \omega_0^{n-1} + (\sqrt{-1}/2)\partial\bar{\partial}\varphi,
$$

and

$$
V = \int_X \omega_0^n.
$$

Equation (3.1) is the same as (1.4) , which is equivalent to the form-type Calabi–Yau equation (1.5). A compatibility condition for (3.1) is

$$
\int_X e^f \omega_0^n = V.
$$

In what follows, we fix k to be an integer greater than $n+3$, and fix a real number α with $0 < \alpha < 1$. We denote by $C^{k,\alpha}(X)$ the usual Hölder space of real-valued functions on X . Recall that

$$
\mathcal{F}^{k,\alpha}(X) = \left\{ g \in C^{k,\alpha}(X); \int_X e^g \,\omega_0^n = V \right\},\,
$$

which is a hypersurface in the Banach space $C^{k,\alpha}(X)$. For any ψ contained in the intersection of $\mathcal{P}(\omega_0)$ and $C^{k+2,\alpha}(\Lambda^{n-2,n-2}(X)),$

$$
M(\psi) \equiv \log \frac{\omega_{\psi}^n}{\omega_0^n} - \log \left(\frac{1}{V} \int_X \omega_{\psi}^n \right) \in \mathcal{F}^{k, \alpha}(X).
$$

By the map M , equation (3.1) can be rewritten as

$$
M(\varphi) = f.
$$

To prove Theorem 3, we first compute the linearization of M.

Proposition 11. Let $G(\varphi) = \omega_{\varphi}^n$ for all $\varphi \in \mathcal{P}(\omega_0)$, and denote by G_{φ} the Fréchet derivative of G at φ . Then, given $\varphi \in \mathcal{P}(\omega_0)$, we have

$$
G_{\varphi}(\psi) = \frac{n\sqrt{-1}}{2(n-1)} \partial \bar{\partial} \psi \wedge \omega_{\varphi},
$$

for all $\psi \in C^{k+2,\alpha}(\Lambda^{n-2,n-2}(X)).$

Proof. For any real $(n-2, n-2)$ –form ψ ,

(3.2)
\n
$$
G_{\varphi}(\psi) = \frac{d}{ds} \left(\omega_{\varphi + s\psi}^{n} \right) \Big|_{s=0}
$$
\n
$$
= n\omega_{\varphi}^{n-1} \wedge \frac{d}{ds} \left(\omega_{\varphi + s\psi} \right) \Big|_{s=0}
$$
\n
$$
= \frac{d}{ds} \left(\omega_{\varphi + s\psi}^{n-1} \right) \Big|_{s=0} \wedge \omega_{\varphi} + \omega_{\varphi}^{n-1} \wedge \frac{d}{ds} \left(\omega_{\varphi + s\psi} \right) \Big|_{s=0}
$$
\n(3.3)
\n
$$
= (\sqrt{-1}/2) \partial \bar{\partial} \psi \wedge \omega_{\varphi} + \omega_{\varphi}^{n-1} \wedge \frac{d}{ds} \left(\omega_{\varphi + s\psi} \right) \Big|_{s=0}.
$$

Comparing (3.2) with (3.3), we obtain that

$$
G_{\varphi}(\psi) = \frac{n}{n-1} (\sqrt{-1}/2) \partial \bar{\partial} \psi \wedge \omega_{\varphi}.
$$

 \Box

Corollary 12. For any $\varphi \in \mathcal{P}(\omega_0)$, the Fréchet derivative of M at φ is given by

$$
M_{\varphi}(\psi) = \frac{n(\sqrt{-1}/2)\partial\bar{\partial}\psi \wedge \omega_{\varphi}}{(n-1)\omega_{\varphi}^{n}} - \frac{n\int_{X}(\sqrt{-1}/2)\partial\bar{\partial}\psi \wedge \omega_{\varphi}}{(n-1)\int_{X}\omega_{\varphi}^{n}},
$$

for all $\psi \in C^{k+2,\alpha}(\Lambda^{n-2,n-2}(X)).$

Next, we recall the Local Surjectivity Theorem (see [13, p. 175 and p. 108], for example).

Theorem (Local Surjectivity Theorem). Let $\mathcal E$ and $\mathcal F$ be Banach manifolds, and $U \subset \mathcal{E}$ be an open subset. If $\mathfrak{F}: U \to \mathcal{F}$ is a C^1 map, and $\mathfrak{F}_\xi \equiv D\mathfrak{F}(\xi)$ is onto from the tangent space $T_{\xi} \mathcal{E}$ to the tangent space $T_{\mathfrak{F}(\xi)} \mathcal{F}$, then \mathfrak{F} is locally onto; that is, there exist open neighborhoods U_1 of ξ and V_1 of $\mathfrak{F}(\xi)$ such that $\mathfrak{F}|_{U_1}: U_1 \to V_1$ is onto.

Thus, to show Theorem 3, it suffices to show that the linearization M_{φ} is surjective from $C^{k+2,\alpha}(\Lambda^{n-2,n-2}(X))$ to $T_f\mathcal{F}^{k,\alpha}(X)$, which denotes the tangent space of $\mathcal{F}^{k,\alpha}(X)$ at f. Now let us introduce the space

$$
\mathcal{E}^{k,\alpha}(X) = \left\{ h \in C^{k,\alpha}(X); \int_X h \,\omega^n_{\varphi} = 0 \right\}.
$$

Note that $\mathcal{E}^{k,\alpha}(X)$ is itself a Banach space, as a closed subspace in $C^{k,\alpha}(X)$. There is another point of view: We can define an equivalence relation on the elements in $C^{k,\alpha}(X)$ by

 $h \sim g$ if and only if $h - g \equiv$ some constant.

In this regard, $\mathcal{E}^{k,\alpha}(X) = C^{k,\alpha}(X) / \sim$. Observe that

$$
T_f \mathcal{F}^{k,\alpha}(X) = \mathcal{E}^{k,\alpha}(X).
$$

To prove the surjectivity of M_{φ} , we consider a special class of the $(n-2, n-2)$ – forms, that is,

(3.4)
$$
\psi = u\eta^{n-2}, \qquad \text{where } u \in \mathcal{E}^{k+2,\alpha}(X).
$$

We recall that η is the Kähler metric on X. For simplicity we denote

$$
L(u) = M_{\varphi}(u\eta^{n-2}).
$$

Then, by Corollary 12,

(3.5)
$$
Lu = \frac{n(\sqrt{-1}/2)\partial\bar{\partial}u \wedge \eta^{n-2} \wedge \omega_{\varphi}}{(n-1)\omega_{\varphi}^n} - \frac{n\int_X (\sqrt{-1}/2)\partial\bar{\partial}u \wedge \eta^{n-2} \wedge \omega_{\varphi}}{(n-1)\int_X \omega_{\varphi}^n}.
$$

We shall prove the following result:

Lemma 13. Let $k \geq n+4$, and $0 < \alpha < 1$. For any $h \in \mathcal{E}^{k,\alpha}(X)$, there exists a unique function $u \in \mathcal{E}^{k+2,\alpha}(X)$ satisfying that

$$
(3.6) \t\t\t Lu = h.
$$

Lemma 13 implies that $M_{\varphi}: C^{k+2,\alpha}(\Lambda^{n-2,n-2}(X)) \to \mathcal{E}^{k,\alpha}(X)$ is surjective, and hence, Theorem 3 follows.

The rest of this section is devoted to prove Lemma 13. We denote by $W^{k,p}(\Omega, \omega_\varphi)$ the usual Sobolev space with respect to ω_{φ} on a domain Ω in X. In the rest of this section, we may denote $W^{k,p}(\Omega) = W^{k,p}(\Omega, \omega_\varphi)$ for simplicity; furthermore, when $\Omega = X$, we abbreviate $W^{k,p} = W^{k,p}(X) = W^{k,p}(X, \omega_{\varphi})$. Notice that $W^{0,2}(X) \equiv$ $L^2(X)$.

We introduce the following spaces:

$$
\mathcal{H} = \left\{ v \in W^{1,2}(X); \int_X v \, \omega_{\varphi}^n = 0 \right\},\,
$$

and

$$
\mathcal{L} = \left\{ v \in L^2(X); \int_X v \, \omega_{\varphi}^n = 0 \right\}.
$$

Clearly, H and L are Hilbert spaces, as closed subspaces in $W^{1,2}(X)$ and $L^2(X)$, respectively. We define a bilinear map $A: \mathcal{H} \times \mathcal{H} \to \mathbb{R}$ by

$$
A(u, v) = \frac{n\sqrt{-1}}{4(n-1)} \int_X \eta^{n-2} \wedge \omega_{\varphi} \wedge (\partial u \wedge \bar{\partial} v + \partial v \wedge \bar{\partial} u)
$$

+
$$
\frac{n\sqrt{-1}}{4(n-1)} \int_X v\eta^{n-2} \wedge (\partial u \wedge \bar{\partial} \omega_{\varphi} + \partial \omega_{\varphi} \wedge \bar{\partial} u)).
$$

Definition 14. Given $h \in \mathcal{L}$, we say that $u \in \mathcal{H}$ is a *weak* solution of the equation

$$
(3.7) \t\t -Lu = h,
$$

if u satisfies that

(3.8)
$$
A(u, v) = \int_X hv \,\omega_{\varphi}^n \equiv \langle h, v \rangle_{L^2}, \quad \text{for all } v \in \mathcal{H}.
$$

Let us remark that, if u is a *classical* solution of (3.7) , i.e., $u \in C^2(X)$, then one can obtain (3.8) by integrating (3.7) by parts with respect to ω_{φ}^n . Conversely, we have the following result:

Proposition 15. If $u \in C^3(X)$ satisfies (3.8) for some $h \in C^1(X) \cap \mathcal{L}$, then

$$
-Lu = h.
$$

Proof. First, we claim the following fact: If $\chi \in C^1(X)$ satisfy that

(3.9)
$$
\int_X \chi v \,\omega_{\varphi}^n = 0, \quad \text{for all } v \in \mathcal{H},
$$

then χ is a constant function on X. To see this, let

$$
v = \chi - \frac{\int_X \chi \omega_{\varphi}^n}{\int_X \omega_{\varphi}^n};
$$

then $v \in \mathcal{H}$ and (3.9) implies that

$$
\int_X |v|^2 \omega_\varphi^n = 0.
$$

This proves the claim. It follows that

$$
\frac{n(\sqrt{-1}/2)\partial\bar{\partial}u\wedge\eta^{n-2}\wedge\omega_{\varphi}}{(n-1)\omega_{\varphi}^{n}}-h=\text{some constant}.
$$

Thus, integrating with respect to ω_{φ}^n yields the result.

The following weak maximum principle is similar to that on a domain in the Euclidean space (see, for example, Gilbarg–Trudinger [7, p. 179]). Proposition 16 is trivial, if $d\omega_{\varphi} = 0$.

Proposition 16. Suppose that $u \in \mathcal{H}$ satisfies

(3.10)
$$
A(u, v) = 0, \quad \text{for all } v \in \mathcal{H}.
$$

Then, $u = 0$.

Proof. It suffices to prove sup_X $u \leq 0$, as one can then replace u by $-u$. (Here sup stands for the essential supremum.) Suppose the contrary. Take a constant δ such that $0 < \delta < \sup_X u$, and define

(3.11)
$$
v = (u - \delta)^{+} - \frac{\int_{X} (u - \delta)^{+} \omega_{\varphi}^{n}}{\int_{X} \omega_{\varphi}^{n}},
$$

in which $(u - \delta)^+ = \max\{u - \delta, 0\}$. Then, $v \in \mathcal{H}$ and

$$
dv = d(u - \delta)^{+} = \begin{cases} du, & \text{if } u > \delta, \\ 0, & \text{if } u \leq \delta. \end{cases}
$$

Let us denote by Γ the compact support of dv. Then, we obtain by (3.10) and metric equivalence of η , ω_{φ} , that

$$
\|\nabla v\|_{L^2}^2 = \int_{\Gamma} |\nabla v|^2 \omega_{\varphi}^n \le C \int_{\Gamma} |v| |\nabla v| \omega_{\varphi}^n.
$$

Here and below, we denote by C a generic positive constant depending only on η , ω_{φ} , and n . Apply Hölder's inequality to get

(3.12)
$$
\|\nabla v\|_{L^2} \leq C \|v\|_{L^2(\Gamma)}.
$$

On the other hand, combining the Sobolev inequality and Poincaré inequality yields that

$$
(3.13) \t\t\t ||v||_{L^{2n/(n-1)}} \leq C(||\nabla v||_{L^2} + ||v||_{L^2}) \leq C||\nabla v||_{L^2}.
$$

Hence, by (3.12) and (3.13),

$$
||v||_{L^{2n/(n-1)}} \leq C||v||_{L^{2}(\Gamma)} \leq C|\Gamma|^{\frac{1}{2n}}||v||_{L^{2n/(n-1)}},
$$

in which $|\Gamma|$ denotes the measure of Γ with respect to ω_{φ} . It follows that

$$
|\Gamma|=|\{u>\delta, |du|>0\}|\geq C^{-1}.
$$

Letting δ tend to sup u implies that $|du| > 0$ on a set of positive measure in $\{x \in$ $X; u(x) = \sup u$, which is evidently impossible by Lemma 7.7 in Gilbarg–Trudinger $[7, p. 152].$

The next two propositions are standard, for which we need the Lax–Milgram Theorem (see Evans [2, p. 297], for example) and the Fredholm alternative (see [2, p. 641] for example). We include them here for completeness.

Theorem (Lax–Milgram Theorem). Let H be a real Hilbert space, and $I : H \times$ $H \to \mathbb{R}$ be a bilinear mapping. Assume that, there exist positive constants β and μ such that

 $|I(u, v)| \leq \beta ||u|| ||v||$, for all $u, v \in H$,

and

$$
I(v, v) \ge \mu ||v||^2, \quad \text{for all } v \in H.
$$

Then, for any bounded linear functional f on H, there exists a unique element $u \in H$ satisfying that

$$
I(u, v) = f(v) \quad \text{for all } v \in H.
$$

Theorem (Fredholm alternative). Let E be a Banach space and $K : E \to E$ be a compact linear operator. Then,

$$
\ker(I - K) = \{0\} \quad \text{ if and only if } \quad \text{Im}(I - K) = E,
$$

where $I: E \to E$ is the identity operator.

Proposition 17. There exists a nonnegative constant γ , depending on ω_{φ} and η , such that for any $h \in \mathcal{L}$, there exists a unique weak solution $u \in \mathcal{H}$ of

(3.14)
$$
-L_{\gamma}u \equiv -Lu + \gamma u = h.
$$

That is, the function u satisfies

(3.15)
$$
A(u, v) + \gamma \langle u, v \rangle_{L^2} = \langle h, v \rangle_{L^2}, \quad \text{for all } v \in \mathcal{H}.
$$

Proof. We have, by the metric equivalence of η and ω_{φ} ,

$$
|A(u, v)| \leq \beta ||u||_{W^{1,2}} ||v||_{W^{1,2}},
$$

and

$$
A(u, u) + \gamma ||u||_{L^2} \ge \mu ||u||_{W^{1,2}}.
$$

Here $\beta > 0$, $\gamma \ge 0$, and $\mu > 0$ are constants depending only on η and ω_{φ} . The result then follows from applying Lax–Milgram Theorem to

$$
I(u, v) = A(u, v) + \gamma \langle u, v \rangle_{L^2}, \quad \text{for all } u, v \in \mathcal{H}.
$$

 \Box

Proposition 18. For any $h \in \mathcal{L}$, there exists a unique weak solution $u \in \mathcal{H}$ of

 $-Lu = h$.

Proof. By Proposition 17 we can define a map L_{γ}^{-1} : $\mathcal{L} \to \mathcal{H}$ as follows: For each $f \in \mathcal{L}$, we define $L_{\gamma}^{-1}(f)$ to be the unique function $w \in \mathcal{H}$ satisfying

$$
A(w, v) + \gamma \langle w, v \rangle_{L^2} = \langle f, v \rangle_{L^2}.
$$

Clearly, L_{γ}^{-1} is linear, and is a compact operator from $\mathcal L$ to $\mathcal L$, in view of Rellich Theorem. To prove the result, it suffices to show that, for a given $h \in \mathcal{L}$, there exists a unique $u \in \mathcal{L}$ satisfying that

$$
u = L_{\gamma}^{-1}(h + \gamma u).
$$

Equivalently, we need to solve a unique $u \in \mathcal{L}$ for the following equation:

$$
(I - \gamma L_{\gamma}^{-1})u = L_{\gamma}^{-1}h.
$$

To invoke the Fredholm alternative, we turn to the kernel of $(I - \gamma L_{\gamma}^{-1})$ in \mathcal{L} , i.e.,

$$
\{u \in \mathcal{L}; u - \gamma L_{\gamma}^{-1}u = 0\}.
$$

This is equivalent to investigate the function $u \in \mathcal{H}$ such that

$$
A(u, v) = 0 \qquad \text{for all } v \in \mathcal{H}.
$$

By Proposition 16, $u = 0$. The result then follows from the Fredholm alternative. \square

Now we are in a position to prove Lemma 13:

Proof of Lemma 13. The uniqueness of (3.6) is an immediate consequence of Proposition 18, since a C^2 solution of (3.6) is in particular a weak solution of $-Lu = -h$.

Given $h \in C^{k,\alpha}(X)$, we have $h \in W^{k,2}(X)$, since X is compact. Then, by Proposition 18, equation (3.6) has a weak solution $u \in W^{1,2}(X)$. Then, we obtain

$$
u \in W^{k+2,2}(X),
$$

by the local regularity theorem (see, for example, Evans [2, p. 314] or Gilbarg– Trudinger [7, p. 186]). Since $k \geq n+4$, $k - 2n/2 - 1 \geq 3$. We apply the Sobolev imbedding theorem to obtain that

$$
u \in C^3(X).
$$

By Proposition 15, u is the classical solution for (3.6) . It follows from the bootstrap argument $([7, p. 109])$ that

$$
u \in C^{k+2,\alpha}(X).
$$

 \Box

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