STABILITY UNDER DEFORMATIONS OF EXTREMAL ALMOST-KÄHLER METRICS IN DIMENSION 4.

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ABSTRACT. Given a path of almost-Kähler metrics compatible with a fixed symplectic form on a compact 4-manifold such that at time zero the almost-Kähler metric is an extremal Kähler one, we prove, for a short time and under a certain hypothesis, the existence of a smooth family of extremal almost-Kähler metrics compatible with the same symplectic form, such that at each time the induced almost-complex structure is diffeomorphic to the one induced by the path.

1. Introduction

An almost-Kähler metric on a 2n-dimensional symplectic manifold (M, ω) is induced by an almost-complex structure J compatible with ω in the sense that the tensor field $g(\cdot, \cdot) = \omega(\cdot, J \cdot)$ is symmetric and positive definite and thus it defines a Riemannian metric on M. The almost-Kähler metric is Kähler if the almost-complex structure J is integrable. Given an almost-Kähler metric, one can define a canonical hermitian connection ∇ (see e.g. [16, 24]). The hermitian scalar curvature s^{∇} is then obtained by taking a trace and contracting the curvature of ∇ with ω . In the Kähler case, the hermitian scalar curvature coincides with the Riemannian scalar curvature.

A key observation, made by Fujiki [13] in the integrable case and by Donaldson [9] in the general almost-Kähler case, asserts that the natural action of the infinite dimensional group $Ham(M,\omega)$ of hamiltonian symplectomorphisms on the space AK_{ω} of ω -compatible almost-Kähler metrics is hamiltonian with moment map $\mu: AK_{\omega} \to (Lie(Ham(M,\omega)))^*$ given by $\mu_J(f) = \int_M s^{\nabla} f \frac{\omega^n}{n!}$. The critical points of the norm $\int_M \left(s^{\nabla}\right)^2 \frac{\omega^n}{n!}$ are called extremal almost-Kähler metrics. It turns out that the symplectic gradient of s^{∇} of such metrics is a holomorphic vector field in the sense that its flow preserves the corresponding almost-complex structure. In particular, extremal Kähler metrics in the sense of Calabi [7] and almost-Kähler metrics with constant hermitian scalar curvature are extremal.

The GIT formal picture in [9] suggests the existence and the uniqueness of an extremal almost-Kähler metric, modulo the action of $Ham(M,\omega)$, in each 'stable complexified' orbit of the action of $Ham(M,\omega)$. However, in this formal infinite dimensional setting, a natural complexification of $Ham(M,\omega)$ does not exist. When $H^1(M,\mathbb{R}) = 0$, an identification of the 'complexified' orbit of a Kähler metric $(J,g) \in AK_{\omega}$ is given by considering all Kähler metrics (J,\tilde{g}) in the Kähler class $[\omega]$ and applying Moser's Lemma [9]. In this setting, Fujiki–Schumacher [14] and LeBrun–Simanca [21] showed, in the abscence of holomorphic vector fileds, that the existence of

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an extremal Kähler metric is an open condition on the space of such orbits. Moreover, Apostolov–Calderbank–Gauduchon–Friedman [3] generalized this result by fixing a maximal torus T in the reduced automorphism group of (M,J) and considering T-invariant ω -compatible Kähler metrics. In general, for an almost-Kähler metric, a description of these 'complexified' orbits is not avaible, see however [10] for the toric case. Nevertheless, the formal picture suggests that the existence of an extremal Kähler metric should persist for smooth almost-Kähler metrics close to an extremal one.

Thus motivated, we consider in this paper the 4-dimensional case where one can introduce a notion of almost-Kähler potential related to the one defined by Weinkove [27, 28]. In the spirit of [14, 21], we shall apply the Banach Implicit Function Theorem for the hermitian scalar curvature of T-invariant ω -compatible almost-Kähler metrics where T is a maximal torus in $Ham(M,\omega)$. The main technical problem is the regularity of a family of Green operators involved in the definition of the almost-Kähler potential. Using a Kodaira–Spencer result [19, 20], one can resolve this problem if we suppose that the dimension of g_t -harmonic J_t -anti-invariant 2-forms, denoted by $h_{J_t}^-$ (see [12]), satisfies the condition $h_{J_t}^- = h_{J_0}^- = b^+(M) - 1$ for $t \in (-\epsilon, \epsilon)$ along the path $(J_t, g_t) \in AK_{\omega}^T$ in the space of T-invariant ω -compatible almost-Kähler metrics. So, our main theorem claims the following

Theorem 1.1. Let (M, ω) be a 4-dimensional compact symplectic manifold and T a maximal torus in $Ham(M, \omega)$. Let (J_t, g_t) be any smooth family of almost-Kähler metrics in AK_{ω}^T such that (J_0, g_0) is an extremal Kähler metric. Suppose that $h_{J_t}^- = h_{J_0}^- = b^+(M) - 1$ for $t \in (-\epsilon, \epsilon)$. Then, there exists a smooth family $(\tilde{J}_t, \tilde{g}_t)$ of extremal almost-Kähler metrics in AK_{ω}^T , defined for sufficiently small t, with $(\tilde{J}_0, \tilde{g}_0) = (J_0, g_0)$ and such that \tilde{J}_t is equivariantly diffeomorphic to J_t .

Remark 1.2. (i) The condition that $h_{J_t}^- = h_{J_0}^- = b^+(M) - 1$ for $t \in (-\epsilon, \epsilon)$ is satisified in the following cases:

- (1) When J_t are integrable almost-complex structures for each t. Then, $h_{J_t}^- = 2h^{2,0}(M,J_t) = b^+(M) 1$ by a well-known result of Kodaira [5]. On the other hand, it is unknown whether or not, for an ω -compatible non-integrable almost-complex J on a compact 4-dimensional symplectic manifold M with $b^+(M) \geq 3$, the equality $h_J^- = b^+(M) 1$ is possible (see [12]).
- (2) When $b^+(M) = 1$, $h_{J_t}^- = 0$ for each t. This condition is satisfied when (M, ω) admits a non trivial torus in $Ham(M, \omega)$ [17].
- (ii) Theorem 1.1 holds under the weaker assumption that the torus $T \subset Ham(M, \omega)$ is maximal in $Ham(M, \omega) \cap Isom_0(M, g_0)$, where $Isom_0(M, g_0)$ denotes the connected component of the isometry group of the initial metric g_0 . By a known result of Calabi [8], any extremal Kähler metric is invariant under a maximal connected compact subgroup of $Ham(M, \omega) \cap \widetilde{Aut}(M, J_0)$, where $\widetilde{Aut}(M, J_0)$ is the reduced automorphism group of (M, J_0) . Hence, Theorem 1.1 generalizes [14, 21] in the 4-dimensional case. (iii) It was kindly pointed out to us by T. Drăghici that using a recent result of Donaldson and Remarks (i) and (ii) above, one can further extend Theorem 1.1 in the case when $b^+(M) = 1$ as follows: Let (M, ω_0, J_0, g_0) be a compact 4-dimensional extremal Kähler manifold with $b^+(M) = 1$ and T be a maximal torus in $Ham(M, \omega) \cap$

 $Isom_0(M, g_0)$. Then, for any smooth family of T-invariant almost-complex structures J(t) with $J(t) = J_0$, J(t) is compatible with an extremal almost-Kähler metric g_t for $t \in (-\epsilon, \epsilon)$. Indeed, as J(t) are tamed by ω_0 for $t \in (-\epsilon, \epsilon)$ and $b^+(M) = 1$, one can use the openess result of Donaldson [11, Proposition 1] (see also [12, Sec. 5]) to show that there exists a smooth family of J(t)-invariant symplectic forms ω_t with $[\omega_t] = [\omega_0]$. Averaging ω_t over the compact group T and using the equivariant Moser Lemma, we obtain a family J_t of T-invariant ω_0 -compatible almost-complex structures such that J_t is T-equivariantly diffeomorphic to J(t). We can then apply Theorem 1.1 to produce compatible extremal metrics

Kim and Sung [18] showed that, in any dimension, if one starts with a Kähler metric of constant scalar curvature with no holomorphic vector fields, one can construct infinite dimensional families of almost-Kähler metrics of constant hermitian scalar curvature which concide with the initial metric away from an open set. Similar existence result was presented in [22] when the initial Kähler metric is locally toric.

2. Preliminaries

Let (M,ω) be a compact symplectic manifold of dimension 2n. An almost-complex structure J is compatible with ω if the tensor field $g(\cdot,\cdot):=\omega(\cdot,J\cdot)$ defines a Riemannian metric on M; then, (J,g) is called an $(\omega$ -compatible) almost-Kähler metric on (M,ω) . If, additionally, the almost-complex structure J is integrable, then (J,g) is a $K\ddot{a}hler$ metric on (M,ω) .

The almost-complex structure J acts on the cotangent bundle $T^*(M)$ by $J\alpha(X) = -\alpha(JX)$, where α is a 1-form and X a vector field on M. Any section ψ of the bundle $\otimes^2 T^*(M)$ admits an orthogonal splitting $\psi = \psi^{J,+} + \psi^{J,-}$, where $\psi^{J,+}$ is the J-invariant part and $\psi^{J,-}$ is the J-anti-invariant part, given by

$$\psi^{J,+}(\cdot,\cdot) = \frac{1}{2} \left(\psi(\cdot,\cdot) + \psi(J\cdot,J\cdot) \right) \text{ and } \psi^{J,-}(\cdot,\cdot) = \frac{1}{2} \left(\psi(\cdot,\cdot) - \psi(J\cdot,J\cdot) \right).$$

In particular, the bundle of 2-forms decomposes under the action of J

(2.1)
$$\Lambda^{2}(M) = \mathbb{R} \cdot \omega \oplus \Lambda_{0}^{J,+}(M) \oplus \Lambda^{J,-}(M),$$

where $\Lambda_0^{J,+}(M)$ is the subbundle of the *primitive J*-invariant 2-forms (i.e. 2-forms pointwise orthogonal to ω) and $\Lambda^{J,-}(M)$ is the subbundle of *J*-anti-invariant 2-forms. Hence, the subbundle of primitive 2-forms $\Lambda_0^2(M)$ admits the splitting

$$\Lambda_0^2(M) = \Lambda_0^{J,+}(M) \oplus \Lambda^{J,-}(M).$$

For an ω -compatible almost-Kähler metric (J, g), the canonical hermitian connection on the complex tangent bundle (T(M), J, g) is defined by

$$\nabla_X Y = D_X^g Y - \frac{1}{2} J \left(D_X^g J \right) Y,$$

where D^g is the Levi-Civita connection with respect to g and X,Y are vector fields on M. Denote by R^{∇} the curvature of ∇ . Then, the hermitian Ricci form ρ^{∇} is the trace of $R_{X,Y}^{\nabla}$ viewed as an anti-hermitian linear operator of (T(M),J,g), i.e.

$$\rho^{\nabla}(X,Y) = -tr(J \circ R_{X,Y}^{\nabla}).$$

Hence, the 2-form ρ^{∇} is a closed (real) 2-form and it is a deRham representative of $2\pi c_1(T(M), J)$ in $H^2(M, \mathbb{R})$, where $c_1(T(M), J)$ is the first (real) Chern class. If

the almost-complex structure J is compatible with a symplectic form $\tilde{\omega}$ such that $\tilde{\omega}^n = e^F \omega^n$ for some smooth real-valued function F on M, then [26, 27]

$$\tilde{\rho}^{\nabla} = -\frac{1}{2}dJdF + \rho^{\nabla},$$

where $\tilde{\rho}^{\nabla}$ is the hermitian Ricci form of the almost-Kähler metric (J, \tilde{g}) (here $\tilde{g}(\cdot, \cdot) = \tilde{\omega}(\cdot, J \cdot)$ is the induced Riemannian metric).

We define the hermitian scalar curvature s^{∇} of an almost-Kähler metric (J, g) as the trace of ρ^{∇} with respect to ω , i.e.

$$(2.3) s^{\nabla} \omega^n = 2n \left(\rho^{\nabla} \wedge \omega^{n-1} \right).$$

The (Riemannian) Hodge operator $*_g: \Lambda^p(M) \to \Lambda^{2n-p}(M)$ is defined to be the unique isomorphism such that $\psi_1 \wedge (*_g \psi_2) = g(\psi_1, \psi_2) \frac{\omega^n}{n!}$, for any p-forms ψ_1, ψ_2 . Then, the codifferential δ^g , defined as the formal adjoint of the exterior derivative d with respect to g, is related to d by the relation [6, 15]

$$\delta^g = - *_q d *_q.$$

It follows that

$$(2.4) d = *_q \delta^g *_q.$$

In dimension 2n = 4, the bundle of 2-forms decomposes as

$$\Lambda^2(M) = \Lambda^+(M) \oplus \Lambda^-(M),$$

where $\Lambda^{\pm}(M)$ correspond to the eigenvalue (±1) under the action of the Hodge operator $*_q$. This decomposition is related to the splitting (2.1) as follows

(2.5)
$$\Lambda^{+}(M) = \mathbb{R} \cdot \omega \oplus \Lambda^{J,-}(M) \text{ and } \Lambda^{-}(M) = \Lambda_{0}^{J,+}(M).$$

3. Extremal almost-Kähler metrics

Let (M, ω) be a compact and connected symplectic manifold of dimension 2n. Any ω -compatible almost-complex structure is identified with the induced Riemannian metric.

Denote by AK_{ω} the Fréchet space of ω -compatible almost-complex structures. The space AK_{ω} comes naturally equipped with a formal Kähler structure. Let $Ham(M,\omega)$ be the group of hamiltonian symplectomorphisms of (M^{2n},ω) . The Lie algebra of $Ham(M,\omega)$ is identified with the space of smooth functions on M with zero mean value.

A key observation, made by Fujiki [13] in the integrable case and by Donaldson [9] in the general almost-Kähler case, asserts that the natural action of $Ham(M,\omega)$ on AK_{ω} is hamiltonian with momentum given by the hermitian scalar curvature. More precisely, the moment map $\mu: AK_{\omega} \to (Lie(Ham(M,\omega)))^*$ is

$$\mu_J(f) = \int_M s^{\nabla} f \frac{\omega^n}{n!}$$

where s^{∇} is the hermitian scalar curvature of (J,g) and f is a smooth function with zero mean value viewed as an element of $Lie(Ham(M,\omega))$. The square-norm of the

hermitian scalar curvature defines a functional on AK_{ω}

(3.1)
$$J \mapsto \int_{M} (s^{\nabla})^{2} \frac{\omega^{n}}{n!}.$$

Definition 3.1. The critical points (J, g) of the functional (3.1) are called *extremal* almost-Kähler metrics.

Proposition 3.2. An almost-Kähler metric (J,g) is a critical point of (3.1) if and only if $\operatorname{grad}_{\omega} s^{\nabla}$ is a Killing vector field with respect to g.

A proof of Proposition 3.2 is given in [4, 15, 22].

3.1. The extremal vector field. We fix a maximal torus T in $Ham(M,\omega)$ and denote by \mathfrak{t}_{ω} the finite dimensional space of real-valued smooth functions on M which are hamiltonians with zero mean value of elements of $\mathfrak{t}=Lie(T)$. Denote by Π^T_{ω} the L^2 -orthogonal projection of T-invariant smooth functions on \mathfrak{t}_{ω} with respect to the volume form $\frac{\omega^n}{n!}$. Let AK^T_{ω} be the space of ω -compatible T-invariant almost-complex structures. Given any $J\in AK^T_{\omega}$, we define $z^T_{\omega}:=\Pi^T_{\omega}s^{\nabla}$, where s^{∇} is the hermitian scalar curvature of (J,g). Then, we have the following (for more details see [3, 15, 22])

Proposition 3.3. The potential z_{ω}^{T} is independent of (J,g). Furthermore, a ω -compatible T-invariant almost-Kähler metric (J,g) is extremal if and only if

$$\mathring{s}^{\nabla} = z_{\omega}^{T},$$

where \mathring{s}^{∇} is the integral zero part of the hermitian scalar curvature s^{∇} of (J,g).

Definition 3.4. The vector field $Z_{\omega}^T := grad_{\omega} z_{\omega}^T$ is called the *extremal vector field* relative to T.

Proposition 3.5. The vector field Z_{ω}^{T} is invariant under T-invariant isotopy of ω .

Remark 3.6. The assumption that $T \subset Ham(M,\omega)$ is a maximal torus is used only in the second part of Proposition 3.3. Indeed, the arguments in [22] show that $z_{\omega}^T = \Pi_{\omega}^T s^{\nabla}$ is independent of (J,g) for any torus $T \subset Ham(M,\omega)$ and Proposition 3.5 still holds true for the corresponding vector field $Z_{\omega}^T = grad_{\omega} z_{\omega}^T$.

4. Almost-Kähler potentials in dimension 4

Let (M, ω) be a compact symplectic manifold of dimension 2n = 4 and (J, g) a ω -compatible almost-Kähler metric. In order to define the almost-Kähler potentials, we consider the following second order linear differential operator [23] on the smooth sections $\Omega^{J,-}(M)$ of the bundle of J-anti-invariant 2-forms.

$$\begin{array}{cccc} P: & \Omega^{J,-}(M) & \longrightarrow & \Omega^{J,-}(M) \\ & \psi & \longmapsto & (d\delta^g \psi)^{J,-}, \end{array}$$

where δ^g is the codifferential with respect to the metric g.

Lemma 4.1. P is a self-adjoint strongly elliptic linear operator with kernel the g-harmonic J-anti-invariant 2-forms.

Proof. The principal symbol of P is given by the linear map $\sigma(P)_{\xi}(\psi) = -\frac{1}{2}|\xi|^2\psi$, $\forall \xi \in T_x^*(M), \psi \in \Omega^{J,-}(M)$. So, P is a self-adjoint elliptic linear operator with respect to the global inner product $\langle \cdot, \cdot \rangle = \int_M g(\cdot, \cdot) \frac{\omega^2}{2}$. Now, let $\psi \in \Omega^{J,-}(M)$ and suppose that $P(\psi) = 0$. Then, $0 = \langle (d\delta^g \psi)^{J,-}, \psi \rangle = \langle d\delta^g \psi, \psi \rangle = \langle \delta^g \psi, \delta^g \psi \rangle$ which means that $\delta^g \psi = 0$. It follows from (2.5) and since ψ is J-anti-invariant that $*_g \psi = \psi$. Using the relation (2.4), we obtain $d\psi = *_g \delta^g *_g \psi = *_g \delta^g \psi = 0$. Hence, $d\psi = \delta^g \psi = 0$ and thus ψ is a g-harmonic J-anti-invariant 2-form.

Corollary 4.2. For $f \in C^{\infty}(M, \mathbb{R})$, there exist a unique $\psi_f \in \Omega^{J,-}(M)$ orthogonal to the kernel of P such that $(d\delta^g \psi_f)^{J,-} = (dJdf)^{J,-}$.

Proof. For a smooth real-valued function $f \in C^{\infty}(M, \mathbb{R})$ and any α in the kernel of P, we have $\langle (dJdf)^{J,-}, \alpha \rangle = \langle dJdf, \alpha \rangle = \langle Jdf, \delta^g \alpha \rangle = 0$. By a standard result of elliptic theory [6, 29] and since P is self-adjoint, there exist a smooth section $\psi_f \in \Omega^{J,-}(M)$ such that $P(\psi_f) = (dJdf)^{J,-}$. Moreover, ψ_f is unique if one requires ψ_f be orthogonal to the kernel of P.

From Corollary 4.2, it follows that, for $f \in C^{\infty}(M, \mathbb{R})$, the symplectic form $\omega_f = \omega + d(Jdf - \delta^g \psi_f)$ is a J-invariant closed 2-form. Then, the function f is called an almost-Kähler potential if the induced symmetric tensor $g_f(\cdot, \cdot) := \omega_f(\cdot, J \cdot)$ is a Riemannian metric. This notion of almost-Kähler potential is closely related but different (in general) from the one defined by Weinkove in [28]. More precisely, if the almost-complex structure J is compatible with a symplectic form $\tilde{\omega}$ which is cohomologous to ω i.e. $\tilde{\omega} - \omega = d\alpha$ (for some 1-form α), then the almost-Kähler potential defined by Weinkove is given by the function \tilde{f} which is uniquely determined (up to the addition of constant) by the Hodge decomposition of α with respect to the (self-adjoint elliptic) twisted Laplace operator $\tilde{\Delta}^c = J\Delta^{\tilde{g}}J^{-1}$, where $\Delta^{\tilde{g}}$ is the (Riemannian) Laplace operator with respect to the induced metric $\tilde{g}(\cdot, \cdot) = \tilde{\omega}(\cdot, J \cdot)$. In other words, we have the decomposition $\alpha = \alpha_{H^c} + \tilde{\Delta}^c \tilde{\mathbb{G}} \alpha$, where $\tilde{\mathbb{G}}$ is the Green operator associated to $\tilde{\Delta}^c$ and α_{H^c} is the harmonic part of α with respect to $\tilde{\Delta}^c$. Thus, $\tilde{f} = -\delta^{\tilde{g}}J\tilde{\mathbb{G}}\alpha$, where $\delta^{\tilde{g}}$ is the codifferential with respect to the metric \tilde{g} .

Thus, $\tilde{f} = -\delta^{\tilde{g}} J \, \tilde{\mathbb{G}} \alpha$, where $\delta^{\tilde{g}}$ is the codifferential with respect to the metric \tilde{g} . Note that $(dJdf)^{J,-} = D^g_{(df)^{\sharp_g}} \omega$ (see e.g. [15]), where \sharp_g stands for the isomorphism between $T^*(M)$ and T(M) induced by g^{-1} . Hence, in the Kähler case, $(dJdf)^{J,-} = 0$ which implies that $\psi_f = 0$ and thus this almost-Kähler potential coincides with the usual Kähler one.

5. Main Theorem

Let (M, ω) be a compact and connected symplectic manifold of dimension 2n = 4 and $J_t \in AK_\omega$ be a smooth path of ω -compatible almost-complex structures. We define the following family of differential operators associated to J_t

$$\begin{array}{cccc} P_t: & \Omega_0^2(M) & \longrightarrow & \Omega_0^2(M) \\ & \psi & \longmapsto & \frac{1}{2}\Delta^{g_t}\psi - \frac{1}{4}g_t(\Delta^{g_t}\psi,\omega)\omega, \end{array}$$

where $\Omega_0^2(M)$ is the space of smooth sections of the bundle $\Lambda_0^2(M)$ of primitive 2-forms (pointwise orthogonal to ω) and Δ^{g_t} is the (Riemannian) Laplacian with respect to the metric $g_t(\cdot, \cdot) = \omega(\cdot, J_t \cdot)$ (here we use the convention $g_t(\omega, \omega) = 2$).

One can easily check that P_t preserves the decomposition

$$\Omega_0^2(M) = \Omega_0^{J,+}(M) \oplus \Omega^{J,-}(M).$$

Furthermore,

$$P_t|_{\Omega^{J_t,-}(M)}(\psi) = (d\delta^{g_t}\psi)^{J_t,-} \text{ and } P_t|_{\Omega^{J_t,+}_0(M)}(\psi) = \frac{1}{2}\Delta^{g_t}\psi.$$

It follows that the kernel of P_t consists of primitive harmonic 2-forms which splits as anti-selfdual and J_t -anti-invariant ones so we have

$$\dim \ker(P_t) = b^-(M) + h_{J_t}^-,$$

where $h_{J_{\star}}^{-}$ is introduced by Drăghici–Li–Zhang in [12].

Moreover, $P_t - \frac{1}{2}\Delta^{g_t}$ is a linear differential operator of order 1. Indeed, a direct computation shows that

$$\left(P_t - \frac{1}{2}\Delta^{g_t}\right)(\psi) = \frac{1}{2} \left[\frac{1}{2}\delta^{g_t}\left(D^{g_t}\omega(\psi)\right) - \frac{1}{2}g_t(D^{g_t}\psi, D^{g_t}\omega) + \frac{s_{g_t}}{6}g_t(\omega, \psi) - W^{g_t}(\omega, \psi)\right]\omega,$$

where W^{g_t} stands for the Weyl tensor (see e.g. [6]), D^{g_t} (resp. δ^{g_t}) for the Levi-Civita connection (resp. the codifferential) with respect to the metric g_t and s_{g_t} for the Riemannian scalar curvature defined as the trace of the (Riemannian) tensor.

The operator P_t is a self-adjoint strongly elliptic linear operator of order 2. We obtain then a family of Green operators \mathbb{G}_t associated to P_t . If $h_{J_t}^- = h_{J_0}^- = b^+(M) - 1$ for $t \in (-\epsilon, \epsilon)$, then \mathbb{G}_t is C^{∞} differentiable in $t \in (-\epsilon, \epsilon)$ [19, 20], meaning that $\mathbb{G}_t(\psi_t)$ is a smooth family of sections of $\Lambda_0^2(M)$ for any smooth sections ψ_t .

To show Theorem 1.1, we consider the extension of \mathbb{G}_t to the Sobolev spaces $W^{k,p}(M,\Lambda^2_0(M))$ involving derivatives up to k.

Lemma 5.1. Let $\mathbb{G}_t: \Omega_0^2(M) \to \Omega_0^2(M)$ the family of the above Green operators associated to P_t and suppose that $h_{J_t}^- = h_{J_0}^- = b^+(M) - 1$ for $t \in (-\epsilon, \epsilon)$. Then, the extension of \mathbb{G}_t to Sobolev spaces, still denoted by \mathbb{G}_t , defines a C^1 map $\mathbb{G}: (-\epsilon, \epsilon) \times W^{p,k}(M, \Lambda_0^2(M)) \to W^{p,k+2}(M, \Lambda_0^2(M))$

Proof. Denote by Π_t the L^2 -orthogonal projection to the kernel of P_t with respect to $\langle \cdot, \cdot \rangle_{L^2_{g_t}} = \int_M g_t(\cdot, \cdot) \frac{\omega^2}{2}$. We claim that $\mathbb{G}_t \circ \Pi_0$ and $\Pi_0 \circ \mathbb{G} : (-\epsilon, \epsilon) \times W^{k,p}(M, \Lambda_0^2(M)) \to W^{k+2,p}(M, \Lambda_0^2(M))$ are C^1 maps. Indeed, let $\{\psi_0^i\}$ be an orthonormal basis of the kernel of P_0 with respect to $\langle \cdot, \cdot \rangle_{L^2_{g_0}}$. Note that ψ_0^i are smooth since P_0 is elliptic.

Then, we have

$$\begin{aligned}
&(\mathbb{G}_{t} \circ \Pi_{0}) (\psi) &= \sum_{i} \left\langle \psi, \psi_{0}^{i} \right\rangle_{L_{g_{0}}^{2}} \mathbb{G}_{t} (\psi_{0}^{i}), \\
&(\Pi_{0} \circ \mathbb{G}_{t}) (\psi) &= \sum_{i} \left\langle \mathbb{G}_{t} (\psi), (\psi_{0}^{i})^{J_{0},+} + (\psi_{0}^{i})^{J_{0},-} \right\rangle_{L_{g_{0}}^{2}} \psi_{0}^{i} \\
&= \sum_{i} \left(\int_{M} -\mathbb{G}_{t} (\psi) \wedge (\psi_{0}^{i})^{J_{0},+} + \mathbb{G}_{t} (\psi) \wedge (\psi_{0}^{i})^{J_{0},-} \right) \psi_{0}^{i} \\
&= \sum_{i} \left(\int_{M} -\mathbb{G}_{t} (\psi) \wedge ((\psi_{0}^{i})^{J_{0},+})^{J_{t},+} - \mathbb{G}_{t} (\psi) \wedge ((\psi_{0}^{i})^{J_{0},+})^{J_{t},-} \right) \\
&+ \mathbb{G}_{t} (\psi) \wedge ((\psi_{0}^{i})^{J_{0},-})^{J_{t},+} + \mathbb{G}_{t} (\psi) \wedge ((\psi_{0}^{i})^{J_{0},-})^{J_{t},-} \right) \psi_{0}^{i} \\
&= \sum_{i} \left[\left\langle \psi, \mathbb{G}_{t} \left(((\psi_{0}^{i})^{J_{0},+})^{J_{t},+} \right) \right\rangle_{L_{g_{t}}^{2}} - \left\langle \psi, \mathbb{G}_{t} \left(((\psi_{0}^{i})^{J_{0},+})^{J_{t},-} \right) \right\rangle_{L_{g_{t}}^{2}} \right] \psi_{0}^{i} \end{aligned}$$

(in the latter equality, we used the fact that \mathbb{G}_t is self-adjoint with respect to $L_{g_t}^2$). The claim follows from the result of Kodaira–Spencer [19, 20].

Denote by $W^{k,p}(M, \Lambda_0^2(M))^{\perp}$ the space of 2-forms in $W^{k,p}(M, \Lambda_0^2(M))$ which are orthogonal to the kernel of P_0 with respect to $L_{g_0}^2$ and consider the map

$$\begin{array}{ccccc} \Phi: & (-\epsilon,\epsilon) \times W^{k+2,p}(M,\Lambda_0^2(M))^{\perp} & \longrightarrow & (-\epsilon,\epsilon) \times W^{k,p}(M,\Lambda_0^2(M))^{\perp} \\ & & (t,\psi) & \longmapsto & (t,(Id-\Pi_0)P_t(\psi)), \end{array}$$

Clearly, the map Φ is of class C^1 and its differential at $(0, \psi)$ is an isomorphism so by the *inverse function theorem* for Banach spaces there exist a neighboorhood V of $(0, \psi)$ such that $\Phi|_V$ admits an inverse of class C^1 . By the The Kodaira–Spencer result [19, 20], the map $\Pi: (-\epsilon, \epsilon) \times W^{k,p}(M, \Lambda_0^2(M)) \to W^{k,p}(M, \Lambda_0^2(M))$ is C^1 and thus the map $P_t(Id - \Pi_0) \mathbb{G}_t(Id - \Pi_0) = (Id - \Pi_t)(Id - \Pi_0) - P_t(\Pi_0 \mathbb{G}_t)(Id - \Pi_0)$ is clearly C^1 since it is a composition of such operators. Then, the map

$$\Phi|_{V}^{-1}(t, (Id - \Pi_{0})P_{t}(Id - \Pi_{0})\mathbb{G}_{t}(Id - \Pi_{0})) = (t, (Id - \Pi_{0})\mathbb{G}_{t}(Id - \Pi_{0}))
= (t, \mathbb{G}_{t} - \Pi_{0}\mathbb{G}_{t} - \mathbb{G}_{t}\Pi_{0} + \Pi_{0}\mathbb{G}_{t}\Pi_{0})$$

is C^1 and hence \mathbb{G}_t is C^1 .

Proof of Theorem 1.1 Let (M, ω) be a 4-dimensional compact and connected symplectic manifold and T a maximal torus in $Ham(M, \omega)$. Let (J_t, g_t) a smooth family of ω -compatible almost-Kähler metrics in AK_{ω}^T such that (J_0, g_0) is an extremal Kähler metric.

Following [21], we consider the almost-Kähler deformations

$$\omega_{t,f} = \omega + d(J_t df - \delta^{g_t} \psi_f^t),$$

where f belongs to the Fréchet space $\widetilde{C}_T^{\infty}(M,\mathbb{R})$ of T-invariant smooth functions (with zero integral), which are L^2 -orthogonal, with respect to $\frac{\omega^2}{2}$, to \mathfrak{t}_{ω} and where the 2-form ψ_f^t is given by Corollary 4.2.

Let \mathcal{U} be an open set in $\mathbb{R} \times \widetilde{C}_T^{\infty}(M, \mathbb{R})$ containing (0,0) such that the symmetric tensor $g_{t,f}(\cdot,\cdot) := \omega_{t,f}(\cdot,J_t\cdot)$ is a Riemannian metric.

By possibly replacing \mathcal{U} with a smaller open set, we may assume as in [21] that the kernel of the operator $(Id - \Pi_{\omega}^T) \circ (Id - \Pi_{\omega_{t,f}}^T)$ is equal to the kernel of $(Id - \Pi_{\omega_{t,f}}^T)$. Indeed, let $\{X_1, \dots, X_n\}$ be a basis of $\mathfrak{t} = Lie(T)$. Then, the corresponding hamiltonians with zero mean value $\{\xi_{\omega}^1, \dots, \xi_{\omega}^n\}$ resp. $\{\xi_{\omega_{t,f}}^1, \dots, \xi_{\omega_{t,f}}^n\}$, with respect to ω resp. $\omega_{t,f}$, form a basis of \mathfrak{t}_{ω} resp. $\mathfrak{t}_{\omega_{t,f}}$. Let $\{\tilde{\xi}_{\omega}^1, \dots, \tilde{\xi}_{\omega}^n\}$ resp. $\{\tilde{\xi}_{\omega_{t,f}}^1, \dots, \tilde{\xi}_{\omega_{t,f}}^n\}$ the corresponding orthonormal basis obtained by the Gram–Schmidt procedure. Since $\det \left[\left\langle \tilde{\xi}_{\omega}^i, \tilde{\xi}_{\omega_{t,f}}^j \right\rangle \right]$ defines a continuous function on \mathcal{U} , then we may suppose that $\det \left[\left\langle \tilde{\xi}_{\omega}^i, \tilde{\xi}_{\omega_{t,f}}^j \right\rangle \right] \neq 0$ on an eventually smaller open set than \mathcal{U} (here $\langle \cdot, \cdot \rangle$ denotes the L^2 product with respect to the volume form $\frac{\omega_{t,f}^2}{2}$). So, if

poduct with respect to the volume form $\frac{-i,j}{2}$). So, i

$$u \in ker\left((Id - \Pi_{\omega}^T) \circ (Id - \Pi_{\omega_{t,f}}^T)\right)$$

then $v \in \mathfrak{t}_{\omega} \cap (\mathfrak{t}_{\omega_{t,f}})^{\perp_{g_{t,f}}}$, where $v = (Id - \Pi_{\omega_{t,f}}^T)u$. But the hypothesis

$$\det\left[\left\langle \tilde{\xi}_{\omega}^{i}, \tilde{\xi}_{\omega_{t,f}}^{j} \right\rangle\right] \neq 0$$

implies that $v \equiv 0$ and then $ker\left((Id - \Pi_{\omega}^T) \circ (Id - \Pi_{\omega_{t,f}}^T)\right) = ker(Id - \Pi_{\omega_{t,f}}^T)$. We then consider the map:

$$\Psi: \qquad \mathcal{U} \longrightarrow \mathbb{R} \times \widetilde{C}_{T}^{\infty}(M, \mathbb{R})$$

$$(t, f) \longmapsto \left(t, (Id - \Pi_{\omega}^{T}) \circ (Id - \Pi_{\omega_{t, f}}^{T}) (\mathring{s}^{\nabla_{t, f}})\right),$$

where $\mathring{s}^{\nabla_{t,f}}$ is the zero integral part of the hermitian scalar curvature $s^{\nabla_{t,f}}$ of $(J_t, g_{t,f})$.

It follows from Proposition 3.3 that $\Psi(t, f) = (t, 0)$ if and only if $(J_t, g_{t,f})$ is an extremal almost-Kähler metric. In particular, $\Psi(0, 0) = (0, 0)$.

Let $\alpha_{t,f} = J_t df - \delta^{g_t} \psi_f^t = J_t df - \delta^{g_t} \mathbb{G}_t \left((dJ_t df)^{J_t,-} \right) = J_t df - \delta^{g_t} \mathbb{G}_t (D_{df}^{g_t}_{g_t} \omega),$ where \mathbb{G}_t is the Green operator associated to the elliptic operator $P_t : \Omega^{J_t,-}(M) \to \Omega^{J_t,-}(M)$. In order to extend the map Ψ to Sobolev spaces, we give an explicit expression of $(Id - \Pi_{\omega_{t,f}}^T)(s^{\nabla_{t,f}})$. A direct computation using (2.2) shows that

(5.1)
$$s^{\nabla_{t,f}} = \Delta^{g_{t,f}} F_{t,f} + g_{t,f}(\rho^{\nabla_t}, \omega_{t,f}),$$

where $F_{t,f} = \log \left(\frac{1}{2} \left(\left(1 + g_t \left(d\alpha_{t,f}, \omega\right)\right)^2 + 1 - g_t \left(d\alpha_{t,f}, d\alpha_{t,f}\right) \right) \right)$ satisfying the relation $\omega_{t,f}^2 = e^{F_{t,f}} \omega^2$. Then

$$(5.2) \qquad (Id - \Pi_{\omega_{t,f}}^T)(s^{\nabla_{t,f}}) = \Delta^{g_{t,f}} F_{t,f} + g_{t,f}(\rho^{\nabla_t}, \omega_{t,f}) - \sum_{s} \left\langle s^{\nabla_{t,f}}, \tilde{\xi}_{\omega_{t,f}}^j \right\rangle \tilde{\xi}_{\omega_{t,f}}^j.$$

Let $\widetilde{W}_T^{p,k}$ be the completion of $\widetilde{C}_T^\infty(M,\mathbb{R})$ with respect to the Sobolev norm $\|\cdot\|_{p,k}$ involving derivatives up to order k. We choose p,k such that pk>2n and the corresponding Sobolev space $\widetilde{W}_T^{p,k}\subset C_T^3(M,\mathbb{R})$ so that all coefficients are $C_T^0(M,\mathbb{R})$. Since $\widetilde{W}_T^{p,k}$ form an algebra relative to the standard multiplication of functions [1], we deduce from the expression (5.2) that the extension of Ψ to the Sobolev completion of $\widetilde{C}_T^\infty(M,\mathbb{R})$ is a map $\Psi^{(p,k)}:\widetilde{\mathcal{U}}\subset\mathbb{R}\times\widetilde{W}_T^{p,k+4}\longrightarrow\mathbb{R}\times\widetilde{W}_T^{p,k}$.

Clearly $\Psi^{(p,k)}$ is a C^1 map (in a small enough open around (0,0)). Indeed, it is obtained by a composition of C^1 maps by Lemma 5.1 and (5.2).

As in [21] and using Proposition 3.3, the differential of $\Psi^{(p,k)}$ at (0,0) is given by

$$\left(\mathbf{T}_{(0,0)}\Psi^{(p,k)}\right)(t,f) = \left(t, t\delta^{g_0}\delta^{g_0}h - 2\delta^{g_0}\delta^{g_0}(D^{g_0}df)^{J_0,-}\right),$$

where $h = \frac{d}{dt}|_{t=0} g_t$.

The operator $L:=\frac{\partial \Psi}{\partial f}|_{(0,0)}$ given by $L(f)=-2\delta^{g_0}\delta^{g_0}(D^{g_0}df)^{J_0,-}$ is called the Lichnerowicz operator. It is a 4-th order self-adjoint T-invariant elliptic linear operator leaving invariant $(\mathfrak{t}_\omega)^\perp$ since L(f)=0 for any $f\in\mathfrak{t}_\omega$. By a known result of the elliptic theory [6,29], we obtain the L^2 -orthogonal splitting $\widetilde{C}_T^\infty(M,\mathbb{R})=\ker(L)\oplus Im(L)$. Following the argument in [3, Lemma 4], any $f\in\ker(L)$ gives rise to a Killing vector field in the centralizer of $\mathfrak{t}=Lie(T)$. By the maximality of the torus $T,f\in\mathfrak{t}_\omega$. It follows that L is an isomorphism of $\widetilde{C}_T^\infty(M,\mathbb{R})$ and also from $\widetilde{W}_T^{p,k+4}$ to $\widetilde{W}_T^{p,k}$. Thus, $\mathbf{T}_{(0,0)}\Psi^{(p,k)}$ is an isomorphism from $\mathbb{R}\oplus\widetilde{W}_T^{p,k+4}$ to $\mathbb{R}\oplus\widetilde{W}_T^{p,k}$. It follows from the inverse function theorem for Banach manifolds that $\Psi^{(p,k)}$ determines an isomorphism from an open neighbourhood V of (0,0) to an open neighbourhood of (0,0). In particular, there exists $\mu>0$ such that for $|t|<\mu$, $\Psi^{(p,k)}(\Psi^{(p,k)}|_V^{-1}(t,0))=(t,0)$. By Sobolev embedding, we can choose a k large enough, such that $\widetilde{W}_T^{p,k+4}\subset\widetilde{C}_T^6(M,\mathbb{R})$. Thus, for $|t|<\mu$, $(J_t,g_{\Psi^{(p,k)}|_V^{-1}(t,0)})$ is an extremal almost-Kähler metric of regularity at least C^4 (so we ensure, in this case, that $\operatorname{grad}_\omega s^{\nabla_{t,f}}$ is of regularity C^1).

By Proposition 3.5, the extremal vector field $Z_{\omega_{t,f}}^T = Z_{\omega}^T$ is smooth for any almost-Kähler metric $(J_t, g_{t,f})$. In particular, for an extremal almost-Kähler metric $(J_t, g_{t,f})$ of regularity C^4 , the dual $ds^{\nabla_{t,f}}$ of Z_{ω}^T with respect to $\omega_{t,f}$ is of regularity C^4 , then the hermitian scalar curvature $s^{\nabla_{t,f}}$ of $(J_t, g_{t,f})$ is of regularity C^5 . From (5.1), it follows that the hermitian scalar curvature is given by the pair of equations

$$(5.3) s^{\nabla_{t,f}} - g_{t,f}(\rho^{\nabla_t}, \omega_{t,f}) = \Delta^{g_{t,f}}(u),$$

$$(5.4) e^u = \frac{\omega_{t,f}^2}{\omega^2}.$$

From (5.3), using the ellipticity [6] of the (Riemannian) Laplacian $\Delta^{g_{t,f}}$ and since the l.h.s of (5.3) is of Hölder class $C^{3,\beta}$ for any $\beta \in (0,1)$, it follows that u is of class $C^{5,\beta}$. Following [11, 28], the linearisation of the equation (5.4) $(\omega + d\alpha) \wedge d\dot{\alpha} = 0$ together with the constraints $\delta^{g_t}\dot{\alpha} = 0$ and $(d\dot{\alpha})^{J_t,-} = 0$ form a linear elliptic system in $\dot{\alpha}$. Elliptic theory [2, 6] ensures that the almost-Kähler metric $g_{t,f}$ is of class $C^{5,\beta}$ as the volume form and we can prove that any extremal almost-Kähler metric of regularity C^4 is smooth by a bootstraping argument (in the Kähler case see [21]).

We obtain then a smooth family of T-invariant extremal almost-Kähler structures $(J_t, \omega_t = \omega + d\alpha_t)$ defined for $|t| < \mu$. The main theorem follows from the Moser Lemma [25].

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