

The Hamiltonian tube of a cotangent-lifted action

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The Marle-Guillemin-Sternberg (MGS) form is local model for a neighborhood of an orbit of a Hamiltonian Lie group action on a symplectic manifold. One of the main features of the MGS form is that it puts simultaneously in normal form the existing symplectic structure and momentum map. The main drawback of the MGS form is that it does not have an explicit expression. We will obtain a MGS form for cotangent- lifted actions on cotangent bundles that, in addition to its defining features, respects the additional fibered structure present. This model generalizes previous results obtained by T. Schmah for orbits with fully-isotropic momentum. In addition, our construction is explicit up to the integration of a differential equation on G . This equation can be easily solved for the groups $SO(3)$ or $SL(2)$, thus giving explicit symplectic coordinates for arbitrary canonical actions of these groups on any cotangent bundle.

1. Introduction

The study of the local geometry of symplectic manifolds equipped with Hamiltonian group actions constitutes a field which originated with the classical papers [6, 11]. In these references the authors obtain a universal model for a tubular neighborhood of the orbit of a point under a Hamiltonian action which puts in normal form both the symplectic structure and the momentum map (Theorem 2.1). This model is known as the *Hamiltonian tube* or *Marle-Guillemin-Sternberg form* and it is the base of almost all the local studies concerning Hamiltonian actions of Lie groups on symplectic manifolds. In fact, since the decade of the 80's almost all the relevant results about the qualitative local dynamics of equivariant Hamiltonian flows have been obtained using techniques based on the Hamiltonian tube, for example [8, 10, 13–18, 21, 24, 25, 30]. Moreover, the Hamiltonian tube was

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a major ingredient for many of the generalizations of the Marsden-Weinstein reduction scheme to singular actions like [2, 9, 19, 31].

In this paper we will focus on the cotangent bundle case. Let T^*Q be a cotangent bundle equipped with its canonical symplectic structure ω_Q and let G be a Lie group that acts smoothly on Q . The canonical lift of this action to T^*Q is automatically a Hamiltonian action. The Marle-Guillemin-Sternberg construction (Theorem 2.1) applied to a canonical action over a cotangent bundle gives, as for every Hamiltonian action, an equivariant local model of (T^*Q, ω_Q) that puts in normal form both the symplectic structure and the momentum map. However, in general this model does not respect the fibration $T^*Q \rightarrow Q$. Moreover, the map given by Theorem 2.1 is not constructive and only some of its properties are known. In the concrete case of cotangent bundles there is a strong motivation coming from geometric mechanics and geometric quantization to obtain explicit or fibred local models. In this paper we obtain a construction of the Hamiltonian tube for a canonical cotangent-lifted action in a cotangent bundle specially adapted to this kind of manifold and that puts the fibration in a normal form (Theorem 5.6). In other words, this assumes that the space that models locally the neighborhood of an orbit of the group in T^*Q has a fibred structure $\tau : Y \rightarrow U$ where U is a local model of the base Q such that

$$\mathcal{T} : Y \rightarrow T^*U$$

is a fibred map. Additionally, the construction of \mathcal{T} will be explicit up to the integration of a differential equation on G . The restricted G -tubes (Definition 4.7) will be the basic building blocks and are the only non-explicit part of the model. Given a Hamiltonian action the restricted G -tube depends only on the group G and its algebraic structure. For example, for $SO(3)$ and $SL(2)$ the expressions of their restricted G -tubes can be obtained explicitly, see Section 7. For larger groups the computation will be more cumbersome but could be done with a computer algebra system. As an additional result of our construction we obtain a fibred analogue of the Lerman-Bates lemma [2] (Proposition 6.1) that characterizes the set of points with prescribed momentum in a neighborhood of the form $\tau^{-1}(U)$. We believe that this result can be used to study in detail the structure of singular reduction for cotangent bundles, generalizing the results of [22] to non-zero momentum. This will be addressed elsewhere.

The first works studying symplectic normal forms in the specific case of cotangent bundles seem to have been [27, 28]. In these references T. Schmah found a Hamiltonian tube around those points $z \in T^*Q$ such that

its momentum $\mu = \mathbf{J}(z)$ is fully-isotropic (that is, $G_\mu = G$ with respect to the coadjoint representation). One of the main differences between her construction and the classical MGS model for symplectic actions is that the one for cotangent bundles was constructive, unlike the general MGS model. The next step came with [23] which provides a general description of the symplectic slice of a cotangent bundle, without the assumption $G = G_\mu$. Recently [29] constructed Hamiltonian tubes for free actions of a Lie group G and showed that this construction can be made explicit for $G = SO(3)$.

This paper is organized as follows: in Section 2 we review some background material regarding proper actions and the classical MGS model. In Section 3 we introduce in Proposition 3.1 a splitting of the Lie algebra that will be needed in all the subsequent development. This splitting already appeared partially in [23] (Theorem 3.4). In Section 4 we introduce simple and restricted G -tubes (Definitions 4.2 and 4.7). Simple G -tubes are, up to technical details, MGS models for the lift of the left action of G on itself to T^*G . Their existence is proved in Proposition 4.3. Restricted G -tubes are defined implicitly in terms of a simple G -tube (Proposition 4.8) and are the technical tool that we will need later to construct the general Hamiltonian tube.

In Section 5 we construct the general Hamiltonian tube for a cotangent-lifted action in such a way that it is explicit up to a restricted G -tube. This general tube will be the composition of two maps. We will first construct a Hamiltonian tube around points in T^*Q with certain maximal isotropy properties (Theorem 5.2) and then an adaptation of the ideas of [28] will be used to construct a Γ map (Proposition 5.4). Together these two maps will give the general Hamiltonian tube in Theorem 5.6. In Section 6 we use a zero section-centered tube to construct a cotangent-bundle version of a result due to Bates and Lerman in Proposition 6.2. One important novelty of this lemma is that our version is global in the vertical direction in the sense that given $z \in \mathbf{J}^{-1}(\mu)$ we can describe, via a Hamiltonian tube, not only a set of the form $\mathbf{J}^{-1}(\mu) \cap U$ where U is a neighborhood of z , but a set $\tau^{-1}(\tau(U)) \cap \mathbf{J}^{-1}(\mu)$ with U is a neighborhood of z . That is, with one Hamiltonian tube we can describe all the points in $\mathbf{J}^{-1}(\mu)$ whose projection is close enough to the projection of the center point. Finally, in Section 7 we present explicit examples of G -tubes for both the groups $SO(3)$ (where we recover the results of [29]) and $SL(2, \mathbb{R})$. In Subsection 7.4 we present an explicit Hamiltonian tube for the natural action of $SO(n)$ on $T^*\mathbb{R}^n$ which generalizes the final example of [28] to the case $\mu \neq 0$.

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2. Preliminaries

This section collects background material from the theory of Hamiltonian actions and MGS normal forms that will be used through this paper. Most material is standard and can be found in greater detail in several references, for instance [19].

2.1. Proper actions and slices

Let G be a Lie group with Lie algebra \mathfrak{g} . We will always denote by e the identity element of the group and by $L_g, R_g; G \rightarrow G$ the left and right multiplications by g respectively. If G acts on M we say that M is a G -space. For $p \in M$ the *isotropy subgroup* of p is

$$G_p = \{g \in G \mid g \cdot p = p\}.$$

A map $f : M_1 \rightarrow M_2$ between two manifolds endowed with G -actions is called G -equivariant if $f(g \cdot p) = g \cdot f(p)$ for all $p \in M_1$ and $g \in G$.

An action is *proper* if the map $G \times M \rightarrow M \times M$ defined by $(g, z) \mapsto (z, g \cdot z)$ is a proper map. For a proper G -action all the isotropy subgroups G_z are compact subgroups of G . For any Lie group G the left and right actions on itself are proper.

If a compact subgroup $H \subset G$ acts on a manifold A then on $G \times A$ we can consider two actions:

- *twisting action* of H : $h \cdot^T (g, a) = (gh^{-1}, h \cdot a)$, $h \in H$
- *left action* of G : $h \cdot^L (g, a) = (hg, a)$, $h \in G$.

If necessary we will use as above the superindexes T or L to indicate the H^T -action (twisting) or the G^L -action (left) on the product $G \times A$. As both actions commute $G \times A$ supports an action of the direct product group $G^L \times H^T$.

The twisting action is free and proper and, therefore, the quotient space $(G \times A)/H^T$ is a manifold and it will be called the *twisted product*. We will denote it as $G \times_H A$ and its elements will be denoted as $[g, a]_H$ $g \in G, a \in A$. The twisted product $G \times_H A$ admits a proper G -action given by $g \cdot [g', a]_H = [gg', a]_H$. In fact the twisted product is exactly the associated bundle for the principal H -bundle $G \rightarrow G/H$ and the H -manifold A .

The Tube Theorem [20] shows that in fact every proper G -space is locally a twisted product. That is, if G acts properly on $M, z \in M$ and A is an G_z -invariant complement of $\mathfrak{g} \cdot z$ in T_zM then there is a G -equivariant diffeomorphism

$$(1) \quad \mathfrak{s} : G \times_{G_z} A \longrightarrow U \subset M$$

satisfying $\mathfrak{s}([e, 0]_{G_z}) = z$ where U is a G -invariant neighborhood of z .

A G_z -invariant complement of $\mathfrak{g} \cdot z$ in T_zM will be called a *linear slice* for the G action at $z \in M$.

2.2. Hamiltonian actions and Hamiltonian tubes

Assume now that G acts symplectically on a symplectic manifold (M, ω) . A momentum map is a function $\mathbf{J} : M \rightarrow \mathfrak{g}^*$ such that

$$i_{\xi_M} \omega = \mathbf{d}\langle \mathbf{J}(\cdot), \xi \rangle \quad \forall \xi \in \mathfrak{g}$$

where $\xi_M \in \mathfrak{X}(M)$ is the fundamental vector field associated with $\xi \in \mathfrak{g}$. If \mathbf{J} is equivariant with respect to the coadjoint action on \mathfrak{g}^* then we will say that the action is *Hamiltonian*. If G acts Hamiltonially on a symplectic manifold (M, ω) there is a symplectic version of the Tube Theorem for proper G -spaces, and this is precisely the content of the *Marle-Guillemin-Sternberg normal form* proven by Marle, Guillemin and Sternberg in [6, 11] for compact groups and extended to proper actions of arbitrary groups in [2].

Theorem 2.1 (Hamiltonian Tube Theorem). *Let (M, ω) be a symplectic manifold endowed with a proper Hamiltonian action of a Lie group G with momentum map $\mathbf{J} : M \rightarrow \mathfrak{g}^*$. Let $z \in M, \mu = \mathbf{J}(z)$, and choose a G_z -invariant splitting $\mathfrak{g}_\mu = \mathfrak{g}_z \oplus \mathfrak{m}$. Let N be a G_z -invariant complement of $\mathfrak{g}_\mu \cdot z$ in $\text{Ker } T_z\mathbf{J}$ and $\mathbf{J}_N : N \rightarrow \mathfrak{g}_z^*$ defined by $\langle \mathbf{J}_N(v), \xi \rangle := \frac{1}{2}\omega(\xi \cdot v, v)$.*

Consider the set $Y := G \times_{G_z} (\mathfrak{m}^* \times N)$ equipped with the two-form

$$(2) \quad \begin{aligned} &\Omega_Y(T_{(g,\nu,v)}\pi_{G_z}(u_1), T_{(g,\nu,v)}\pi_{G_z}(u_2)) \\ &= \langle \dot{\nu}_2 + T_v\mathbf{J}_N(\dot{\nu}_2), \xi_1 \rangle - \langle \dot{\nu}_1 + T_v\mathbf{J}_N(\dot{\nu}_1), \xi_2 \rangle \\ &\quad + \langle \nu + \mathbf{J}_N(v) + \mu, [\xi_1, \xi_2] \rangle + \omega(\dot{\nu}_1, \dot{\nu}_2) \end{aligned}$$

where $u_i = (T_e L_g \xi_i; \dot{\nu}_i, \dot{\nu}_i) \in T_{(g,\nu,v)}(G \times \mathfrak{m}^* \times N)$ and $\pi_{G_z} : G \times (\mathfrak{m}^* \times N) \rightarrow G \times_{G_z} (\mathfrak{m}^* \times N)$. There is a neighborhood Y_r of the zero section of Y such that the restriction (Y_r, Ω_Y) is a symplectic manifold equipped with a Hamiltonian action of G (seen as a twisted product) for which the momentum map is

$$(3) \quad \mathbf{J}_Y[g, \nu, v]_{G_z} = \text{Ad}_{g^{-1}}^*(\mu + \nu + \mathbf{J}_N(v)).$$

Additionally, there is a map

$$\mathcal{T} : Y_r \longrightarrow M$$

such that:

- $\mathcal{T} : Y_r \longrightarrow \mathcal{T}(Y_r) \subset M$ is a G -equivariant diffeomorphism with $\mathcal{T}([e, 0, 0]_{G_z}) = z$.
- $\mathcal{T}^*\omega = \Omega_Y$.

The pair (Y_r, Ω_Y) is called the MGS model at $z \in M$, the G -equivariant symplectomorphism \mathcal{T} is called a Hamiltonian tube around z and the space N a symplectic slice at z .

Note that whereas in the Palais model (1) the twisted product depends only on a G_z -invariant complement to $\mathfrak{g} \cdot z$ in $T_z M$ in the Hamiltonian tube the normal form depends on a G_z -invariant complement N of $\mathfrak{g}_\mu \cdot z$ in $\text{Ker } T_z \mathbf{J}$ and on the complement \mathfrak{m} of \mathfrak{g}_z in \mathfrak{g}_μ . Note also that N , equipped with the restriction of $\omega(z)$ is a symplectic linear space supporting a linear Hamiltonian representation of G_z which admits \mathbf{J}_N as equivariant momentum map. In addition, since \mathcal{T} is a symplectomorphism equivariant with respect to the Hamiltonian actions of G on Y_r and M , by general geometric arguments we have

$$\mathbf{J}_Y = \mathbf{J} \circ \mathcal{T}$$

when the above expression is well defined.

2.3. The MGS model as a reduced space

We can interpret the symplectic form Ω_Y in (2) as the reduced symplectic form for a more basic structure. In this section we are going to recall some well-known facts about this interpretation of the MGS model that will be used throughout the paper. Let G be a Lie group, $\mu \in \mathfrak{g}^*$ and $K \subset G_\mu$ a compact subgroup. Since K is compact we can choose a K -invariant complement of \mathfrak{g}_μ in \mathfrak{g} and this choice induces a K -equivariant linear inclusion $\iota : \mathfrak{g}_\mu^* \rightarrow \mathfrak{g}^*$. Consider the product $T_\mu := G \times \mathfrak{g}_\mu^*$ and the map $T_\mu \rightarrow T^*G$ given by $(g, \nu) \mapsto T_e L_{g^{-1}}(\mu + \iota(\nu)) \in T^*G$. With this map we can pull-back the canonical symplectic form of T^*G obtaining the two-form ω_{T_μ} given by

$$\omega_{T_\mu}(g, \nu)(v_1, v_2) = \langle \dot{v}_2, \xi_1 \rangle - \langle \dot{v}_1, \xi_2 \rangle + \langle \mu + \iota(\nu), [\xi_1, \xi_2] \rangle,$$

where $v_i = (T_e L_g \xi_i, \dot{v}_i) \in T_{(g, \nu)}G \times \mathfrak{g}_\mu^*$. This form satisfies $\omega_{T_\mu} = -\mathbf{d}\theta_{T_\mu}$ where

$$(4) \quad \theta_{T_\mu}(g, \nu)(v_1) = \langle \mu + \iota(\nu), \xi_1 \rangle.$$

It can be checked that, for any $g \in G$ the two-form $\omega_{T_\mu}(g, 0)$ is non-degenerate and, therefore there is an open K -invariant neighborhood $(\mathfrak{g}_\mu^*)_r$ of $0 \in \mathfrak{g}_\mu^*$ such that $(G \times (\mathfrak{g}_\mu^*)_r, \omega_{T_\mu})$ is a symplectic space (see Proposition 7.2.2 of [19]). Let (N, ω_N) be a symplectic linear space with a K -Hamiltonian linear action with momentum map

$$(5) \quad \langle \mathbf{J}_N(v), \xi \rangle := \frac{1}{2} \omega_N(\xi \cdot v, v).$$

The product $Z := G \times ((\mathfrak{g}_\mu^*)_r \times N)$ equipped with the two-form $\omega_{T_\mu} + \omega_N$ is a symplectic space and the natural G^L and K^T -actions are free and Hamiltonian with momentum maps

$$(6) \quad \mathbf{K}_{K^T}(g, \nu, v) = -\nu|_{\mathfrak{k}} + \mathbf{J}_N(v), \quad \text{and} \quad \mathbf{K}_{G^L}(g, \nu, v) = \text{Ad}_{g^{-1}}^* \nu.$$

By the Marsden-Weinstein reduction procedure [12] the quotient $\mathbf{K}_{K^T}^{-1}(0)/K^T$ is a symplectic manifold. Since the G^L and K^T actions commute then the induced G -action on this quotient is also Hamiltonian.

Let now \mathfrak{m} be a K -invariant complement of $\mathfrak{k} = \text{Lie}(K)$ in \mathfrak{g}_μ . There are small enough open neighborhoods \mathfrak{m}_r^* and N_r of the origin in \mathfrak{m}^* and N such

that $\nu + \mathbf{J}_N(v) \in (\mathfrak{g}_\mu^*)_r$ for every $\nu \in \mathfrak{m}_r^*$ and $v \in N_r$. In this setting

$$(7) \quad \begin{aligned} L : G \times_K ((\mathfrak{m}_r^*)_r \times N_r) &\longrightarrow \mathbf{K}_{K^T}^{-1}(0)/K^T \\ [g, \nu, v]_K &\longmapsto [g, \nu + \mathbf{J}_N(v), v]_K \end{aligned}$$

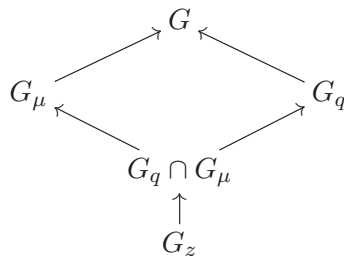
is a well-defined G -equivariant symplectomorphism between the MGS model $Y_r = G \times_K (\mathfrak{m}_r^* \times N_r)$ equipped with the symplectic form (2) and the reduced space $\mathbf{K}_{K^T}^{-1}(0)/K^T$ equipped with its canonical reduced symplectic form.

3. The symplectic slice for cotangent-lifted actions

If G acts on a manifold Q then the natural lift of the action to T^*Q will be called the *cotangent-lifted action*. The cotangent-lifted action of a proper action is again proper and it is Hamiltonian with respect to the canonical symplectic form with momentum map

$$(8) \quad \langle \mathbf{J}_{T^*Q}(q, p), \xi \rangle = \langle p, \xi_Q(q) \rangle$$

Fix $z = (q, p) \in T^*Q$ with $\mu := \mathbf{J}_{T^*Q}(z)$. By the equivariance of the cotangent bundle projection $\tau : T^*Q \rightarrow Q$ and of \mathbf{J}_{T^*Q} we have $G_z \subset G_q$ and $G_z \subset G_\mu$. Moreover, from (8) we see that μ is annihilated by any $\xi \in \mathfrak{g}_q$. Graphically, the several groups involved can be put into lattice form as



where each arrow represents an inclusion. As we shall see, one of the main problems of building a Hamiltonian tube for a cotangent-lifted action is that the relationship between the subgroups G_q , G_z and G_μ can be complicated in general. In this section we will first introduce a splitting of \mathfrak{g} that has good properties with respect to μ and G_q and then will use this splitting to restate the characterization of the symplectic slice for cotangent-lifted actions obtained in [23].

3.1. A Lie algebra decomposition

We start by giving a useful invariant splitting of \mathfrak{g} that will be the starting point of many of the constructions we will make in the next sections.

Proposition 3.1. *Let G be a Lie group, H a compact subgroup and $\mu \in \mathfrak{g}^*$ with $[\mathfrak{h}, \mathfrak{h}] \in \langle \mu \rangle^\circ$. Let Ω^μ be the bilinear form on \mathfrak{g} given by $\Omega^\mu(\xi_1, \xi_2) = -\langle \mu, [\xi_1, \xi_2] \rangle$. Then there is a H_μ -invariant splitting*

$$(9) \quad \mathfrak{g} = \mathfrak{g}_\mu \oplus \mathfrak{o} \oplus \mathfrak{l} \oplus \mathfrak{n}$$

satisfying

- 1) $\mathfrak{h} = \mathfrak{h}_\mu \oplus \mathfrak{l}$.
- 2) $\Omega^\mu|_{\mathfrak{o}}$ is non-degenerate.
- 3) \mathfrak{l} and \mathfrak{n} are Ω^μ -isotropic subspaces and $\Omega^\mu|_{\mathfrak{l} \oplus \mathfrak{n}}$ is non-degenerate.
- 4) $\Omega^\mu(\xi_1, \xi_2) = 0$ if $\xi_1 \in \mathfrak{o}$, $\xi_2 \in \mathfrak{l} \oplus \mathfrak{n}$.

Proof. As H is compact we can endow \mathfrak{g} with an Ad_H -invariant metric. The two-form Ω^μ restricted to \mathfrak{g}_μ^\perp is non-degenerate because if $\xi \in \text{Ker } \Omega^\mu|_{\mathfrak{g}_\mu^\perp}$ then $\langle \mu, [\xi, \eta] \rangle = 0$ for all $\eta \in \mathfrak{g}_\mu^\perp$ but if now $\eta \in \mathfrak{g}_\mu$ then $0 = \langle \text{ad}_\eta^* \mu, \xi \rangle = -\langle \mu, [\xi, \eta] \rangle = \langle \text{ad}_\xi^* \mu, \eta \rangle$ for any η in \mathfrak{g} but this implies that $\text{ad}_\xi^* \mu = 0$ and as $\xi \in \mathfrak{g}_\mu^\perp$ then $\xi = 0$. Denote by $\omega = \Omega^\mu|_{\mathfrak{g}_\mu^\perp}$ the restriction. The form ω is symplectic on \mathfrak{g}_μ^\perp .

Define now $\mathfrak{l} := \mathfrak{h} \cap \mathfrak{g}_\mu^\perp$ and

$$\mathfrak{o} = \{ \lambda \in \mathfrak{g}_\mu^\perp \cap \mathfrak{h}^\perp \subset \mathfrak{g} \mid \langle \text{ad}_\lambda^* \mu, \eta \rangle = 0 \quad \forall \eta \in \mathfrak{h} \}.$$

If $\xi \in \mathfrak{g}_\mu^\perp$ is ω -orthogonal to \mathfrak{l} then it must lie in $\mathfrak{o} \oplus \mathfrak{l}$ because ξ can be decomposed as $\xi = \xi_1 + \xi_2$ with $\xi_1 \in \mathfrak{h} \cap \mathfrak{g}_\mu^\perp = \mathfrak{l}$ and $\xi_2 \in \mathfrak{h}^\perp \cap \mathfrak{g}_\mu^\perp$ but then as $\langle \mu, [\xi_2, \eta] \rangle = \langle \mu, [\xi, \eta] \rangle = 0$ for any $\eta \in \mathfrak{h}$ then $\xi_2 \in \mathfrak{o}$. That is, $\mathfrak{l}^\omega \subset \mathfrak{o} \oplus \mathfrak{l}$. Conversely, if $\xi \in \mathfrak{o}$ then by definition of \mathfrak{o} $\xi \in \mathfrak{l}^\omega$, and if $\xi \in \mathfrak{h}$ for any $\eta \in \mathfrak{l}$ we have $\langle \mu, [\xi, \eta] \rangle = 0$ because $\mathfrak{l} \in \mathfrak{h}$ and $\mu \in [\mathfrak{h}, \mathfrak{h}]^\circ$ so $\xi \in \mathfrak{l}^\omega$ and, therefore, $\mathfrak{l}^\omega = \mathfrak{o} \oplus \mathfrak{l}$.

Let $\xi \in \mathfrak{o} \cap \mathfrak{o}^\omega$. Noting that $\xi \in \mathfrak{l}^\omega$ we have $\xi \in \mathfrak{o}^\omega \cap \mathfrak{l}^\omega = (\mathfrak{o} \oplus \mathfrak{l})^\omega = (\mathfrak{l}^\omega)^\omega = \mathfrak{l}$ but as $\mathfrak{o} \cap \mathfrak{l} = 0$ this implies that $\xi = 0$, thus the restriction $\omega|_{\mathfrak{o}}$ is non-degenerate.

To build the space \mathfrak{n} we will need a preliminary standard result in linear algebra.

Lemma 3.2. *Let $A, B, C \subset E$ be three linear subspaces of a linear space E such that $A \subset B$ and $A \cap C = 0$. Then*

$$B \cap (C \oplus A) = (B \cap C) \oplus A$$

Note that $\mathfrak{l} \subset \mathfrak{o}^\omega$ and as $\langle \mu, [\xi, \eta] \rangle = 0$ for any $\xi, \eta \in \mathfrak{l}$, then \mathfrak{l} is an isotropic subset of the symplectic subspace \mathfrak{o}^ω , but in fact

$$\mathfrak{l}^\omega \cap \mathfrak{o}^\omega = \mathfrak{o}^\omega \cap (\mathfrak{o} \oplus \mathfrak{l}) = (\mathfrak{o}^\omega \cap \mathfrak{o}) \oplus \mathfrak{l} = \mathfrak{l}$$

where we applied the previous lemma with $A = \mathfrak{l}$, $B = \mathfrak{o}^\omega$ and $C = \mathfrak{o}$. This implies that \mathfrak{l} is actually a Lagrangian subspace of \mathfrak{o}^ω and it is clearly H_μ -invariant. By Lemma 7.1.2 of [19] there must exist a H_μ -invariant complement $\mathfrak{n} \subset \mathfrak{g}_\mu^\perp$ of \mathfrak{l} . That is, we have $\mathfrak{g} = \mathfrak{g}_\mu \oplus \mathfrak{o} \oplus \mathfrak{l} \oplus \mathfrak{n}$ and with respect to this splitting Ω^μ block diagonalizes as

$$\Omega^\mu = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & \Omega|_{\mathfrak{o}} & 0 & 0 \\ 0 & 0 & 0 & * \\ 0 & 0 & * & 0 \end{bmatrix}$$

where the entries $*$ will not be important in our discussion. □

Remark 3.3. The subspace \mathfrak{o} was introduced in [23] by a different procedure. In that work \mathfrak{o} was constructed as a symplectic slice for the H -action on the coadjoint orbit $\mathcal{O}_\mu \subset \mathfrak{g}^*$. In fact the subspaces \mathfrak{l} and \mathfrak{n} can be understood as part of a Witt-Artin decomposition $T_\mu \mathcal{O}_\mu = \mathfrak{l} \cdot \mu \oplus \mathfrak{n} \cdot \mu \oplus \mathfrak{o} \cdot \mu$ (see [19]). Note that as vector spaces supporting a H_μ -action both \mathfrak{l} and \mathfrak{n} are isomorphic to the quotient $\mathfrak{h}/\mathfrak{h}_\mu$ and \mathfrak{o} is isomorphic to $\mathfrak{h}^{\Omega_\mu}/(\text{Ker } \Omega_\mu + \mathfrak{h})$.

3.2. The symplectic slice

The first step towards the construction of a Hamiltonian tube for cotangent-lifted actions is to describe the symplectic slice at a point $z = (q, p) \in T^*Q$. This was done in [28] under the assumption $G_q \subset G_\mu$. Later, in [23] the symplectic slice for the general case was worked out. Before stating the result we will introduce some new notation that will be used throughout the paper. Let V be a linear space supporting a linear representation of a group

G and let $a \in V$ and $b \in V^*$. The *diamond product* (see [7]) $a \diamond \alpha \in \mathfrak{g}^*$ is defined as

$$\langle a \diamond \alpha, \eta \rangle := \langle \alpha, \eta \cdot a \rangle$$

for all $\eta \in \mathfrak{g}$. With this notation the cotangent lift of a linear G -action on V has momentum map $\mathbf{J}(a, b) = a \diamond b \in \mathfrak{g}^*$. Note that if we consider $G = SO(3)$ acting on \mathbb{R}^3 the diamond product is just the usual cross product. If $\mathfrak{h} \subset \mathfrak{g}$ is a subspace then $a \diamond_{\mathfrak{h}} b := (a \diamond b)|_{\mathfrak{h}} \in \mathfrak{h}^*$. If \mathfrak{h} is the Lie algebra of a subgroup $H \subset G$, $a \diamond_{\mathfrak{h}} b$ is the momentum map for the H -action on T^*V induced by restriction of the original G -action, which is in turn the same as the lift of the restricted H -action on V . In this context, the next result gives a characterization of the symplectic slice for a cotangent-lifted action.

Theorem 3.4 (Theorem 6.1 of [23]). *Let $z = (q, p) \in T^*Q$ and S be a G_q -invariant complement of $\mathfrak{g} \cdot q$ in T_qQ . Define $H := G_q$, $\mu := \mathbf{J}(z)$, $\alpha := z|_S \in S^*$, $B := (\mathfrak{h}_\mu \cdot \alpha)^\circ \subset S$. Let \mathfrak{o} be as defined in (9)*

*The symplectic slice N at z is linearly and G_z -equivariantly symplectomorphic to the product $\mathfrak{o} \times T^*B$ equipped with the symplectic form*

$$(10) \quad \Omega_N((\lambda_1, v_1, w_1)(\lambda_2, v_2, w_2)) = -\langle \mu, [\lambda_1, \lambda_2] \rangle + \langle w_2, v_1 \rangle - \langle w_1, v_2 \rangle$$

with $\lambda_j \in \mathfrak{o}$ and $(v_j, w_j) \in B \times B^$. The corresponding momentum map for the linear G_z -action is*

$$\mathbf{J}_N(\lambda, (a, \beta)) = \frac{1}{2} \lambda \diamond_{\mathfrak{g}_z} \text{ad}_\lambda^* \mu + a \diamond_{\mathfrak{g}_z} \beta.$$

4. G -tubes

In this section we will define both simple and restricted G -tubes. These maps will be the building blocks needed to find an explicit Hamiltonian tube for cotangent-lifted actions.

4.1. Simple G -tubes

From now on we will identify TG with $G \times \mathfrak{g}$ and T^*G with $G \times \mathfrak{g}^*$ using left trivializations

$$\begin{aligned} G \times \mathfrak{g} &\longrightarrow TG & G \times \mathfrak{g}^* &\longrightarrow T^*G \\ (g, \xi) &\longmapsto T_e L_g(\xi) & (g, \nu) &\longmapsto T_e^* L_{g^{-1}}(\nu). \end{aligned}$$

Combining them we can trivialize $T(T^*G) \cong G \times \mathfrak{g} \times \mathfrak{g}^* \times \mathfrak{g}^*$.

We will need the following well-known properties of the symplectic structure and the cotangent-lifted actions of G on T^*G (see [1])

Proposition 4.1. *Let G act on itself by the left multiplication action and by cotangent lifts on T^*G then we have*

- Symplectic structure: Let $u_i := (\xi_i, \beta_i) \in T_{(g,\nu)}T^*G$ with $i = 1, 2$, the canonical one-form of T^*G is

$$(11) \quad \theta_{T^*G}(u_1) = \langle \nu, \xi_1 \rangle$$

and the canonical symplectic form $\omega_{T^*G} = -\mathbf{d}\theta_{T^*G}$ is

$$(12) \quad \omega_{T^*G}(u_1, u_2) = \langle \beta_2, \xi_1 \rangle - \langle \beta_1, \xi_2 \rangle + \langle \nu, [\xi_1, \xi_2] \rangle.$$

- Cotangent-lifted left multiplication: The G -action given by $h \cdot^L(g, \nu) = (hg, \nu)$ has as infinitesimal generator $\eta_{T^*G}^L(g, \nu) = (\text{Ad}_{g^{-1}}\eta, 0)$ and is Hamiltonian with momentum map $\mathbf{J}_L(g, \nu) = \text{Ad}_{g^{-1}}^*\nu$.
- Cotangent-lifted right multiplication: The G -action given by

$$h \cdot^R(g, \nu) = (gh^{-1}, \text{Ad}_{h^{-1}}^*\nu)$$

has as infinitesimal generator $\eta_{T^*G}^R(g, \nu) = (-\eta, -\text{ad}_\eta^*\nu)$ and is Hamiltonian with momentum map $\mathbf{J}_R(g, \nu) = -\nu$.

Notice that if we think of \mathfrak{g}^* as a manifold endowed with the G -action $g \cdot \nu = \text{Ad}_{g^{-1}}^*\nu$ then the above left and right actions of G on T^*G are exactly the actions G^L and G^R on the product manifold $G \times \mathfrak{g}^*$ (see Subsection 2.1).

We now define simple G -tubes.

Definition 4.2. Let H be a compact subgroup of G and $\mu \in \mathfrak{g}^*$. Given a splitting $\mathfrak{g} = \mathfrak{g}_\mu \oplus \mathfrak{q}$ invariant under the H_μ -action, a simple G -tube is a map

$$\Theta : G \times U \subset G \times (\mathfrak{g}_\mu^* \times \mathfrak{q}) \longrightarrow G \times \mathfrak{g}^* \cong T^*G$$

such that:

- 1) U is a connected H_μ -invariant neighborhood of 0 in $\mathfrak{g}_\mu^* \times \mathfrak{q}$.
- 2) Θ is a G^L -equivariant diffeomorphism onto $\Phi(G \times U)$ satisfying $\Theta(e, 0) = (e, \mu)$.

3) Let $u_i := (\xi_i, \dot{\nu}_i, \dot{\lambda}_i) \in T_{(g,\nu,\lambda)}G \times \mathfrak{g}_\mu^* \times \mathfrak{q}$ with $i = 1, 2$, then

$$(13) \quad (\Theta^* \omega_{T^*G})(u_1, u_2) = \langle \dot{\nu}_2, \xi_1 \rangle - \langle \dot{\nu}_1, \xi_2 \rangle + \langle \nu + \mu, [\xi_1, \xi_2] \rangle - \langle \mu, [\dot{\lambda}_1, \dot{\lambda}_2] \rangle$$

4) Θ is H_μ^T -equivariant.

5) $T_{(e,0,0)}\Theta(\xi, \dot{\nu}, \dot{\lambda}) = (\xi + \dot{\lambda}; \dot{\nu} + \text{ad}_\lambda^* \mu) \in T_{(e,0)}(T^*G)$.

If \mathfrak{q} is defined as above, note that the symplectic slice for the cotangent-lifted left multiplication of G on T^*G at $(e, \mu) \in T^*G$ is precisely \mathfrak{q} . Indeed, as $T_{(e,\mu)}\mathbf{J}_L(e, \mu) \cdot (\xi, \dot{\nu}) = -\text{ad}_\xi^* \mu + \dot{\nu}$ then a complement to $\mathfrak{g}_\mu \cdot (e, \mu)$ can be chosen to be the space $\{(\xi, \text{ad}_\xi^* \mu) \mid \xi \in \mathfrak{q}\}$, and using (12), this linear space is symplectomorphic to $(\mathfrak{q}, \Omega^\mu|_{\mathfrak{q}})$. According to Theorem 2.1, the MGS model at $(e, \mu) \in T^*G$ for the free cotangent-lifted left multiplication of G on T^*G will be of the form $G \times \mathfrak{g}_\mu^* \times \mathfrak{q}$ and, in this case, the symplectic form (2) is precisely the one given by (13). In other words, a simple G -tube is a Hamiltonian tube for T^*G at (e, μ) (properties 1–3) but we further require H_μ^T -equivariance and a prescribed property on its linearization (properties 4–5).

The next result ensures the existence of simple G -tubes. The idea is that an equivariant version of the Moser trick can be used to construct them. This part follows closely Theorem 6 in [2] and Theorem 7.3.1 in [19]. We are going to apply Moser’s trick to an explicit, well-behaved, family of symplectic potentials.

Proposition 4.3 (Existence of simple G -tubes). *Given a H_μ -invariant splitting $\mathfrak{g} = \mathfrak{g}_\mu \oplus \mathfrak{q}$ there exists a H_μ -invariant open neighborhood D of $0 \in \mathfrak{g}_\mu^* \times \mathfrak{q}$ and a simple G -tube $\Theta : G \times D \subset G \times \mathfrak{g}_\mu^* \times \mathfrak{q} \rightarrow G \times \mathfrak{g}^*$.*

Proof. As a first approximation we will consider the map

$$(14) \quad \begin{aligned} F : G \times \mathfrak{g}_\mu^* \times \mathfrak{q} &\longrightarrow G \times \mathfrak{g}^* \\ (g, \nu, \lambda) &\longmapsto (g \exp(\lambda), \text{Ad}_{\exp(\lambda)}^*(\nu + \mu)) \end{aligned}$$

defined only for λ small enough so that it is contained in the injectivity domain of the group exponential $\exp : \mathfrak{g} \rightarrow G$. The map F is G^L -equivariant and also H_μ^T -equivariant because

$$F(g'g, \nu, \lambda) = (g'g \exp(\lambda), \text{Ad}_{\exp(\lambda)}(\nu + \mu)), \quad \text{and}$$

$$\begin{aligned} F(gh^{-1}, \text{Ad}_{h^{-1}}^* \nu, \text{Ad}_h \lambda) &= (gh^{-1} \exp(\text{Ad}_h \lambda), \text{Ad}_{\exp(\text{Ad}_h \lambda)}^* (\text{Ad}_{h^{-1}}^* \nu + \mu)) \\ &= (g \exp(\lambda) h^{-1}, \text{Ad}_{h^{-1}}^* \text{Ad}_{\exp(\lambda)}^* (\nu + \mu)). \end{aligned}$$

Consider now the one-form on $G \times (\mathfrak{g}_\mu^* \times \mathfrak{q})$ given by $\theta_Y(g, \nu, \lambda)(\xi, \dot{\nu}, \dot{\lambda}) = \langle \nu + \mu, \xi \rangle + \frac{1}{2} \langle \mu, \text{ad}_\lambda \dot{\lambda} \rangle + \langle \mu, \dot{\lambda} \rangle$. It is clearly G^L -invariant and H_μ^T -invariant because

$$\begin{aligned} &\theta_Y(gh^{-1}, \text{Ad}_{h^{-1}}^* \nu, \text{Ad}_h \lambda)(\text{Ad}_h \xi, \text{Ad}_{h^{-1}}^* \dot{\nu}, \text{Ad}_h \dot{\lambda}) \\ &= \langle \text{Ad}_{h^{-1}}^* (\nu) + \mu, \text{Ad}_h \xi \rangle + \frac{1}{2} \langle \mu, [\text{Ad}_h \lambda, \text{Ad}_h \dot{\lambda}] \rangle + \langle \mu, \text{Ad}_h \dot{\lambda} \rangle \\ &= \theta_Y(g, \nu, \lambda)(\xi, \dot{\nu}, \dot{\lambda}). \end{aligned}$$

Let $u_i := (\xi_i, \dot{\nu}_i, \dot{\lambda}_i) \in T_{(g, \nu, \lambda)}(G \times \mathfrak{g}_\mu^* \times \mathfrak{q})$ with $i = 1, 2$. Note that

$$(-\mathbf{d}\theta_Y)(u_1, u_2)$$

is the right-hand side of equation (13). Consider now the family of $G^L \times H_\mu^T$ -invariant one-forms

$$\theta_t = tF^* \theta_{T^*G} + (1 - t)\theta_Y$$

and define $\omega_t := -\mathbf{d}\theta_t$. Using (12) and $T_{(e, 0, 0)}F(\xi, \dot{\nu}, \dot{\lambda}) = (\xi + \dot{\lambda}, \dot{\nu} + \text{ad}_\lambda^* \mu)$ it can be checked that

$$\begin{aligned} &(-\mathbf{d}\theta_t)(g, 0, 0)(\xi_1, \dot{\nu}_1, \dot{\lambda}_1)(\xi_2, \dot{\nu}_2, \dot{\lambda}_2) \\ &= \langle \dot{\nu}_2, \xi_1 \rangle - \langle \dot{\nu}_1, \xi_2 \rangle + \langle \mu, [\xi_1, \xi_2] \rangle - \langle \mu, [\dot{\lambda}_1, \dot{\lambda}_2] \rangle, \end{aligned}$$

but this two-form is non-degenerate because it corresponds precisely to Ω_Y of Theorem 2.1. This implies that Moser's equation $i_{X_t} \omega_t = \frac{\partial \theta_t}{\partial t}$ defines a time-dependent vector field X_t on an open set $G \times V \subset G \times \mathfrak{g}_\mu^* \times \mathfrak{q}$. If Ψ_t is the local flow of X_t then $\Psi_t^* \omega_t = \omega_0$ (see Theorem 7.3.1 of [19] for technical details). As θ_t and $-\mathbf{d}\theta_t$ are $G^L \times H_\mu^T$ invariant differential forms, then the vector field X_t is $G^L \times H_\mu^T$ invariant and, therefore, the local flow Ψ_t is $G^L \times H_\mu^T$ -equivariant for any t . Note that $\theta_Y(g, \nu, 0) = \langle \mu + \nu, \xi \rangle + \langle \mu, \dot{\lambda} \rangle$ and $F^* \theta_{T^*G}(g, \nu, 0) = \langle \mu + \nu, \xi \rangle + \langle \mu + \nu, \dot{\lambda} \rangle = \langle \mu + \nu, \xi \rangle + \langle \mu, \dot{\lambda} \rangle$. This implies that $\frac{\partial \theta_t}{\partial t} \Big|_{(g, \nu, 0)} = 0$ and $X_t(g, \nu, 0) = 0$ so $\Psi_t(g, \nu, 0) = (g, \nu, 0)$ for any $t \in \mathbb{R}$ and then there is a H_μ -invariant open set $U \subset V$ such that Ψ_1 is a diffeomorphism with domain $G \times U$.

The simple G -tube will then be the composition $\Theta = F \circ \Psi_1 : G \times U \rightarrow T^*G$. It is $G^L \times H_\mu^T$ -equivariant and it satisfies

$$\begin{aligned} \omega_Y &= \omega_0 = \Psi_1^* \omega_1 = \Psi_1^* F^* \omega_{T^*G} = \Theta^* \omega_{T^*G} \\ \text{and } \Theta(e, 0, 0) &= (e, \mu). \end{aligned}$$

Let Ψ_t be the local flow of X_t and η_t be any time-dependent tensor field then

$$(15) \quad \frac{d}{dt} \Psi_t^* \eta_t = \Psi_t^* \left(\mathcal{L}_{X_t} \eta_t + \frac{d}{dt} \eta_t \right).$$

This expression can be used to compute $T_{(e,0,0)}\Theta$. To do so let Y be any time-independent vector field on $G \times \mathfrak{g}_\mu^* \times \mathfrak{q}$ not vanishing at $(e, 0, 0)$. As X_t vanishes at $(e, 0, 0)$ then $\mathcal{L}_{X_t} Y \Big|_{(e,0,0)} = 0$. Setting $\eta_t = Y$ in (15) it gives $\frac{d}{dt} \Psi_t^* Y = 0$ but this implies $T_{(e,0,0)}\Psi_1 = \text{Id}$ and, therefore, $T_{(e,0,0)}\Theta(\xi, \dot{\nu}, \dot{\lambda}) = (\xi + \dot{\lambda}, \dot{\nu} + \text{ad}_\lambda^* \mu)$. That is, Θ satisfies all the five required conditions for a simple G -tube. \square

The main shortcoming with the previous existence result is that, as happens with Theorem 2.1, it does not produce an explicit map and it relies on the integration of a time-dependent field. However, we will see in Section 7 that in some particular cases we can explicitly describe these objects. Nevertheless, using momentum maps we can still find a simpler expression for the simple G -tube Θ . Decompose Θ as

$$\Theta(g, \nu, \lambda) = (A(g, \nu, \lambda), B(g, \nu, \lambda)) \in G \times \mathfrak{g}^*.$$

The property of G^L -equivariance implies that $A(g, \nu, \lambda) = gA(e, \nu, \lambda)$. As $\Theta(e, 0, 0) = (e, \mu)$ then $A(e, 0, 0) = e$ and $B(e, 0, 0) = \mu$.

Using Section 2.3 we have that the product $G \times \mathfrak{g}_\mu^* \times \mathfrak{q}$ is equipped with G^L and H_μ^T Hamiltonian actions with momentum maps \mathbf{K}_{G^L} and $\mathbf{K}_{H_\mu^T}$ respectively (see (6)). We also have G^L and H_μ^T Hamiltonian actions on $G \times \mathfrak{g}^*$ and their momentum maps are \mathbf{J}_{G^L} and $\mathbf{J}_{H_\mu^T}$ (see Proposition 4.1). As the difference between two momentum maps is a locally constant function and both \mathbf{J}_{G^L} and \mathbf{K}_{G^L} are equivariant then $\mathbf{J}_{G^L} \circ \Theta = \mathbf{K}_{G^L}$, that is

$$\text{Ad}_{A(g,\nu,\lambda)^{-1}}^* B(g, \nu, \lambda) = \text{Ad}_{g^{-1}}^* (\nu + \mu)$$

so

$$\begin{aligned} B(g, \nu, \lambda) &= \text{Ad}_{A(g, \nu, \lambda)}^* \text{Ad}_{g^{-1}}^*(\nu + \mu) \\ &= \text{Ad}_{g^{-1}A(g, \nu, \lambda)}^*(\nu + \mu) \\ &= \text{Ad}_{A(e, \nu, \lambda)}^*(\nu + \mu). \end{aligned}$$

If we denote $E(\nu, \lambda) = A(e, \nu, \lambda)$ then we can write

$$(16) \quad \begin{aligned} \Theta : G \times \mathfrak{g}_\mu \times \mathfrak{q} &\longrightarrow T^*G \\ (g, \nu, \lambda) &\longmapsto (gE(\nu, \lambda), \text{Ad}_{E(\nu, \lambda)}^*(\nu + \mu)). \end{aligned}$$

Therefore, a simple G -tube is determined by a function $E : U \subset \mathfrak{g}_\mu^* \times \mathfrak{q} \rightarrow G$. Rewriting Definition 4.2 in terms of the function E we could obtain necessary and sufficient conditions for E .

Remark 4.4. Note that if $\mathfrak{g} = \mathfrak{g}_\mu$, which is the hypothesis used in [28], then $\mathfrak{q} = 0$ and the shifting map

$$\begin{aligned} G \times \mathfrak{g}^* &\longrightarrow G \times \mathfrak{g}^* \\ (g, \nu) &\longmapsto (g, \nu + \mu) \end{aligned}$$

is a simple G -tube.

Remark 4.5. As Θ is H_μ^T -equivariant, the momentum preservation argument gives

$$(17) \quad \mathbf{J}_{H_\mu^T}(\Theta(g, \nu, \lambda)) = \mathbf{K}_{H_\mu^T}(g, \nu, \lambda) = -\nu|_{\mathfrak{h}_\mu} + \frac{1}{2}\lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu.$$

That is, we have the condition $(\text{Ad}_{E(\nu, \lambda)}^*(\nu + \mu))|_{\mathfrak{h}_\mu} = \nu|_{\mathfrak{h}_\mu} - \frac{1}{2}\lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu$. This property will be useful later during the proof of Proposition 6.2.

Remark 4.6. In the proof of Proposition 4.3 we noted that $\Psi_t(g, \nu, 0) = (g, \nu, 0)$ for any t , hence for any simple G -tube Θ given by the previous proposition it always satisfies $\Theta(g, \nu, 0) = (g, \nu + \mu) \in T^*G$.

4.2. Restricted G -tubes

If G acts freely on Q we will see in Section 5.4 that the simple G -tube is enough to construct explicitly the Hamiltonian tube for T^*Q but for non-free actions we will need to adapt a simple G -tube to the corresponding isotropy subgroup, the result being the restricted G -tube.

Definition 4.7. Given an adapted splitting $\mathfrak{g} = \mathfrak{g}_\mu \oplus \mathfrak{o} \oplus \mathfrak{l} \oplus \mathfrak{n}$ as in Proposition 3.1 a restricted G -tube is a map

$$\Phi : G \times U \subset G \times \mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{l}^* \longrightarrow T^*G$$

such that:

- 1) U is a connected H_μ -invariant neighborhood of 0 in $\mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{l}^*$.
- 2) Φ is a $G^L \times H_\mu^T$ -equivariant diffeomorphism between $G \times U$ and $\Phi(G \times U)$ such that $\Phi(e, 0, 0; 0) = (e, \mu)$.
- 3) Let $u_i := (\xi_i, \dot{\nu}_i, \dot{\lambda}_i, \dot{\epsilon}_i) \in T_{(g, \nu, \lambda, \epsilon)} G \times \mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{l}^*$ with $i = 1, 2$, then $\Phi^* \omega_{T^*G}$ is

$$(18) \quad \omega_{\text{restr}}(u_1, u_2) = \langle \dot{\nu}_2, \xi_1 \rangle - \langle \dot{\nu}_1, \xi_2 \rangle + \langle \nu + \mu, [\xi_1, \xi_2] \rangle - \langle \mu, [\dot{\lambda}_1, \dot{\lambda}_2] \rangle.$$

- 4) $\mathbf{J}_R(\Phi(g, \nu, \lambda, \epsilon)) \Big|_{\mathfrak{l}} = -\epsilon$ for any $(g, \nu, \lambda, \epsilon)$ where \mathbf{J}_R is the momentum map for the G^R -action on T^*G (see Proposition 4.1).

If we are given a simple G -tube Θ then we can build a restricted G -tube Φ solving a non-linear equation. In fact, the restricted G -tube will be of the form $\Phi(g, \nu, \lambda, \epsilon) = \Theta(g, \nu, \lambda + \zeta(\nu, \lambda, \epsilon))$ for some map $\zeta : \mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{l}^* \rightarrow \mathfrak{n}$. This is the main idea behind the following result.

Proposition 4.8 (Existence of restricted G -tubes). *Given an adapted splitting $\mathfrak{g} = \mathfrak{g}_\mu \oplus \mathfrak{o} \oplus \mathfrak{l} \oplus \mathfrak{n}$ as in Proposition 3.1 there is an H_μ -invariant open neighborhood D of $0 \in \mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{l}^*$ and a restricted G -tube $\Phi : G \times D \rightarrow T^*G$.*

Proof. Define $\mathfrak{q} = \mathfrak{o} \oplus \mathfrak{l} \oplus \mathfrak{n}$. Using Proposition 4.3 there exists a simple G -tube Θ defined on the symplectic space $Y := G \times U \subset G \times (\mathfrak{g}_\mu^* \times \mathfrak{q})$ with symplectic form ω_Y (13). As $U \subset \mathfrak{g}_\mu^* \times \mathfrak{q}$ is a neighborhood of 0 there are H_μ -invariant neighborhoods of the origin $(\mathfrak{g}_\mu^*)_r \subset \mathfrak{g}_\mu^*$, $\mathfrak{o}_r \subset \mathfrak{o}$ and $\mathfrak{n}_r \subset \mathfrak{n}$ such that $(\mathfrak{g}_\mu^*)_r \times (\mathfrak{o}_r + \mathfrak{n}_r) \subset U$. Consider now the map

$$\begin{aligned} \iota_W : W = G \times ((\mathfrak{g}_\mu^*)_r \times \mathfrak{o}_r \times \mathfrak{n}_r) &\longrightarrow Y = G \times U \subset G \times (\mathfrak{g}_\mu^* \times \mathfrak{q}) \\ (g, \nu, \lambda, \zeta) &\longmapsto (g, \nu, \lambda + \zeta) \end{aligned}$$

This map is a $G^L \times H_\mu^T$ -equivariant embedding. By the properties of the adapted splitting (see Proposition 3.1) $\Omega^\mu(\lambda, \zeta) = 0$ if $\lambda \in \mathfrak{o}$ and $\zeta \in \mathfrak{n}$, therefore,

$$(\iota_{W^*}^* \omega_Y)(u_1, u_2) = \langle \dot{\nu}_2, \xi_1 \rangle - \langle \dot{\nu}_1, \xi_2 \rangle + \langle \nu + \mu, [\xi_1, \xi_2] \rangle - \langle \mu, [\dot{\lambda}_1, \dot{\lambda}_2] \rangle$$

where $u_i := (\xi_i, \dot{\nu}_i, \dot{\lambda}_i, \dot{\zeta}_i) \in T_{(g, \nu, \lambda, \zeta)} G \times \mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{n}$ with $i = 1, 2$. In order to obtain the restricted G -tube we will need to impose the relationship between ε and \mathbf{J}_R . To do so define the map

$$\begin{aligned} \psi : W &\longrightarrow G \times \mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{l}^* \\ (g, \nu, \lambda, \zeta) &\longmapsto (g, \nu, \lambda; -\mathbf{J}_R(\Theta(g, \nu, \lambda + \zeta))|_{\mathfrak{l}}). \end{aligned}$$

Note that this map is $G^L \times H_\mu^T$ -equivariant because

$$\begin{aligned} &(g'g, \nu, \lambda; -\mathbf{J}_R(\Theta(g'g, \nu, \lambda + \zeta))|_{\mathfrak{l}}) \\ &= (g'g, \nu, \lambda; -\mathbf{J}_R(g'\Theta(g, \nu, \lambda + \zeta))|_{\mathfrak{l}}) \\ &= (g'g, \nu, \lambda; -\mathbf{J}_R(\Theta(g, \nu, \lambda + \zeta))|_{\mathfrak{l}}) \end{aligned}$$

and

$$\begin{aligned} \psi(h \cdot^T (g, \nu, \lambda, \zeta)) &= (gh^{-1}, \text{Ad}_{h^{-1}}^* \nu, \text{Ad}_h \lambda; -\mathbf{J}_R(\Theta(h \cdot^T (g, \nu, \lambda + \zeta))|_{\mathfrak{l}})) \\ &= (gh^{-1}, \text{Ad}_{h^{-1}}^* \nu, \text{Ad}_h \lambda; -\mathbf{J}_R(h \cdot^T \Theta((g, \nu, \lambda + \zeta))|_{\mathfrak{l}})) \\ &= (gh^{-1}, \text{Ad}_{h^{-1}}^* \nu, \text{Ad}_h \lambda; -\text{Ad}_{h^{-1}}^* \mathbf{J}_R(\Theta((g, \nu, \lambda + \zeta))|_{\mathfrak{l}})) \\ &= h \cdot^T (g, \nu, \lambda; -\mathbf{J}_R(\Theta(g, \nu, \lambda + \zeta))|_{\mathfrak{l}}). \end{aligned}$$

Moreover, if we endow $G \times \mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{l}^*$ with the two-form (18) then $\psi^* \omega_{\text{restr}} = \iota_{W^*}^* \omega_Y$. We will now check that Ψ is invertible. Let

$$v := (\xi, \dot{\nu}, \dot{\lambda}, \dot{\zeta}) \in T_{(e, 0, 0, 0)} G \times \mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{n},$$

then

$$\begin{aligned} T_{(e, 0, 0, 0)} \psi \cdot v &= \left(\xi, \dot{\nu}, \dot{\lambda}; -T_{(e, 0, 0)} (\mathbf{J}_R|_{\mathfrak{l}} \circ \Theta) \cdot (\xi, \dot{\nu}, \dot{\lambda} + \dot{\zeta}) \right) \\ &= \left(\xi, \dot{\nu}, \dot{\lambda}; -T_{(e, 0)} (\mathbf{J}_R|_{\mathfrak{l}}) \cdot (\xi + \dot{\lambda} + \dot{\zeta}, \dot{\nu} + \text{ad}_{\dot{\lambda} + \dot{\zeta}}^* \mu) \right) \\ &= \left(\xi, \dot{\nu}, \dot{\lambda}; -(\dot{\nu} + \text{ad}_{\dot{\lambda} + \dot{\zeta}}^* \mu)|_{\mathfrak{l}} \right) \\ &= \left(\xi, \dot{\nu}, \dot{\lambda}; -(\text{ad}_{\dot{\zeta}}^* \mu)|_{\mathfrak{l}} \right) \end{aligned}$$

where we have used the expression for $T_{(e,0,0)}\Theta$ given in Definition 4.2 and that $\text{ad}_\lambda^*\mu|_{\mathfrak{l}} = 0$ since \mathfrak{o} and \mathfrak{l} are Ω^μ -orthogonal (see Proposition 3.1).

Since \mathfrak{l} and \mathfrak{n} are complementary isotropic subspaces, the map $\sigma : \mathfrak{n} \rightarrow \mathfrak{l}^*$ given by $\sigma(\zeta) = \text{ad}_\zeta^*\mu|_{\mathfrak{l}}$ is a linear H_μ -equivariant isomorphism. This implies that $T_{(e,0,0,0)}\psi$ is invertible. By the Inverse Function Theorem there is a neighborhood of $(e, 0, 0, 0) \in G \times \mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{l}^*$ on which ψ^{-1} is well defined. Due to $G^L \times H_\mu^T$ equivariance of Ψ this neighborhood must be of the form $G \times D$ with $D \subset \mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{l}^*$ a H_μ -invariant neighborhood of zero.

Note that the composition $\Theta \circ \iota_W \circ \psi^{-1}$ is a restricted G -tube because it satisfies

$$(\Theta \circ \iota_W \circ \psi^{-1})^* \omega_{T^*G} = (\iota_W \circ \psi^{-1})^* \omega_Y = \omega_{\text{restr}},$$

it is $G^L \times H_\mu^T$ -equivariant (because it is the composition of $G^L \times H_\mu^T$ -equivariant maps), the origin $(e, 0, 0, 0)$ is mapped to $(e, \mu) \in T^*G$, and it is a diffeomorphism onto its image (because it is a composition of diffeomorphisms onto its images). Finally, if $(g, \nu, \lambda, \varepsilon) = \psi(g, \nu, \lambda, \zeta)$ then

$$\mathbf{J}_R(\Theta(g, \nu, \lambda + \zeta))|_{\mathfrak{l}} = -\varepsilon,$$

that is $(\mathbf{J}_R|_{\mathfrak{l}} \circ \Theta \circ \iota_W \circ \psi^{-1})(g, \nu, \lambda, \varepsilon) = -\varepsilon$, which is the condition needed for a restricted G -tube.

To sum up, the composition $\Phi := \Theta \circ \iota_W \circ \psi^{-1} : G \times D \rightarrow T^*G$ is a restricted G -tube. This map can also be written as

$$(19) \quad \begin{aligned} \Phi : G \times D \subset G \times \mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{l}^* &\longrightarrow T^*G \\ (g, \nu, \lambda; \varepsilon) &\longmapsto \Theta(g, \nu, \lambda + \zeta(\nu, \lambda; \varepsilon)) \end{aligned}$$

where $\zeta : D \subset \mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{l}^* \rightarrow \mathfrak{n}$ is determined by the equation $\mathbf{J}_R|_{\mathfrak{l}}(\Phi(g, \nu, \lambda + \zeta, \varepsilon)) = -\varepsilon$. In fact, using Remark 4.6 and this characterization of Φ it follows that

$$(20) \quad \Phi(g, \nu, 0, 0) = (g, \nu + \mu) \in T^*G.$$

□

5. Cotangent bundle Hamiltonian tubes

Let G be a Lie group acting properly on Q , and fix $z \in T^*Q$. In this section we will construct a Hamiltonian tube for the cotangent-lifted action of G on

T^*Q around z that will be explicit except for the computation of a restricted G -tube. This Hamiltonian tube will be a generalization of the construction in [28] under the hypothesis $G_\mu = G$.

5.1. Cotangent-lifted Palais model

We will first reduce the problem on T^*Q to a problem on $T^*(G \times_H S)$. This first simplification is already discussed in [28]. Recall the well-known regular reduction theorem for cotangent bundles at zero momentum [26].

Theorem 5.1 (Regular cotangent reduction at zero). *Let G act freely and properly by cotangent lifts on T^*Q with momentum map \mathbf{J} . Denote by π_G the projection $Q \rightarrow Q/G$ and consider the map $\varphi : \mathbf{J}^{-1}(0) \rightarrow T^*(Q/G)$ defined by $\langle \varphi(z), T\pi_G(v) \rangle = \langle z, v \rangle$ for every $z \in T_q^*Q$ and $v \in T_qQ$. Let π_0 and ι by the natural projection $\pi_0 : \mathbf{J}^{-1}(0) \rightarrow \mathbf{J}^{-1}(0)/G$ and the inclusion $\iota : \mathbf{J}^{-1}(0) \rightarrow T^*Q$. The map φ is a G -invariant surjective submersion that induces a symplectomorphism*

$$\bar{\varphi} : \mathbf{J}^{-1}(0)/G \longrightarrow T^*(Q/G)$$

where $\mathbf{J}^{-1}(0)/G$ is endowed with the unique symplectic form ω_0 satisfying and $\pi_0^*\omega_0 = \iota^*\omega_{T^*Q}$.

Let G be a Lie group and H a compact subgroup that acts linearly on the linear space S . Let S_r be a H -invariant open neighborhood of $0 \in S$. We will consider the symplectic space $T^*(G \times S_r)$ that can be identified with $G \times \mathfrak{g}^* \times S_r \times S^*$ using the left-trivialization of G and the linear structure of S . In Section 2.1 we introduced the G^L and H^T actions on the space $G \times S_r$. These actions can be lifted to Hamiltonian actions on $T^*(G \times S_r)$. More explicitly, using Proposition 4.1 and the diamond notation we have

- cotangent-lifted G^L -action: $g' \cdot^L (g, \nu, a, b) = (g'g, \nu, a, b)$ with momentum map

$$\mathbf{J}_{G^L}(g, \nu, a, b) = \text{Ad}_{g^{-1}}^* \nu.$$

- cotangent-lifted H^T -action: $h \cdot^T (g, \nu, a, b) = (gh^{-1}, \text{Ad}_{h^{-1}}^* \nu, h \cdot a, h \cdot b)$ with momentum map

$$\mathbf{J}_{H^T}(g, \nu, a, b) = -\nu|_{\mathfrak{h}} + a \diamond b.$$

Theorem 5.1 applied to $G \times S_r$ with the H^T -action gives the diagram

$$(21) \quad \begin{array}{ccc} \mathbf{J}_{H^T}^{-1}(0) & \hookrightarrow & T^*(G \times S_r) \\ \downarrow & \searrow \varphi & \\ \mathbf{J}_{H^T}^{-1}(0)/H^T & \xrightarrow{\bar{\varphi}} & T^*(G \times_H S_r) \end{array}$$

but, since $\mathbf{J}_{H^T}^{-1}(0)$ is G^L -invariant the quotient $\mathbf{J}_{H^T}^{-1}(0)/H^T$ supports a G -action that it can be checked to be Hamiltonian with momentum map $\mathbf{J}_{\text{red}}([g, \nu, a, b]_{H^T}) = \text{Ad}_{g^{-1}}^* \nu$. That is, $\bar{\varphi}$ is in fact a G -equivariant symplectomorphism and φ is a G^L -equivariant surjective submersion.

Consider now a Lie group G acting properly on a general manifold Q and by cotangent lifts on T^*Q with momentum map \mathbf{J}_{T^*Q} . Our goal is to construct a Hamiltonian tube around an arbitrary point $z \in T^*Q$. For that, let $q = \tau(z)$ where $\tau : T^*Q \rightarrow Q$ is the projection, define $H := G_q$ and consider a linear slice at q , that is a H -invariant complement S to $\mathfrak{g} \cdot q$ in $T_q Q$. Using Palais' model (1) there is a H -invariant neighborhood $S_r \subset S$ and a G -equivariant diffeomorphism $\mathbf{s} : G \times_H S_r \rightarrow U \subset T^*Q$ satisfying $\mathbf{s}([e, 0]_H) = q$. As \mathbf{s} is a diffeomorphism the cotangent lift $T^*\mathbf{s}^{-1} : T^*(G \times_H S_r) \rightarrow \pi^{-1}(U) \subset Q$ is a G -equivariant symplectomorphism onto $T^*U \cong \tau^{-1}(U) \subset T^*Q$.

If we denote $\alpha = z|_S$ and $\mu = \mathbf{J}_{T^*Q}(z)$, then $T^*\mathbf{s}^{-1}(\varphi(e, \mu, 0, \alpha)) = z$ since $\tau(\varphi(e, \mu, 0, \alpha)) = [e, 0]_H = \mathbf{s}^{-1}(q)$, and as any $v \in T_q Q$ can be decomposed as $v = \xi \cdot q + \dot{a}$ with $\xi \in \mathfrak{g}$ and $\dot{a} \in S$ then

$$\begin{aligned} \langle T_q^* \mathbf{s}^{-1}(\varphi(e, \mu, 0, \alpha)), v \rangle &= \langle \varphi(e, \mu, 0, \alpha), T_q \mathbf{s}^{-1} \cdot v \rangle = \langle \varphi(e, \mu, 0, \alpha), (\xi, \dot{a}) \rangle \\ &= \langle (\mu, \alpha), (\xi, \dot{a}) \rangle = \langle \mu, \xi \rangle + \langle \alpha, \dot{a} \rangle \\ &= \langle \mathbf{J}_{T^*Q}(z), \xi \rangle + \langle z, \dot{a} \rangle \\ &= \langle z, \xi \cdot q \rangle + \langle z, \dot{a} \rangle = \langle z, v \rangle. \end{aligned}$$

Therefore, from now on we will assume without loss of generality $Q = G \times_H S_r$ and $z = \varphi([e, \mu, 0, \alpha]_H)$ with $\mu \in \mathfrak{g}^*$ and $\alpha \in S^*$. Moreover, this simplification is explicit up to the exponential of a metric.

The G -isotropy of z is $G_z = H_\alpha \cap G_\mu$ since $G_z = G_{[e, \mu, 0, \alpha]_H} = H_{(\mu, \alpha)} = H_\mu \cap H_\alpha = G_\mu \cap H_\alpha$. If \mathfrak{p} is a H_μ -invariant complement of \mathfrak{h}_μ in \mathfrak{g}_μ , \mathfrak{s} is a complement of \mathfrak{g}_z in \mathfrak{h}_μ and $B = (\mathfrak{h}_\mu \cdot \alpha)^\circ \subset S$ then, using Theorems 3.4 and 2.1, and the adapted splitting of Proposition 3.1, the Hamiltonian tube

around z has to be a map of the form

$$(22) \quad \mathcal{T} : G \times_{G_z} \left(\underbrace{(\mathfrak{g}^* \oplus \mathfrak{p}^*)}_{\mathfrak{m}^*} \times \underbrace{\mathfrak{o} \times B \times B^*}_N \right) \longrightarrow T^*(G \times_H S).$$

The first difficulty that we find is that the MGS model is a G_z -quotient but the target space is a H -quotient. For this reason, instead of constructing directly the tube we are going to split it as the composition of two maps: one that goes from an H_μ -quotient to an H -quotient (essentially this is done by a restricted G -tube) and another that goes from a G_z -quotient to an H_μ -quotient. We will carefully explain this process in the following sections.

5.2. The $\alpha = 0$ case

In this section we will construct a Hamiltonian around a point of the form $z_0 = \varphi(e, \mu, 0, 0) \in T^*(G \times_H S)$ which is explicit up to a restricted G -tube. Notice that $G_{z_0} = G_\mu \cap H_\alpha = H_\mu$ and by Theorem 3.4 and the adapted splitting of Proposition 3.1 the symplectic slice is $N_0 = \mathfrak{o} \times S \times S^*$ with symplectic form (10). Then the map (22) reduces in this case to

$$\mathcal{T}_0 : G \times_{H_\mu} \left(\underbrace{\mathfrak{p}^*}_{\mathfrak{m}^*} \times \underbrace{\mathfrak{o} \times S \times S^*}_{N_0} \right) \longrightarrow T^*(G \times_H S)$$

where $G \times_{H_\mu} (\mathfrak{p}^* \times \mathfrak{o} \times S \times S^*)$ is equipped with the symplectic form (2). As the general Hamiltonian tube will eventually factor through \mathcal{T}_0 we will need to ensure that its domain is large enough, and in particular it should contain all the points of the form $[e, 0, 0, 0, b]_{H_\mu}$ for all $b \in S^*$.

Theorem 5.2. *Consider the point $z_0 = \varphi(e, \mu, 0, 0) \in T^*(G \times_H S)$. Let $\mathfrak{g} = \mathfrak{g}_\mu \oplus \mathfrak{o} \oplus \mathfrak{l} \oplus \mathfrak{n}$ be an adapted splitting in the sense of Proposition 3.1 and let $\Phi : G \times U_\Phi \longrightarrow T^*G$ be an associated restricted G -tube. In this context, there are H_μ -invariant open neighborhoods of zero: $\mathfrak{p}_r^* \subset \mathfrak{p}^*$, $\mathfrak{o}_r \subset \mathfrak{o}$ and an H -invariant open neighborhood of zero $\mathfrak{h}_r^* \subset \mathfrak{h}^*$ such that the map*

$$(23) \quad \mathcal{T}_0 : G \times_{H_\mu} (\mathfrak{p}_r^* \times \mathfrak{o}_r \times (T^*S)_r) \longrightarrow T^*(G \times_H S)$$

$$[g, \nu, \lambda; a, b]_{H_\mu} \longmapsto \varphi(\Phi(g, \tilde{\nu}, \lambda; a \diamond_{\mathfrak{l}} b); a, b)$$

is a Hamiltonian tube around the point z_0 , where

$$\tilde{\nu} = \nu + \underbrace{\frac{1}{2} \lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu + a \diamond_{\mathfrak{h}_\mu} b}_{\mathbf{J}_{N_0}(\lambda, a, b)}$$

and $(T^*S)_r := \{(a, b) \in T^*S \mid a \diamond_{\mathfrak{h}} b \in \mathfrak{h}_r^*\}$.

Proof. If we assume the existence of $(\mathfrak{p}^*)_r, \mathfrak{o}_r$ and \mathfrak{h}_r such that the map \mathcal{T}_0 is well defined then it follows from the properties of Φ that

$$\mathcal{T}_0([e, 0, 0; 0, 0]_{H_\mu}) = \varphi(\Phi(e, 0, 0; 0); 0, 0) = \varphi(e, \mu; 0, 0)$$

and by the G -equivariance of Φ it is also clear that

$$\begin{aligned} \mathcal{T}_0(g' \cdot [g, \nu, \lambda; a, b]_{H_\mu}) &= \mathcal{T}_0([g'g, \nu, \lambda; a, b]_{H_\mu}) \\ &= \varphi(\Phi(g'g, \tilde{\nu}, \lambda; a \diamond_{\mathfrak{l}} b); a, b) \\ &= \varphi(g' \cdot \Phi(g, \tilde{\nu}, \lambda; a \diamond_{\mathfrak{l}} b); a, b) \\ &= g' \cdot \varphi(\Phi(g, \tilde{\nu}, \lambda; a \diamond_{\mathfrak{l}} b); a, b) \\ &= g' \cdot \mathcal{T}_0([g, \nu, \lambda; a, b]_{H_\mu}). \end{aligned}$$

We will divide the rest of the proof in three steps. In the first one we prove that there is a set $G \times_{H_\mu} (\mathfrak{p}_{\text{dom}}^* \times \mathfrak{o}_{\text{dom}} \times (T^*S)_{\text{dom}})$ such that the map \mathcal{T}_0 is well defined, it pulls-back the natural symplectic form of $T^*(G \times_H S)$ to the MGS form $G \times_{H_\mu} (\mathfrak{p}_{\text{dom}}^* \times \mathfrak{o}_{\text{dom}} \times (T^*S)_{\text{dom}})$ and it is a local diffeomorphism. In the second one we will show that it is injective in a certain subset and in the third we will prove that it is a diffeomorphism onto its image.

1- \mathcal{T}_0 is a local symplectomorphism. Let $N_0 = \mathfrak{o} \times S \times S^*$ be the symplectic slice at $z_0 = \varphi(e, \mu, 0, 0)$. As in Section 2.3, there must be an H_μ -invariant neighborhood $(\mathfrak{g}_\mu^*)_r$ such that the product $Z := G \times (\mathfrak{g}_\mu^*)_r \times (\mathfrak{o} \times S \times S^*)$ with $\omega_Z := \omega_{T_\mu} + \Omega_{N_0}$ is a symplectic manifold with G^L and H_μ^T Hamiltonian actions with momentum maps \mathbf{K}_{G^L} and $\mathbf{K}_{H_\mu^T}$ (see (6)).

We will now use the restricted G -tube (see Definition 4.7) $\Phi : G \times U_\Phi \subset G \times \mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{l}^* \rightarrow T^*G$ to relate Z with $T^*(G \times S)$. As Φ is only defined on $G \times U_\Phi$ we will define the open set

$$D := \{(\nu, \lambda, a, b) \mid (\nu, \lambda, a \diamond_{\mathfrak{l}} b) \in U_\Phi, \nu \in (\mathfrak{g}_\mu^*)_r\} \subset \mathfrak{g}_\mu^* \times \mathfrak{o} \times S \times S^*$$

and the map

$$(24) \quad \begin{aligned} f : G \times D &\longrightarrow T^*G \times T^*S \\ (g, \nu, \lambda, a, b) &\longmapsto (\Phi(g, \nu, \lambda, a \diamond_{\mathfrak{l}} b), a, b). \end{aligned}$$

The pullback of $\omega_{T^*(G \times S)}$ by f is ω_Z , because

$$\begin{aligned} & (f^* \omega_{T^*(G \times S)})(u_1, u_2) \\ &= \langle \Phi^* \omega_{T^*G}(g, \nu, \lambda, a \diamond_l b)(v_1, v_2) + \omega_{T^*S}(a, b)(w_1, w_2) \rangle \\ &= \underbrace{\langle \dot{\nu}_2, \xi_1 \rangle - \langle \dot{\nu}_1, \xi_2 \rangle + \langle \nu + \mu, [\xi_1, \xi_2] \rangle}_{\omega_{T_\mu}} \\ &\quad - \underbrace{\langle \mu, [\dot{\lambda}_1, \dot{\lambda}_2] \rangle + \langle \dot{b}_2, \dot{a}_1 \rangle - \langle b_1, \dot{b}_1 \rangle}_{\Omega_{N_0}} \end{aligned}$$

where $u_i = (\xi_i, \dot{\nu}_i, \dot{\lambda}_i, \dot{a}_i, \dot{b}_i) \in T_{(g, \nu, \lambda, a, b)}(G \times D)$.

Note that on $G \times D$ there is a $G^L \times H_\mu^T$ action but on $T^*G \times T^*S \cong G \times \mathfrak{g}^* \times S \times S^*$ there is a $G^L \times H^T$ action. As the map f is $G^L \times H_\mu^T$ -equivariant it preserves the H_μ -momentum, that is, $\mathbf{K}_{H_\mu^T} = \mathbf{J}_{H_\mu^T} \circ f$. In particular $f(\mathbf{K}_{H_\mu^T}^{-1}(0)) \subset \mathbf{J}_{H_\mu^T}^{-1}(0)$. However, the \mathfrak{l} -momentum property (see Definition 4.7) of restricted G -tubes allows us to improve this since for any $\xi \in \mathfrak{l}$

$$\begin{aligned} \langle \mathbf{J}_{H^T}(f(g, \nu, \lambda; a, b)), \xi \rangle &= \langle \mathbf{J}_R(\Phi(g, \nu, \lambda; a \diamond_l b)) + a \diamond_{\mathfrak{h}} b, \xi \rangle \\ &= \langle \mathbf{J}_R(\Phi(g, \nu, \lambda; a \diamond_l b)) \Big|_{\mathfrak{l}} + a \diamond_l b, \xi \rangle \\ &= \langle -a \diamond_l b + a \diamond_l b, \xi \rangle = 0. \end{aligned}$$

This means that f can be restricted to a map

$$\tilde{f} : \mathbf{K}_{H_\mu^T}^{-1}(0) \longrightarrow \mathbf{J}_{H^T}^{-1}(0)$$

and this is the key condition that will allow us to relate the H_μ -quotient $G \times_{H_\mu} (\mathfrak{p}^* \times N_0)$ with the H -quotient $\mathbf{J}_{H^T}^{-1}(0)/H^T \cong T^*(G \times_H S)$. To do so consider the diagram

$$\begin{array}{ccccc} & & G \times D & \xrightarrow{f} & T^*(G \times S) \\ & & \uparrow & & \uparrow \\ G \times \mathfrak{p}_{\text{dom}}^* \times \mathfrak{o}_{\text{dom}} \times (T^*S)_{\text{dom}} & \xrightarrow{l} & \mathbf{K}_{H_\mu^T}^{-1}(0) & \xrightarrow{\tilde{f}} & \mathbf{J}_{H^T}^{-1}(0) \\ & \downarrow & \downarrow & \searrow & \downarrow \varphi \\ G \times_{H_\mu} (\mathfrak{p}_{\text{dom}}^* \times \mathfrak{o}_{\text{dom}} \times (T^*S)_{\text{dom}}) & \xrightarrow{L} & \mathbf{K}_{H_\mu^T}^{-1}(0)/H_\mu^T & \xrightarrow{F} & \mathbf{J}_{H^T}^{-1}(0)/H^T \xrightarrow{\bar{\varphi}} T^*(G \times_H S) \\ & & \searrow & \swarrow & \uparrow \tau_0 \end{array}$$

Composing \tilde{f} with the projection by H^T in the target we get a smooth map $\mathbf{K}_{H_\mu^T}^{-1}(0) \longrightarrow \mathbf{J}_{H^T}^{-1}(0)/H^T$ which is G^K -equivariant and H_μ^T -invariant

and so it induces the smooth mapping

$$F : \mathbf{K}_{H_\mu^T}^{-1}(0)/H_\mu^T \longrightarrow \mathbf{J}_{H^T}^{-1}(0)/H^T.$$

If $\mathbf{K}_{H_\mu^T}^{-1}(0)/H_\mu^T$ is endowed with the reduced form $(\omega_Z)_{\text{red}}$ and $\mathbf{J}_{H^T}^{-1}(0)/H^T$ with $(\omega_{T^*(G \times S)})_{\text{red}}$ then $F^*(\omega_{T^*(G \times S)})_{\text{red}} = (\omega_Z)_{\text{red}}$ because $f^*\omega_{T^*(G \times S)} = \omega_Z$. In particular F is an immersion. Also, as the H_μ^T -action on $G \times D$ is free

$$\begin{aligned} \dim \mathbf{K}_{H_\mu^T}^{-1}(0)/H_\mu^T &= \dim \mathbf{K}_{H_\mu^T}^{-1}(0) - \dim \mathfrak{h}_\mu \\ &= \dim(G \times (\mathfrak{g}_\mu^* \times \mathfrak{o})) + 2 \dim S - 2 \dim \mathfrak{h}_\mu \\ &= \dim \mathfrak{g} + \dim \mathfrak{g}_\mu + \dim \mathfrak{o} - 2 \dim \mathfrak{h}_\mu + 2 \dim S \\ &= 2 \dim \mathfrak{p} + 2 \dim \mathfrak{o} + \dim \mathfrak{l} + \dim \mathfrak{n} + 2 \dim S \\ &= 2 \dim \mathfrak{p} + 2 \dim \mathfrak{o} + 2 \dim \mathfrak{l} + 2 \dim S \\ &= 2(\dim \mathfrak{g} - \dim \mathfrak{h} + \dim S). \end{aligned}$$

Analogously,

$$\dim \mathbf{J}_H^{-1}(0)/H^T = \dim \mathbf{J}_H^{-1}(0) - \dim \mathfrak{h} = 2(\dim \mathfrak{g} - \dim \mathfrak{h} + \dim S).$$

This implies that F is a local diffeomorphism because it is an immersion between spaces of the same dimension.

By continuity we can choose H_μ -invariant neighborhoods of the origin $\mathfrak{p}_{\text{dom}}^* \subset \mathfrak{p}^*$, $\mathfrak{o}_{\text{dom}} \subset \mathfrak{o}$ and an H -invariant neighborhood of the origin $\mathfrak{h}_{\text{dom}}^* \subset \mathfrak{h}^*$ such that $(\nu + \underbrace{\frac{1}{2} \lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu + a \diamond_{\mathfrak{h}_\mu} b, \lambda, a, b)}_{\mathbf{J}_{N_0}}) \in D$ for any $\nu \in \mathfrak{p}_{\text{dom}}^*$, $\lambda \in \mathfrak{o}_{\text{dom}}$ and $a, b \in T^*S$ with $a \diamond_{\mathfrak{h}} b \in \mathfrak{h}_{\text{dom}}^*$. The map

$$L : G \times_{H_\mu} (\mathfrak{p}_{\text{dom}}^* \times \mathfrak{o}_{\text{dom}} \times (T^*S)_{\text{dom}}) \rightarrow \mathbf{K}_{H_\mu^T}^{-1}(0)/H^T$$

given by $[g, \nu, \lambda, a, b]_{H_\mu} \mapsto [g, \nu + \mathbf{J}_{N_0}(\lambda, a, b), \lambda, a, b]_{H_\mu}$ is well defined and, as in (7), $L^*(\omega_Z)_{\text{red}} = \Omega_Y$. The conclusion of this first step is that the composition $\mathcal{T}_0 := \bar{\varphi} \circ F \circ L$ is then a local diffeomorphism that pulls-back the canonical form of $T^*(G \times_H S)$ to the MGS form on the set $G \times_{H_\mu} (\mathfrak{p}_{\text{dom}}^* \times \mathfrak{o}_{\text{dom}} \times (T^*S)_{\text{dom}})$.

2- \mathcal{T}_0 is locally injective. As $\mathcal{T}_0 : G \times_{H_\mu} (\mathfrak{p}_{\text{dom}}^* \times \mathfrak{o}_{\text{dom}} \times (T^*S)_{\text{dom}}) \rightarrow T^*(G \times_H S)$ is a local diffeomorphism, there is a neighborhood of $[e, 0, 0, 0, 0]_{H_\mu}$ such that \mathcal{T}_0 is injective on it. Using that \mathcal{T}_0 is G -equivariant and that

the action is proper this neighborhood can be chosen to be G -invariant (see for example the proof of Theorem 2.3.28 in [19]). That is, \mathcal{T}_0 will be injective when restricted to the set $G \times_{H_\mu} \left(\mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_{\text{inj}} \right)$ where $\mathfrak{p}_{\text{inj}}^* \subset \mathfrak{p}_{\text{dom}}^*$, $\mathfrak{o}_{\text{inj}} \subset \mathfrak{o}_{\text{dom}}$ are H_μ -invariant neighborhoods and $(T^*S)_{\text{inj}}$ is an H -invariant neighborhood of 0 on $(T^*S)_{\text{dom}}$. Note that we cannot ensure that $(T^*S)_{\text{inj}}$ will be big enough to contain all the points of the form $(0, b) \in T^*S$. This issue will be addressed in the next step.

3- \mathcal{T}_0 is injective. In this step, we will define an open set $(T^*S)_r \subset (T^*S)_{\text{dom}}$ such that the restriction

$$\mathcal{T}_0 : G \times_{H_\mu} \left(\mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r \right) \longrightarrow \mathcal{T}_0 \left(G \times_{H_\mu} \left(\mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r \right) \right)$$

is a proper map

The key result that we will use to prove the properness of \mathcal{T}_0 is the following topological result.

Proposition 5.3 (Lemma 5 of [16]). *Let H be a compact Lie group acting on a symplectic vector space (W, ω_W) denote by $\mathbf{J} : W \rightarrow \mathfrak{h}^*$ the associated homogeneous momentum map (5). Then \mathbf{J} is H -open relative to its image. That is, if U is an H -invariant open set of W then $\mathbf{J}(U)$ is an H -invariant open set of the topological space $\mathbf{J}(W) \subset \mathfrak{h}^*$.*

Let $U_1 \subset S$ and $U_2 \subset S^*$ be H -invariant neighborhoods of the origin such that $\overline{U_1 \times U_2} \subset (T^*S)_{\text{inj}}$. Using Proposition 5.3 there is \mathfrak{h}_r^* an open neighborhood of $0 \in \mathfrak{h}^*$ such that

$$\mathfrak{h}_r^* \cap (S \diamond_{\mathfrak{h}} S^*) = U_1 \diamond_{\mathfrak{h}} U_2 \subset \mathfrak{h}^*.$$

In this setting, define $(T^*S)_r := \{(a, b) \in T^*S \mid a \diamond_{\mathfrak{h}} b \in \mathfrak{h}_r^*\}$. From the first part of the proof we have the following commutative diagram

$$\begin{array}{ccc} G \times \mathfrak{p}_{\text{dom}}^* \times \mathfrak{o}_{\text{dom}} \times (T^*S)_{\text{dom}} & \xrightarrow{\tilde{f} \circ l} & \mathbf{J}_{H^T}^{-1}(0) \\ \downarrow \pi_{H_\mu} & & \downarrow \overline{\varphi} \circ \pi_H \\ G \times_{H_\mu} \left(\mathfrak{p}_{\text{dom}}^* \times \mathfrak{o}_{\text{dom}} \times (T^*S)_{\text{dom}} \right) & \xrightarrow{\mathcal{T}_0} & T^*(G \times_H S) \end{array}$$

The problem is that $\tilde{f} \circ l$ is an injective embedding but it is not clear if it is proper. We will now show that $\mathcal{T}_0 \circ \pi_{H_\mu}$ is a proper map onto its image when

restricted to $G \times \mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r$. To do so let $x_n = (g_n, \nu_n, \lambda_n; a_n, b_n)$ be a sequence in $G \times \mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r$ such that

$$\mathcal{T}_0(\pi_{H_\mu}(x_n)) \longrightarrow \mathcal{T}_0(\pi_{H_\mu}(\bar{g}, \bar{\nu}, \bar{\lambda}; \bar{a}, \bar{b}))$$

with $(\bar{g}, \bar{\nu}, \bar{\lambda}; \bar{a}, \bar{b}) \in G \times \mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r$. We will construct a subsequence $\{x_{\sigma_3(n)}\} \subset \{x_n\}$ which is convergent on $G \times \mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r$.

The map $\bar{\varphi} \circ \pi_H : \mathbf{J}_{H^T}^{-1}(0) \rightarrow T^*(G \times_H S)$ is proper because it is a composition of an homeomorphism and the projection by a compact group. Since $\mathcal{T}_0 \circ \pi_{H_\mu} = \bar{\varphi} \circ \pi_H \circ \tilde{f} \circ l$ then there is a increasing map $\sigma_1 : \mathbb{N} \rightarrow \mathbb{N}$ such that the sequence $\{(\tilde{f} \circ l)(x_{\sigma_1(n)})\}_n$ converges in $\mathbf{J}_{H^T}^{-1}(0) \subset T^*(G \times S)$ (we are just taking a subsequence). But then by uniqueness of the limit there is $h \in H$ such that

$$(\tilde{f} \circ l)(x_{\sigma_1(n)}) \longrightarrow h \cdot^T ((\tilde{f} \circ l)(\bar{g}, \bar{\nu}, \bar{\lambda}; \bar{a}, \bar{b})),$$

but using the expression of f (24) this implies that $a_{\sigma_1(n)} \rightarrow h \cdot \bar{a}$ and $b_{\sigma_1(n)} \rightarrow h \cdot \bar{b}$. By the definition of $(T^*S)_r$ we can choose for each n a pair $(\alpha_n, \beta_n) \in U_1 \times U_2$ satisfying

$$\alpha_n \diamond_{\mathfrak{h}} \beta_n = a_n \diamond_{\mathfrak{h}} b_n.$$

Since $U_1 \times U_2$ is a relatively compact subset of $(T^*S)_{\text{inj}}$ we can find an increasing map $\sigma_2 : \mathbb{N} \rightarrow \mathbb{N}$ such that $\sigma_2(\mathbb{N}) \subset \sigma_1(\mathbb{N})$ and $(\alpha_{\sigma_2(n)}, \beta_{\sigma_2(n)}) \rightarrow (\alpha_\infty, \beta_\infty)$ but then $(\tilde{f} \circ l)(g_{\sigma_2(n)}, \nu_{\sigma_2(n)}, \lambda_{\sigma_2(n)}, \alpha_{\sigma_2(n)}, \beta_{\sigma_2(n)})$ converges and

$$\begin{aligned} & (\tilde{f} \circ l)(g_{\sigma_2(n)}, \nu_{\sigma_2(n)}, \lambda_{\sigma_2(n)}, \alpha_{\sigma_2(n)}, \beta_{\sigma_2(n)}) \\ & \longrightarrow h \cdot^T \left((\tilde{f} \circ l)(\bar{g}, \bar{\nu}, \bar{\lambda}, h^{-1}\alpha_\infty, h^{-1}\beta_\infty) \right). \end{aligned}$$

As $(g_{\sigma_2(n)}, \nu_{\sigma_2(n)}, \lambda_{\sigma_2(n)}, \alpha_{\sigma_2(n)}, \beta_{\sigma_2(n)})$ lies in $G \times \mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_{\text{inj}}$ by the facts that \mathcal{T}_0 restricted to $G \times \mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_{\text{inj}}$ is a diffeomorphism and π_{H_μ} is proper then there is an increasing map $\sigma_3 : \mathbb{N} \rightarrow \mathbb{N}$ with $\sigma_3(\mathbb{N}) \subset \sigma_2(\mathbb{N})$ such that $(g_{\sigma_3(n)}, \nu_{\sigma_3(n)}, \lambda_{\sigma_3(n)}, \alpha_{\sigma_3(n)}, \beta_{\sigma_3(n)})$ converges in $G \times \mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_{\text{inj}}$. Therefore, $\{x_{\sigma_3(n)}\}$ is a convergent sequence on $G \times \mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r$.

This proves that

$$\begin{aligned} & \mathcal{T}_0 \circ \pi_{H_\mu} : G \times \mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r \\ & \rightarrow (\mathcal{T}_0 \circ \pi_{H_\mu})(G \times \mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r) \end{aligned}$$

is proper. But since

$$\pi_{H_\mu} : G \times \mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r \rightarrow G \times_{H_\mu} (\mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r)$$

is surjective and continuous this implies that

$$\mathcal{T}_0 : G \times_{H_\mu} (\mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r) \longrightarrow \mathcal{T}_0(G \times_{H_\mu} (\mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r))$$

is a proper map.

As \mathcal{T}_0 is a local homeomorphism which is also proper it is a covering map. Then if $\varphi(e, \mu, 0, 0)$ has only one preimage, this implies that the covering map is in fact injective. But if $\mathcal{T}_0([g, \nu, \lambda, a, b]_{H_\mu}) = \varphi(e, \mu, 0, 0)$ then it is clear from the expression of \mathcal{T}_0 (23) that $a = b = 0$ and then $[g, \nu, \lambda, 0, 0]_{H_\mu} \in G \times_{H_\mu} (\mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_{\text{inj}}) \subset G \times_{H_\mu} (\mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r)$. Therefore, by injectivity we have $[g, \nu, \lambda, 0, 0]_{H_\mu} = [e, 0, 0, 0, 0]_{H_\mu}$. But as in the first step of this proof we showed that \mathcal{T}_0 is a local symplectomorphism the injectivity implies that \mathcal{T}_0 is a diffeomorphism. That is, the map

$$\begin{aligned} \mathcal{T}_0 : G \times_{H_\mu} (\mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r) &\longrightarrow \mathcal{T}_0(G \times_{H_\mu} (\mathfrak{p}_{\text{inj}}^* \times \mathfrak{o}_{\text{inj}} \times (T^*S)_r)) \\ &\subset T^*(G \times_H S) \\ [g, \nu, \lambda; a, b]_{H_\mu} &\longmapsto \varphi(\Phi(g, \tilde{\nu}, \lambda; a \diamond_{\mathfrak{h}_\mu} b); a, b) \end{aligned}$$

where $\tilde{\nu} = \nu + \mathbf{J}_{N_0}(\lambda, a, b) = \nu + \frac{1}{2}\lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu + a \diamond_{\mathfrak{h}_\mu} b$, is a diffeomorphism. □

5.3. The Γ map

In this section we are going to introduce the Γ map, a technical tool used in [28] to build the Hamiltonian cotangent tube when $G = G_\mu$. Here we will use it as the final step towards generalizing the previous tube at $\alpha = 0$ to the general case $\alpha \neq 0$.

Let $\varphi(e, \mu, 0, \alpha) \in T^*(G \times_H S)$ and define $K = H_\mu \cap H_\alpha$. Recall that in (22) we defined $B := (\mathfrak{h}_\mu \cdot \alpha)^\circ \subset S$ and a K -invariant complement \mathfrak{s} of \mathfrak{k} in \mathfrak{h}_μ . As K is compact we can choose a K -invariant splitting $S = B \oplus C$ inducing the K -invariant splitting $S^* = B^* \oplus C^*$.

However the previous splitting of S^* is not in principle H_μ -invariant. The following technical result studies how it behaves with respect to the infinitesimal H_μ -action on S . The next result is a straightforward generalization to the case $\mathfrak{g}_\mu \neq \mathfrak{g}$ of Lemmas 27 and 28 of [28].

Proposition 5.4. *In the above situation:*

- If $a \in B$, $c \in C$ and $b \in B^*$ then

$$(a + c) \diamond_{\mathfrak{h}_\mu} (\alpha + b) = a \diamond_{\mathfrak{h}_\mu} b + c \diamond_{\mathfrak{s}} (\alpha + b).$$

- There is a K -invariant neighborhood $(B^*)_r$ of the origin in B^* and a K -equivariant map

$$\Gamma : \mathfrak{s}^* \times (B^*)_r \longrightarrow S$$

defined by

$$\langle \Gamma(\nu; b), \xi \cdot (\alpha + b) + \beta \rangle = \langle \nu, \xi \rangle \quad \forall \beta \in B^*, \quad \forall \xi \in \mathfrak{s}.$$

That is, Γ satisfies $\Gamma(\nu; b) \diamond_{\mathfrak{s}} (b + \alpha) = \nu$ and $\Gamma(\nu; b) \in C$ for any $\nu \in \mathfrak{s}^*$ and $b \in (B^*)_r$.

With the notation that we have already introduced the symplectic slice at $\varphi(e, \mu, 0, 0)$ is $N_0 = \mathfrak{o} \times T^*S$ whereas the symplectic slice at $\varphi(e, \mu, 0, \alpha)$ is $N_\alpha = \mathfrak{o} \times T^*B$ (see Theorem 3.4). The abstract MGS models at $\varphi(e, \mu, 0, 0)$ and $\varphi(e, \mu, 0, \alpha)$ are $G \times_{H_\mu} (\mathfrak{p}^* \times N_0)$ and $G \times_K (\mathfrak{s}^* \oplus \mathfrak{p}^* \times N_\alpha)$ respectively. The next result shows that Γ can be used to build a well-behaved map between both spaces.

Theorem 5.5. *In the above context there is an open K -invariant neighborhood W of zero in $(\mathfrak{s}^* \oplus \mathfrak{p}^*) \times \mathfrak{o} \times B \times B^*$ such that the G -equivariant map*

$$\begin{aligned} \mathcal{F} : G \times_K W &\longrightarrow G \times_{H_\mu} (\mathfrak{p}^* \times \mathfrak{o} \times S \times S^*) \\ [g, \nu_{\mathfrak{s}} + \nu_{\mathfrak{p}}, \lambda, a, b]_K &\longmapsto [g, \nu_{\mathfrak{p}}, \lambda, \tilde{a}, b + \alpha]_{H_\mu}, \end{aligned}$$

where $\tilde{a} = a + \Gamma(\nu_{\mathfrak{s}} - a \diamond_{\mathfrak{s}} b - \frac{1}{2} \lambda \diamond_{\mathfrak{s}} \text{ad}_\lambda^* \mu; b)$, is a local symplectomorphism.

Proof. As in the first part of the proof of Theorem 5.2 there is a neighborhood $(\mathfrak{g}_\mu^*)_r$ of $0 \in \mathfrak{g}^*$ such that $Z_0 := G \times (\mathfrak{g}_\mu^*)_r \times (\mathfrak{o} \times S \times S^*)$ is a symplectic space with $\omega_{Z_0} := \omega_{T_\mu} + \Omega_{N_0}$. We are in the same setting as in Subsection 2.3, therefore, Z_0 supports G^L and H_μ^T Hamiltonian actions with momentum maps that were denoted as \mathbf{K}_{G^L} and $\mathbf{K}_{H_\mu^T}$.

Similarly $Z_\alpha := G \times (\mathfrak{g}_\mu^*)_r \times (\mathfrak{o} \times B \times (B^*)_r)$ is a symplectic space with symplectic form $\omega_{Z_\alpha} := \omega_{T_\mu} + \Omega_{N_\alpha}$ because $N_\alpha = \mathfrak{o} \times B \times B^*$. Note that Z_α has G^L and K^T Hamiltonian actions with momentum maps \mathbf{M}_{G^L} and \mathbf{M}_{K^T} . Consider now the map

$$f : Z_\alpha \longrightarrow Z_0$$

$$(g, \nu, \lambda; a, b) \longmapsto (g, \nu, \lambda; a + \Gamma(\eta; b), b + \alpha)$$

where $\eta = \nu|_{\mathfrak{s}} - a \diamond_{\mathfrak{s}} b - \frac{1}{2} \lambda \diamond_{\mathfrak{s}} \text{ad}_\lambda^* \mu$. As Γ is K -equivariant then f is $G^L \times K^T$ equivariant. Note that the potential $\theta_{Z_0}(g, \nu, \lambda; a, b)(\xi, \dot{\nu}, \dot{\lambda}; \dot{a}, \dot{b}) = \langle \nu + \mu, \xi \rangle + \frac{1}{2} \langle \mu, [\lambda, \dot{\lambda}] \rangle + \langle b, \dot{a} \rangle - \langle \alpha, \dot{a} \rangle$ generates the symplectic structure ω_{Z_0} (see (4)) and

$$(f^* \theta_{Z_0})(g, \nu, \lambda, a, b) \cdot v = \theta_{Z_0}(f(g, \nu, \lambda, a, b))(T_{(g, \nu, \lambda, a, b)} f \cdot v)$$

$$= \langle \mu + \nu, \xi \rangle + \frac{1}{2} \langle \mu, [\lambda, \dot{\lambda}] \rangle$$

$$+ \langle b, \dot{a} + T_{(\nu, \lambda, a, b)} \Gamma \cdot (\dot{\nu}, \dot{\lambda}, \dot{a}, \dot{b}) \rangle$$

$$= \langle \mu + \nu, \xi \rangle + \frac{1}{2} \langle \mu, [\lambda, \dot{\lambda}] \rangle + \langle b, \dot{a} \rangle$$

where $v = (\xi, \dot{\nu}, \dot{a}, \dot{b}) \in T_{(g, \nu, \lambda, a, b)} Z_\alpha$. Taking the exterior derivative of this equality we get $f^* \omega_{Z_0} = \omega_{Z_\alpha}$.

Additionally, the H_μ^T -momentum evaluated at $f(g, \nu, \lambda, a, b)$ is

$$(25) \quad \mathbf{K}_{H_\mu^T}(f(g, \nu, \lambda; a, b))$$

$$= -\nu + \frac{1}{2} \lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu + (a + \Gamma(\eta; b)) \diamond_{\mathfrak{h}_\mu} (b + \alpha)$$

$$= -\nu|_{\mathfrak{h}_\mu} + \frac{1}{2} \lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu + a \diamond_{\mathfrak{h}_\mu} b + \Gamma(\eta; b) \diamond_{\mathfrak{s}} (b + \alpha)$$

$$= -\nu|_{\mathfrak{h}_\mu} + \frac{1}{2} \lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu + a \diamond_{\mathfrak{h}_\mu} b + \eta$$

$$= -\nu|_{\mathfrak{h}_\mu} + \frac{1}{2} \lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu + a \diamond_{\mathfrak{h}_\mu} b + \nu|_{\mathfrak{s}} - a \diamond_{\mathfrak{s}} b - \frac{1}{2} \lambda \diamond_{\mathfrak{s}} \text{ad}_\lambda^* \mu$$

$$= -\nu|_{\mathfrak{k}} + \frac{1}{2} \text{ad}_\lambda^* \mu \diamond_{\mathfrak{k}} \lambda + a \diamond_{\mathfrak{k}} b.$$

This means that if $\xi \in \mathfrak{s}$ then $\langle \mathbf{K}_{H_\mu^T}(f(g, \nu, \lambda, a, b)), \xi \rangle = 0$ and as

$$\mathbf{M}_{K^T}(g, \nu, \lambda, a, b) = -\nu|_{\mathfrak{k}} + \frac{1}{2} \text{ad}_\lambda^* \mu \diamond_{\mathfrak{k}} \lambda + a \diamond_{\mathfrak{k}} b,$$

f can be restricted to

$$\tilde{f} : \mathbf{M}_{K^T}^{-1}(0) \longrightarrow \mathbf{K}_{H_\mu}^{-1}(0).$$

As we did in Theorem 5.2 we can construct from the top to the bottom all the arrows of the diagram

$$\begin{array}{ccc} Z_\alpha & \xrightarrow{f} & Z_0 \\ \uparrow & & \uparrow \\ \mathbf{M}_{K^T}^{-1}(0) & \xrightarrow{\tilde{f}} & \mathbf{K}_{H_\mu}^{-1}(0) \\ \downarrow & \searrow & \downarrow \\ \mathbf{M}_{K^T}^{-1}(0)/K^T & \xrightarrow{F} & \mathbf{K}_{H_\mu}^{-1}(0)/H_\mu^T \end{array}$$

using the same arguments as in the first part of Theorem 5.2. Therefore, as F is an immersion between spaces of the same dimension, it is a local diffeomorphism onto its image.

Adapting the construction of map (7) to this setting define the K -invariant open set

$$W := \left\{ (\nu_{\mathfrak{s}} + \nu_{\mathfrak{p}}, \lambda, a, b) \in (\mathfrak{s}^* \oplus \mathfrak{s}^*) \times \mathfrak{o} \times B \times (B^*)_r \mid \nu_{\mathfrak{s}} + \nu_{\mathfrak{p}} + \underbrace{\frac{1}{2} \text{ad}^*_{\lambda} \mu \diamond_{\mathfrak{k}} \lambda + a \diamond_{\mathfrak{k}} b}_{\mathbf{J}_{N_\alpha}} \in (\mathfrak{g}^*_\mu)_r \right\}$$

and the map $L_\alpha : G \times_K W \longrightarrow \mathbf{M}_{K^T}^{-1}(0)/K^T$ given by

$$L_\alpha([g, \nu_{\mathfrak{s}} + \nu_{\mathfrak{p}}, \lambda, a, b]_K) = [g, \nu_{\mathfrak{s}} + \nu_{\mathfrak{p}} + \mathbf{J}_{N_\alpha}(\lambda, a, b), \lambda, a, b]_K,$$

then L_α is a G -equivariant symplectomorphism. And similarly,

$$R_0 : \mathbf{K}_{H_\mu}^{-1}(0)/H_\mu^T \rightarrow G \times_{H_\mu} (\mathfrak{p}^* \times \mathfrak{o} \times S \times S^*)$$

defined by $[g, \nu, \lambda, a, b]_{H_\mu} \mapsto [g, \nu|_{\mathfrak{p}}, \lambda, a, b]_{H_\mu}$ is a G -equivariant symplectomorphism.

Finally, we can conclude that the composition $\mathcal{F} = R_0 \circ F \circ L_\alpha : G \times_K W \rightarrow G \times_{H_\mu} (\mathfrak{p}^* \times \mathfrak{o} \times S \times S^*)$ is a G -equivariant local diffeomorphism that

pulls-back the MGS symplectic form of the target to the MGS symplectic form of $G \times_K W$. □

5.4. General tube

In this section we will deal with the most general situation and will construct a Hamiltonian tube around an arbitrary point $\varphi(e, \mu, 0, \alpha)$. To do so we will use Theorem 5.2 to obtain a Hamiltonian tube around $\varphi(e, \mu, 0, 0)$ and then we will compose it with the map \mathcal{F} of Theorem 5.5. The result of this composition will be the desired Hamiltonian tube around $\varphi(e, \mu, 0, \alpha)$.

Theorem 5.6. *Consider the point $z \in T^*(G \times_H S)$ defined by $z = \varphi(e, \mu, 0, \alpha)$. Let $\mathfrak{g} = \mathfrak{g}_\mu \oplus \mathfrak{o} \oplus \mathfrak{l} \oplus \mathfrak{n}$ be an adapted splitting in the sense of Proposition 3.1 and let $\Phi : G \times U_\Phi \rightarrow T^*G$ be an associated restricted G -tube. Let $\mathfrak{g}_\mu = \mathfrak{h}_\mu \oplus \mathfrak{p}$ be a H_μ -invariant splitting and $\mathfrak{h}_\mu = \mathfrak{g}_z \oplus \mathfrak{s}$ a G_z invariant splitting. Define $B = (\mathfrak{h}_\mu \cdot \alpha)^\circ \subset S$ and let the map $\Gamma : \mathfrak{s}^* \times B_r^* \rightarrow S$ be the one defined in Proposition 5.4.*

In these conditions there are small enough G_z -invariant neighborhoods of zero $\mathfrak{s}_r^ \subset \mathfrak{s}^*$, $\mathfrak{p}_r^* \subset \mathfrak{p}^*$, $\mathfrak{o}_r \subset \mathfrak{o}$, $B_r \subset B$ and $B_r^* \subset B^*$ such that the map*

$$(26) \quad \mathcal{T} : G \times_{G_z} ((\mathfrak{s}_r^* \oplus \mathfrak{p}_r^*) \times \mathfrak{o}_r \times B_r \times B_r^*) \longrightarrow T^*(G \times_H S)$$

$$[g, \nu_{\mathfrak{s}} + \nu_{\mathfrak{p}}, \lambda; a, b]_{G_z} \longmapsto \varphi(\Phi(g, \tilde{\nu}, \lambda; \varepsilon); \tilde{a}, b + \alpha)$$

where

$$\tilde{a} = a + \Gamma \left(\nu_{\mathfrak{s}} - a \diamond_{\mathfrak{s}} b - \frac{1}{2} \lambda \diamond_{\mathfrak{s}} \text{ad}_\lambda^* \mu; b \right)$$

$$\tilde{\nu} = \nu_{\mathfrak{p}} + \nu_{\mathfrak{s}} + \frac{1}{2} \lambda \diamond_{\mathfrak{g}_z} \text{ad}_\lambda^* \mu + a \diamond_{\mathfrak{g}_z} b$$

$$\varepsilon = \tilde{a} \diamond_{\mathfrak{l}} (b + \alpha)$$

is a Hamiltonian tube around the point $z = \varphi(e, \mu, 0, \alpha)$.

Proof. Using Theorem 5.2, there is a Hamiltonian tube $\mathcal{T}_0 : G \times_{H_\mu} V \rightarrow T^*(G \times_H S)$ with $V \subset \mathfrak{p}^* \times \mathfrak{o} \times S \times S^*$ satisfying $\mathcal{T}_0([e, 0]_{H_\mu}) = \varphi(e, \mu, 0, 0)$. By Theorem 5.5, there is a map $\mathcal{F} : G \times_{G_z} W \rightarrow G \times_{H_\mu} (\mathfrak{p}^* \times \mathfrak{o} \times T^*S)$ with $W \subset (\mathfrak{p}^* \oplus \mathfrak{s}^*) \times \mathfrak{o} \times B \times B^*$.

Note that $G \times_{H_\mu} V$ is an open G -invariant subset of $G \times_{H_\mu} (\mathfrak{p}^* \times \mathfrak{o} \times T^*S)$. Since \mathcal{F} is continuous then the preimage of $G \times_{H_\mu} V$ by \mathcal{F} is open

and it contains the point $[e, 0, 0, 0, 0]_{G_z}$ because

$$\mathcal{F}([e, 0, 0, 0, 0]_{G_z}) = [e, 0, 0, 0, \alpha]_{H_\mu} \subset G \times_{H_\mu} V$$

since $0 \diamond_{\mathfrak{h}} \alpha = 0 \in \mathfrak{h}_r^*$. Therefore, we can choose small enough G_z -invariant neighborhoods of zero $\mathfrak{s}_r^* \subset \mathfrak{s}^*$, $\mathfrak{p}_r^* \subset \mathfrak{p}^*$, $\mathfrak{o}_r \subset \mathfrak{o}$, $B_r \subset B$ and $B_r^* \subset B^*$ such that the composition $\mathcal{T} := \mathcal{T}_0 \circ \mathcal{F} : G \times_{G_z} (\mathfrak{s}_r^* \oplus \mathfrak{p}_r^*) \times \mathfrak{o}_r \times B_r \times B_r^* \rightarrow T^*(G \times_H S)$ is well-defined and injective. Using Theorems 5.2 and 5.5 we conclude that \mathcal{T} is a Hamiltonian tube around $\varphi(e, \mu, 0, \alpha)$.

More precisely, as $\mathcal{F}([g, \nu_{\mathfrak{s}} + \nu_{\mathfrak{p}}, \lambda, a, b]_{G_z}) = [g, \nu_{\mathfrak{p}}, \lambda, \tilde{a}, b + \alpha]_{H_\mu}$ with $\tilde{a} = a + \Gamma(\nu_{\mathfrak{s}} - a \diamond_{\mathfrak{s}} b - \frac{1}{2} \lambda \diamond_{\mathfrak{s}} \text{ad}_\lambda^* \mu; b)$ then $(\mathcal{T}_0 \circ \mathcal{F})([g, \nu_{\mathfrak{s}} + \nu_{\mathfrak{p}}, \lambda; a, b]_{G_z}) = \varphi(\Phi(g, \tilde{\nu}, \lambda; \varepsilon); \tilde{a}, b + \alpha)$ where $\varepsilon = \tilde{a} \diamond_{\mathfrak{l}} (b + \alpha)$ and

$$\tilde{\nu} = \nu_{\mathfrak{p}} + \frac{1}{2} \lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu + \tilde{a} \diamond_{\mathfrak{h}_\mu} (b + \alpha),$$

but using exactly the same computations that as in (25) we get $\tilde{a} \diamond_{\mathfrak{h}_\mu} (b + \alpha) = a \diamond_{\mathfrak{g}_z} b + \nu_{\mathfrak{s}} - \frac{1}{2} \lambda \diamond_{\mathfrak{s}} \text{ad}_\lambda^* \mu$, that is $\tilde{\nu} = \nu_{\mathfrak{p}} + \nu_{\mathfrak{s}} + \frac{1}{2} \lambda \diamond_{\mathfrak{g}_z} \text{ad}_\lambda^* \mu + a \diamond_{\mathfrak{g}_z} b$. □

Note that if we assume that $\mu \in \mathfrak{g}^*$ satisfies $\mathfrak{g}_\mu = \mathfrak{g}$ then $\mathfrak{o} = 0$ and the Hamiltonian tube \mathcal{T} will be of the form

$$\begin{aligned} \mathcal{T} : G \times_{G_z} ((\mathfrak{s}_r^* \oplus \mathfrak{p}_r^*) \times B_r \times B_r^*) &\longrightarrow T^*(G \times_H S) \\ [g, \nu_{\mathfrak{s}} + \nu_{\mathfrak{p}}; a, b]_{G_z} &\longmapsto \varphi(g, \mu + \tilde{\nu}; \tilde{a}, b + \alpha) \end{aligned}$$

where

$$\tilde{\nu} = \nu_{\mathfrak{p}} + \nu_{\mathfrak{s}} + a \diamond_{\mathfrak{g}_z} b, \quad \tilde{a} = a + \Gamma(\nu_{\mathfrak{s}} - a \diamond_{\mathfrak{s}} b; b).$$

This map is the content of Theorem 31 of [28]. That is, the map \mathcal{T} coincides with the results of [28] when we restrict to their totally isotropic hypothesis $\mathfrak{g} = \mathfrak{g}_\mu$. What happens in this case is that the map given by Theorem 5.2 becomes the trivial μ -shift

$$\begin{aligned} \mathcal{T}_0 : G \times_{H_\mu} (\mathfrak{p}^* \times S \times S^*) &\longrightarrow T^*(G \times_H S) \\ [g, \nu; a, b]_{H_\mu} &\longmapsto \varphi(g, \mu + \nu + a \diamond_{\mathfrak{h}_\mu} b; a, b). \end{aligned}$$

The other extreme case is when Γ becomes trivial. This will happen for example if $S = 0$, which is equivalent to assuming that locally $Q = G/H$. Fix a point $z = \varphi(e, \mu) \in T^*(G/H)$, as $G_z = H_\mu$ then $\mathfrak{s} = 0$ and as $S = 0$

then $B = 0$, therefore, according to (26) \mathcal{T} becomes

$$\begin{aligned} \mathcal{T} : G \times_{H_\mu} (\mathfrak{p}_r^* \times \mathfrak{o}_r) &\longrightarrow T^*(G/H) \\ [g, \nu_{\mathfrak{p}}, \lambda]_{H_\mu} &\longmapsto \varphi(\Phi(g, \nu_{\mathfrak{p}} + \frac{1}{2}\lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu, \lambda; 0)). \end{aligned}$$

6. A fibered Bates-Lerman Lemma

One of the most important consequences of the MGS model is that it provides a local description of the set of points with momentum μ which is very useful in the theory of singular reduction. This is the content of the following result that appeared in [2] and previously in [31] for $\mu = 0$. More precisely, instead of using the orbit-reduction-approach of [2] we are going to need the point-reduction-approach of Proposition 8.1.2 of [19].

Proposition 6.1. *Let (M, ω) be a symplectic manifold supporting a proper Hamiltonian G -action with momentum map \mathbf{J} . Let $m \in M$, $\mu = \mathbf{J}(m)$ and $\mathcal{T} : G \times_{G_m} (\mathfrak{m}_r^* \times N_r) \rightarrow M$ a Hamiltonian tube around m . There is an open G_μ -invariant neighborhood U_M of $G_\mu \cdot m$ such that*

$$U_M \cap \mathbf{J}^{-1}(\mu) = \mathcal{T}(Z)$$

where

$$Z = \{[g, \nu, \lambda, v]_{G_m} \in \mathcal{T}^{-1}(U_M) \mid g \in G_\mu, \nu = 0, \mathbf{J}_N(v) = 0\}.$$

Recall that, in some cases, we saw that the domain of the cotangent bundle Hamiltonian tube is unbounded in the S^* direction (see Theorem 5.2). In this section we will present an important consequence of this which is that for cotangent-lifted actions the open neighborhood U_M can be global in the vertical direction. That is, it will be of the form $\tau^{-1}(\mathcal{U}_Q)$ where $\tau : T^*Q \rightarrow Q$ is the natural projection and \mathcal{U}_Q is a neighborhood in Q .

Proposition 6.2. *Consider the Hamiltonian tube $\mathcal{T}_0 : G \times_{H_\mu} (\mathfrak{p}_r^* \times \mathfrak{o}_r \times (T^*S)_r) \rightarrow T^*(G \times_H S)$ of Theorem 5.2 at the point $\varphi(e, \mu, 0, 0) \in T^*(G \times_H S)$. There is a G_μ -invariant neighborhood \mathcal{U}_Q of $[e, 0]_H \in G \times_H S$ such that*

$$(27) \quad \tau^{-1}(\mathcal{U}_Q) \cap \mathbf{J}^{-1}(\mu) = \mathcal{T}_0(Z)$$

where

$$Z = \left\{ [g, \nu, \lambda, a, b]_{H_\mu} \in \mathcal{T}_0^{-1}(\tau^{-1}(\mathcal{U}_Q)) \mid \right. \\ \left. g \in G_\mu, \nu = 0, \frac{1}{2}\lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu + a \diamond_{\mathfrak{h}_\mu} b = 0 \right\}.$$

Proof. As in the previous section, $N_0 = \mathfrak{o} \times T^*S$ will be the symplectic slice at $\varphi(e, \mu, 0, 0)$ with the symplectic form (10). The Hamiltonian tube puts the momentum map \mathbf{J} in the normal form (3), that is $\mathbf{J} \circ \mathcal{T}_0 = \mathbf{J}_Y$. We will now proceed as in the proof of Proposition 6.1 in [2] and we will factorize $\mathbf{J}_Y = \gamma \circ \beta$, with

$$\beta : G \times_{H_\mu} (\mathfrak{p}^* \times N_0) \longrightarrow G \times_{H_\mu} \mathfrak{g}_\mu^*, \quad \gamma : G \times_{H_\mu} \mathfrak{g}_\mu^* \longrightarrow \mathfrak{g}^* \\ [g, \nu, v]_{H_\mu} \longmapsto [g, \nu + \mathbf{J}_{N_0}(v)] \quad [g, \nu]_{H_\mu} \longmapsto \text{Ad}_{g^{-1}}^*(\mu + \nu).$$

Using this factorization it will be easy to describe $\mathbf{J}_Y^{-1}(\mu)$. Note that, since the map

$$T_{[e,0]_{H_\mu}} \gamma \cdot (\xi, \dot{\nu}) = -\text{ad}_\xi^* \mu + \dot{\nu}$$

is onto, γ is a submersion near $[e, 0]_{H_\mu}$ but by G -equivariance there is a G -invariant open set $U_{\text{subm}} \subset G \times_{H_\mu} \mathfrak{g}^*$ where γ is a submersion. Therefore, $\gamma^{-1}(\mu) \cap U_{\text{subm}}$ is a manifold of dimension $\dim G_\mu - \dim H_\mu$. As $G_\mu \times_{H_\mu} \{0\} \subset \gamma^{-1}(\mu)$, $G_\mu \times_{H_\mu} \{0\}$ has to be an open submanifold of $\gamma^{-1}(\mu) \cap U$, that is, there is an open set $U_{\text{BL}} \subset U_{\text{subm}}$ with $G_\mu \times_{H_\mu} \{0\} = \gamma^{-1}(\mu) \cap U_{\text{BL}}$. By equivariance of γ we can assume that U_{BL} is G_μ -invariant. Applying β^{-1} on the equality $G_\mu \times_{H_\mu} \{0\} = \gamma^{-1}(\mu) \cap U_{\text{BL}}$, we get

$$(28) \quad \{[g, 0, v]_{H_\mu} \in G \times_{H_\mu} (\mathfrak{p}^* \times N_0) \mid \mathbf{J}_{N_0}(v) = 0\} = \mathbf{J}_Y^{-1}(\mu) \cap \beta^{-1}(U_{\text{BL}}).$$

In this setting, let U_G be a $G_\mu^L \times H_\mu^R$ -invariant neighborhood of $e \in G$, $\mathfrak{p}_0^* \subset \mathfrak{p}^*$, $\mathfrak{o}_0 \subset \mathfrak{o}$ H_μ -invariant neighborhoods of zero and $\mathfrak{h}_0^* \subset \mathfrak{h}^*$ a H -invariant neighborhood of zero such that

$$\{[g, \nu, \lambda, a, b]_{H_\mu} \mid g \in U_G, \nu \in \mathfrak{p}_0^*, \lambda \in \mathfrak{o}_0, a \diamond_{\mathfrak{h}} b \in \mathfrak{h}_0^*\} \\ \subset (G \times_{H_\mu} (\mathfrak{p}_r^* \times \mathfrak{o}_r \times (T^*S)_r)) \cap \beta^{-1}(U_{\text{BL}}).$$

Let now Φ be the restricted tube used in the definition of \mathcal{T}_0 (see Theorem 5.2) and consider the map

$$f : U_G \times \mathfrak{p}_0^* \times \mathfrak{o}_0 \times \mathfrak{h}_0^* \longrightarrow T^*G/H^T$$

$$(g, \nu, \lambda, \rho) \longmapsto \pi_{H^T}(\Phi(g, \nu + \frac{1}{2}\lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu + \rho|_{\mathfrak{h}_\mu}, \lambda, \rho|_{\mathfrak{l}})).$$

Note that this expression is very similar to (23) but we have changed the dependence on T^*S by a dependence on \mathfrak{h}_μ^* . We claim that this map is a submersion at $(e, 0, 0, 0)$. Using the notation of the proof of Proposition 4.8, if $v = (\xi, \dot{\nu}, \dot{\lambda}, \dot{\varepsilon}) \in T_{(e,0,0,0)}(G \times \mathfrak{g}_\mu^* \times \mathfrak{o} \times \mathfrak{l}^*)$,

$$T_{(e,0,0,0)}\Phi \cdot v = T_{(e,0,0,0)}(\Theta \circ \iota_W \circ \psi^{-1}) \cdot v = T_{(e,0,0,0)}\Theta \cdot (\xi, \dot{\nu}, \dot{\lambda} + \sigma^{-1} \cdot \varepsilon)$$

$$= (\xi + \dot{\lambda} + \sigma^{-1} \cdot \dot{\varepsilon}, \dot{\nu} - \dot{\varepsilon} + \text{ad}_\lambda^* \mu)$$

where $\sigma : \mathfrak{n} \rightarrow \mathfrak{l}^*$ is the linear isomorphism $\sigma \cdot \zeta = \text{ad}_\zeta^* \mu|_{\mathfrak{l}}$. Applying this result to f we get

$$(29) \quad T_{(e,0,0,0)}f \cdot (\xi, \dot{\nu}, \dot{\lambda}, \dot{\rho}) = T_{(e,\mu)}\pi_{H^T} \cdot (\xi + \dot{\lambda} + \sigma^{-1} \cdot \rho|_{\mathfrak{l}}, \dot{\nu} + \text{ad}_\lambda^* \mu + \dot{\rho}).$$

But, since the splitting of Proposition 3.1 induces the dual decomposition $\mathfrak{g}^* = \mathfrak{h}_\mu^* \oplus \mathfrak{p}^* \oplus \mathfrak{o}^* \oplus \mathfrak{l}^* \oplus \mathfrak{n}^*$ each element of \mathfrak{g}^* can be expressed as $\dot{\rho}|_{\mathfrak{h}_\mu} + \dot{\nu} + \text{ad}_\lambda^* \mu + \dot{\rho}|_{\mathfrak{l}} + \text{ad}_\eta^* \mu$ for some $\dot{\rho} \in \mathfrak{h}^*$, $\dot{\nu} \in \mathfrak{p}^*$, $\dot{\lambda} \in \mathfrak{o}^*$ and $\eta \in \mathfrak{h}$. Finally $\text{Ker}(T_{(e,\mu)}\pi_{H^T}) = \{(\eta, \text{ad}_\eta^* \mu) \mid \eta \in \mathfrak{h}\}$ and (29) imply that $T_{(e,0,0,0)}f$ is a surjective linear map.

Since f is G -equivariant f is a submersion on a neighborhood of $G \cdot (e, 0, 0, 0)$ and since submersions are open maps the image of f contains an open neighborhood of $G \cdot [e, \mu]_H$ so there must exist a neighborhood $U_{\mathfrak{g}^*}$ of $\mu \in \mathfrak{g}^*$ such that

$$(30) \quad \pi_H(G \times_H U_{\mathfrak{g}^*}) \subset f(U_G \times \mathfrak{p}_0^* \times \mathfrak{o}_0 \times \mathfrak{h}_0^*).$$

Define $\mathcal{U}_G := \{g \in U_G \mid \text{Ad}_g^* \mu \in U_{\mathfrak{g}^*}\} \cap \{g \in U_G \mid \text{Ad}_g^* \mu|_{\mathfrak{h}} \in \mathfrak{h}_0^*\}$ which is an open neighborhood of $e \in G$ and let $\mathcal{U}_Q := \mathcal{U}_G \times_H S$ which is an open neighborhood of $[e, 0]_H \in Q$. We will now check that \mathcal{U}_Q satisfies (27).

- $\tau^{-1}(\mathcal{U}_Q) \cap \mathbf{J}^{-1}(\mu) \supset \mathcal{T}_0(Z)$.

This inclusion is trivial because if $[g, 0, \lambda, a, b]_{H_\mu} \in Z$ then

$$\begin{aligned} \mathcal{T}_0([g, 0, \lambda, a, b]_{H_\mu}) &\in \tau^{-1}(\mathcal{U}_Q) \quad \text{and} \\ (\mathbf{J} \circ \mathcal{T}_0)([g, 0, \lambda, a, b]_{H_\mu}) &= \mathbf{J}_Y([g, 0, \lambda, a, b]_{H_\mu}) \\ &= \text{Ad}_{g^{-1}}^*(\mu + \mathbf{J}_{N_0}(\lambda, a, b)) = \text{Ad}_{g^{-1}}^*\mu = \mu. \end{aligned}$$

- $\tau^{-1}(\mathcal{U}_Q) \cap \mathbf{J}^{-1}(\mu) \subset \mathcal{T}_0(Z)$.

Let $z \in \tau^{-1}(\mathcal{U}_Q) \cap \mathbf{J}^{-1}(\mu)$. Using the cotangent reduction map φ (see (21)), there is an element (g, ν, a, b) such that $\varphi(g, \nu, a, b) = z$, but, as $\tau(z) \in \mathcal{U}_Q$ then $g \in \mathcal{U}_G$. Since $\varphi(g, \nu, a, b) \in \mathbf{J}^{-1}(\mu)$, using (6) we have $\text{Ad}_{g^{-1}}^*\nu = \mu$. Additionally, as $(g, \nu, a, b) \in \mathbf{J}_{H^T}^{-1}(0)$, then $\nu|_{\mathfrak{h}} = a \diamond_{\mathfrak{h}} b$. Using $\nu = \text{Ad}_g^*\mu$ this implies the relation $(\text{Ad}_g^*\mu)|_{\mathfrak{h}} = a \diamond_{\mathfrak{h}} b$.

As $g \in \mathcal{U}_G$, using the definition of \mathcal{U}_G we have $(g, \text{Ad}_g^*\mu) \in G \times U_{\mathfrak{g}^*}$. Equation (30) implies that there is a point $(g', \nu', \lambda, \rho) \in U_G \times \mathfrak{p}_0^* \times \mathfrak{o}_0 \times \mathfrak{h}_0^*$ such that $f(g', \nu', \lambda, \rho) = [g, \text{Ad}_g^*\mu]_H$. Therefore there is $h \in H$ such that

$$(31) \quad (gh^{-1}, \text{Ad}_{h^{-1}}^*\text{Ad}_g^*\mu) = \Phi \left(g', \nu' + \frac{1}{2}\lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^*\mu + \rho|_{\mathfrak{h}_\mu}, \rho|_{\mathfrak{t}} \right).$$

Moreover, using (17) and (19) it can be checked that the H_μ^T -momentum of a restricted G -tube is $\mathbf{J}_{H_\mu^T}(\Phi(g, \nu, \lambda, \varepsilon)) = -\nu|_{\mathfrak{h}_\mu} + \frac{1}{2}\lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^*\mu$. Therefore taking the H_μ^T -momentum on the previous equation

$$-(\text{Ad}_{h^{-1}}^*\text{Ad}_g^*\mu)|_{\mathfrak{h}_\mu} = -\frac{1}{2}\lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^*\mu - \rho|_{\mathfrak{h}_\mu} + \frac{1}{2}\lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^*\mu = -\rho|_{\mathfrak{h}_\mu}.$$

Now, using item 4. in Definition 4.7 we have that H^T -momentum restricted to $\mathfrak{l}^* \subset \mathfrak{g}^*$ in (31) becomes the equality

$$-(\text{Ad}_{h^{-1}}^*\text{Ad}_g^*\mu)|_{\mathfrak{t}} = -\rho|_{\mathfrak{t}}.$$

That is $(\text{Ad}_{h^{-1}}^*\text{Ad}_g^*\mu)|_{\mathfrak{h}} = \rho$. But, as $\text{Ad}_g^*\mu|_{\mathfrak{h}} = \nu|_{\mathfrak{h}} = a \diamond_{\mathfrak{h}} b$, it follows that $\rho = \text{Ad}_{h^{-1}}^*(a \diamond_{\mathfrak{h}} b) = (h \cdot a) \diamond_{\mathfrak{h}} (h \cdot b)$ and therefore

$$\begin{aligned}
 & \mathcal{T}_0([g', \nu', \lambda, h \cdot a, h \cdot b]_{H_\mu}) \\
 &= \varphi(\Phi(g', \nu' + \frac{1}{2}\lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu + \rho|_{\mathfrak{h}_\mu}, \lambda, (h \cdot a) \diamond_{\mathfrak{l}} (h \cdot b); h \cdot a, h \cdot b) \\
 &= \varphi(\Phi(g', \nu' + \frac{1}{2}\lambda \diamond_{\mathfrak{h}_\mu} \text{ad}_\lambda^* \mu + \rho|_{\mathfrak{h}_\mu}, \lambda, \rho|_{\mathfrak{l}}; h \cdot a, h \cdot b) \\
 &= \varphi(gh^{-1}, \text{Ad}_{h^{-1}}^* \text{Ad}_g^* \mu, h \cdot a, h \cdot b) \\
 &= \varphi(g, \text{Ad}_g^* \mu, a, b) = \varphi(g, \nu, a, b).
 \end{aligned}$$

Finally, as $g \in \mathcal{U}_G$, $\text{Ad}_g^* \mu|_{\mathfrak{h}} \in \mathfrak{h}_0^*$ and \mathfrak{h}_0^* is H -invariant, then $(h \cdot a) \diamond_{\mathfrak{h}} (h \cdot b) = \text{Ad}_{h^{-1}}^*(\text{Ad}_g^* \mu|_{\mathfrak{h}}) \in \mathfrak{h}_0^*$. This observation implies that $(g', \nu', \lambda', h \cdot a, h \cdot b) \in \pi_{H_\mu}^{-1}(\beta^{-1}(U_{\text{BL}}))$. Using the characterization (28), $\nu' = 0$, $g \in G_\mu$ and $\mathbf{J}_{N_0}(\lambda', h \cdot a, h \cdot b) = 0$, that is $[g', 0, \lambda', h \cdot a, h \cdot b]_{H_\mu} \in Z$ as we wanted to show. \square

Remark 6.3. Note that we have started with a tube centered around $\varphi(e, \mu, 0, 0)$. In general if we consider a tube around $\varphi(e, \mu, 0, \alpha)$ we cannot expect its image to be global in the B^* direction. This is because all the points in the model space $G \times_{G_z} ((\mathfrak{s}_r^* \oplus \mathfrak{p}_r^*) \times \mathfrak{o}_r \times B_r \times B_r^*)$ must have G -isotropy conjugate to a subgroup of G_z . From this observation we can conclude that, in general, $(\mathfrak{s}_r^* \oplus \mathfrak{p}_r^*) \times \mathfrak{o}_r \times B_r \times B_r^*$ will not be an open neighborhood containing containing points of the form $(0, 0, 0, b)$ for arbitrary large $b \in B^*$. Indeed, if that was true, we would have $(0, 0, 0, \alpha) \in (\mathfrak{s}_r^* \oplus \mathfrak{p}_r^*) \times \mathfrak{o}_r \times B_r \times B_r^*$ which would imply $\mathcal{T}([e, 0, 0, 0, -\alpha]_{G_z}) = \varphi(e, \mu, 0, 0)$. But this is a point with G -isotropy H_μ and in general, $G_z \subsetneq H_\mu$, producing a contradiction.

7. Explicit examples

In Proposition 4.3 we proved the existence of simple G -tubes using Moser’s trick. In this section we will write down the actual differential equation that must be solved. We will see that if $\dim \mathfrak{q} = 2$ then the simple G -tube will be a scaling of an exponential map and we will compute it explicitly for $SO(3)$ and $SL(2, \mathbb{R})$. Using a different approach, without the explicit $SO(n)$ restricted tube in Subsection 7.4 we will present the Hamiltonian tube for the natural action of $SO(3)$ on $T^*\mathbb{R}^3$ generalizing the final example of [28] to $\mu \neq 0$.

We will compute explicitly the flow that determines a simple G -tube. For that we are going to use the notation of the proof of Proposition 4.3.

Recall that we constructed the simple G -tube as the composition $\Theta = F \circ \Psi_1$ where

$$F(g, \nu, \lambda) = (g \exp(\lambda), \text{Ad}_{\exp(\lambda)}^*(\nu + \mu))$$

(see (14)) and Ψ_1 is the time-1 flow of the time dependent vector field X_t that satisfies the Moser equation associated with $\theta_t = tF^*\theta_{T^*G} + (1-t)\theta_Y$, that is, $i_{X_t}(-\mathbf{d}\theta_t) = \frac{\partial\theta_t}{\partial t}$.

Moser's equation can be written explicitly in this case. Using the above expression of F and (11) we have

$$\begin{aligned} F^*\theta_{T^*G}(g, \nu, \lambda)(\xi, \dot{\nu}, \dot{\lambda}) &= \langle \text{Ad}_{\exp(\lambda)}^*(\nu + \mu), \text{Ad}_{\exp(\lambda)}^{-1}\xi + T_e L_{\exp(\lambda)}^{-1} T_\lambda \exp(\dot{\lambda}) \rangle \\ &= \langle \nu + \mu, \xi \rangle + \langle \nu + \mu, \text{Ad}_{\exp(\lambda)} T_e L_{\exp(\lambda)}^{-1} T_\lambda \exp(\dot{\lambda}) \rangle \\ &= \langle \nu + \mu, \xi \rangle + \langle \nu + \mu, T_e R_{\exp(\lambda)}^{-1} T_\lambda \exp(\dot{\lambda}) \rangle. \end{aligned}$$

The last term can be expressed as a series of Lie brackets (see for example [5])

$$M(\lambda) \cdot \dot{\lambda} := T_e R_{\exp(\lambda)}^{-1} T_\lambda \exp(\dot{\lambda}) = \sum_{n \geq 0} \frac{1}{(n+1)!} \text{ad}_\lambda^n \dot{\lambda}.$$

And, in fact, this series is just the pullback of the right Maurer-Cartan form $\varpi^R(g) = T_e R_g^{-1}$ by the restricted exponential $\exp|_{\mathfrak{q}} : \mathfrak{q} \rightarrow G$. Therefore using the Maurer-Cartan relation

$$\begin{aligned} (32) \quad (\mathbf{d}M)(X, Y) &= \mathbf{d}(\exp^* \varpi^R)(X, Y) = \exp^*(\mathbf{d}\varpi^R)(X, Y) \\ &= [\exp^* \varpi^R(X), \exp^* \varpi^R(Y)] = [M(X), M(Y)]. \end{aligned}$$

Now, since

$$\theta_t(g, \nu, \lambda)(\xi, \dot{\nu}, \dot{\lambda}) = \langle \mu + \nu, \xi \rangle + t \langle \mu + \nu, M(\lambda) \cdot \dot{\lambda} \rangle + (1-t) \frac{1}{2} \langle \mu, [\lambda, \dot{\lambda}] \rangle$$

(using (32)) the exterior derivative $\omega_t = -\mathbf{d}\theta_t$ simplifies as

$$\begin{aligned} \omega_t(g, \nu, \lambda)(\xi_1, \dot{\nu}_1, \dot{\lambda}_1)(\xi_2, \dot{\nu}_2, \dot{\lambda}_2) &= \langle \dot{\nu}_2, \xi_1 \rangle - \langle \dot{\nu}_1, \xi_2 \rangle + \langle \nu + \mu, [\xi_1, \xi_2] \rangle \\ &\quad + t \langle \dot{\nu}_2, M(\lambda) \dot{\lambda}_1 \rangle - t \langle \dot{\nu}_1, M(\lambda) \dot{\lambda}_2 \rangle \\ &\quad + t \langle \nu + \mu, -[M(\lambda) \cdot \dot{\lambda}_1, M(\lambda) \cdot \dot{\lambda}_2] \rangle - (1-t) \langle \mu, [\dot{\lambda}_1, \dot{\lambda}_2] \rangle. \end{aligned}$$

Also, the expression $\frac{\partial\theta_t}{\partial t} = \theta_1 - \theta_0$ can be written as

$$(33) \quad \frac{\partial\theta_t}{\partial t}(g, \nu, \lambda)(\xi, \dot{\nu}, \dot{\lambda}) = \langle \nu + \mu, M(\lambda) \cdot \dot{\lambda} \rangle - \langle \mu, \frac{1}{2} \text{ad}_\lambda \dot{\lambda} + \dot{\lambda} \rangle.$$

From now on we will assume that $\dim \mathfrak{q} = 2$. Note that the one-form

$$\begin{aligned} & \omega_t(g, \nu, \lambda)(0, 0, \lambda)(\xi_2, \dot{\nu}_2, \dot{\lambda}_2) \\ &= t\langle \dot{\nu}_2, \lambda \rangle + t\langle \nu + \mu, -[\lambda, M(\lambda) \cdot \dot{\lambda}_2] \rangle - (1 - t)\langle \mu, [\lambda, \dot{\lambda}_2] \rangle \\ &= t\langle \nu + \mu, -M(\lambda)[\lambda, \dot{\lambda}_2] \rangle - (1 - t)\langle \mu, [\lambda, \dot{\lambda}_2] \rangle \end{aligned}$$

and (33) have the same kernel $\mathfrak{g} \oplus \mathfrak{g}_\mu^* \oplus \mathbb{R} \cdot \lambda \subset T_{(g, \nu, \lambda)}(G \times \mathfrak{g}_\mu^* \times \mathfrak{q})$. Therefore, there is a real-valued function f such that

$$\omega_t(g, \nu, \lambda)(0, 0, f(\nu, \lambda, t)\lambda)(\xi_2, \dot{\nu}_2, \dot{\lambda}_2) = \frac{\partial \theta_t}{\partial t}(g, \nu, \lambda)(\xi_2, \dot{\nu}_2, \dot{\lambda}_2).$$

That is, $X_t(g, \nu, \lambda) = f(\nu, \lambda, t) \frac{\partial}{\partial \lambda}$ and in particular

$$\Psi_t(g, \nu, \lambda) = (g, \nu, m_t(\nu, \lambda)\lambda)$$

for certain scaling factor $m_t : \mathfrak{g}_\mu^* \times \mathfrak{q} \rightarrow \mathbb{R}$. We will obtain an equation that fully determines m_1 and therefore the map Ψ_1 and the simple G -tube.

Taking the time-derivative of the time-dependent pull-back (see (15))

$$\begin{aligned} \frac{\partial}{\partial t} (\Psi_t^* \theta_t) &= \Psi_t^* \left((\mathbf{d}i_{X_t} + i_{X_t} \mathbf{d}) \theta_t + \frac{\partial \theta_t}{\partial t} \right) \\ &= \Psi_t^* (\mathbf{d}i_{X_t} \theta_t) = \Psi_t^* (\mathbf{d}i_{X_t} (\theta_0 + t(\theta_1 - \theta_0))) \\ &= \Psi_t^* (\mathbf{d}i_{X_t} \theta_0) = \mathbf{d}(\Psi_t^* (i_{X_t} \theta_0)). \end{aligned}$$

Additionally

$$\begin{aligned} \Psi_t^* (i_{X_t} \theta_0) &= \Psi_t^* (\langle \mu, \lambda \rangle f(\nu, \lambda, t)) \\ &= \langle \mu, \lambda \rangle f(\nu, m_t(\nu, \lambda)\lambda, t) = \frac{\partial}{\partial t} \langle \mu, m_t(\nu, \lambda)\lambda \rangle \end{aligned}$$

from where we get

$$\frac{\partial}{\partial t} (\Psi_t^* \theta_t - \mathbf{d} \langle \mu, m_t(\nu, \lambda)\lambda \rangle) = 0.$$

This equation implies that Ψ_t satisfies the following equation on one-forms

$$(34) \quad \Psi_1^* \theta_1 - \mathbf{d} \langle \mu, m_1(\nu, \lambda)\lambda \rangle = \theta_0 - \mathbf{d} \langle \mu, \lambda \rangle.$$

But this equation does not depend on the derivatives of the scaling factor m_1 because

$$\begin{aligned} \Psi_1^* \theta_1(\xi, \dot{\nu}, \dot{\lambda}) &= \langle \mu + \nu, (\mathbf{D}_\nu m_1 \cdot \dot{\nu} + \mathbf{D}_\lambda m_1 \cdot \dot{\lambda}) \lambda + M(m_1 \lambda) \cdot (m_1 \dot{\lambda}) \rangle \\ &= \langle \mu, \lambda \rangle (\mathbf{D}_\nu m_1 \cdot \dot{\nu} + \mathbf{D}_\lambda m_1 \cdot \dot{\lambda}) + \langle \mu + \nu, M(m_1 \lambda) \cdot (m_1 \dot{\lambda}) \rangle \end{aligned}$$

and

$$\mathbf{d}\langle \mu, m_1 \lambda \rangle(\xi, \dot{\nu}, \dot{\lambda}) = \langle \mu, \lambda \rangle (\mathbf{D}_\nu m_1 \cdot \dot{\nu} + \mathbf{D}_\lambda m_1 \cdot \dot{\lambda}) + \langle \mu, m_1 \dot{\lambda} \rangle.$$

Since $\langle \mu, [\lambda, \dot{\lambda}] \rangle$ is a non-vanishing one-form on the two dimensional space \mathfrak{q} with kernel $\mathbb{R} \cdot \lambda$ and $\langle \mu + \nu, M(\lambda) \cdot \dot{\lambda} \rangle - \langle \mu, \dot{\lambda} \rangle$ has also kernel $\mathbb{R} \cdot \lambda$ then there is real-valued function $h(\lambda, \nu)$ such that

$$\langle \mu + \nu, M(\lambda) \cdot \dot{\lambda} \rangle - \langle \mu, \dot{\lambda} \rangle = h(\lambda, \nu) \langle \mu, [\lambda, \dot{\lambda}] \rangle.$$

With this notation, (34) becomes the non-linear equation

$$(35) \quad h(m_1 \lambda, \nu) m_1^2 = \frac{1}{2}.$$

To find the solution $m_1(\lambda, \nu)$ it is enough to write down explicitly the simple G -tube

$$\Theta(g, \nu, \lambda) = (g \exp(m_1(\lambda, \nu) \lambda), \text{Ad}_{\exp(m_1(\lambda, \nu) \lambda)}(\nu + \mu)).$$

In the following lemmas we will see that under some algebraic assumptions on \mathfrak{g} we can write down m_1 in terms of elementary functions.

Lemma 7.1. *Assume that the subspace \mathfrak{q} defined by the splitting $\mathfrak{g} = \mathfrak{g}_\mu \oplus \mathfrak{q}$ is a 2-dimensional subalgebra. Then the equation (35) has the solution $m_1(\nu, \lambda) = \mathcal{E}(-\text{tr}(\text{ad}_\lambda|_{\mathfrak{q}}))$ where $\mathcal{E} : \mathbb{R} \rightarrow \mathbb{R}^+$ the unique analytic function that satisfies*

$$(36) \quad e^{-x\mathcal{E}(x)} = 1 - x\mathcal{E}(x) + \frac{x^2}{2}.$$

Proof. As the dimension of \mathfrak{q} is two and ad_ξ is singular it follows that $\text{ad}_\eta^2|_{\mathfrak{q}} - \text{tr}(\text{ad}_\eta|_{\mathfrak{q}})\text{ad}_\eta|_{\mathfrak{q}} = 0$ for any $\eta \in \mathfrak{q}$. Therefore,

$$\sum_{k \geq 0} \frac{1}{(k+1)!} \text{ad}_\xi^k = \text{Id} + \sum_{k \geq 0} \frac{(\text{tr}(\text{ad}_\xi|_{\mathfrak{q}}))^k}{(k+2)!} \text{ad}_\xi = \text{Id} + \frac{e^x - x - 1}{x^2} \text{ad}_\xi$$

where $x = \text{tr}(\text{ad}_\xi|_{\mathfrak{q}})$. Then (34) becomes

$$\langle \mu + \nu, M(\lambda) \cdot \dot{\lambda} \rangle = \langle \mu, M(\lambda) \cdot \lambda \rangle = \langle \mu, \dot{\lambda} \rangle + \langle \mu, [\lambda, \dot{\lambda}] \rangle \frac{e^x - x - 1}{x^2}.$$

Comparing with (35) we see that $h(\nu, \lambda) = \frac{e^{-x} + x - 1}{x^2}$ with $x = -\text{tr}(\text{ad}_\xi|_{\mathfrak{q}})$. \square

Remark 7.2. The function \mathcal{E} can be written in terms of the *Lambert W function* (see [4])

$$\mathcal{E}(x) = \begin{cases} \frac{x}{2} + \frac{W_0(-\exp(-1 - \frac{1}{2}x^2)) + 1}{x} & \text{if } x > 0 \\ \frac{x}{2} + \frac{W_{-1}(-\exp(-1 - \frac{1}{2}x^2)) + 1}{x} & \text{if } x < 0 \end{cases}$$

where W_0 and W_{-1} are defined in the same reference. It can be checked that $\mathcal{E}(x)$ is positive and strictly increasing for all $x \in \mathbb{R}$. Additionally $\mathcal{E}(x)$ is asymptotic to $\frac{x}{2}$ as $x \rightarrow \infty$, and satisfies $\mathcal{E}(0) = 1$ and $\mathcal{E}(x) \rightarrow 0$ as $x \rightarrow -\infty$.

Lemma 7.3. *Assume that the splitting $\mathfrak{g} = \mathfrak{g}_\mu \oplus \mathfrak{q}$ satisfies*

- 1) $\text{ad}_\xi^3 + a(\xi)\text{ad}_\xi = 0 \ \forall \xi \in \mathfrak{q}$ for a certain smooth function $a : \mathfrak{q} \rightarrow \mathbb{R}$, and
- 2) $\langle \mu + \nu, \text{ad}_\xi^2 \eta \rangle = 0$ for any $\xi, \eta \in \mathfrak{q}$ and $\nu \in \mathfrak{q}^\circ$.

In addition, let $b : (\mathfrak{g}_\mu^)_r \rightarrow \mathbb{R}$ be the function satisfying*

$$\langle \nu + \mu, [\xi, \eta] \rangle = b(\nu) \langle \mu, [\xi, \eta] \rangle$$

for any $\xi, \eta \in \mathfrak{q}$. Then equation (35) has the solution

$$m_1(\nu, \lambda) = \mathcal{F} \left(\frac{a(\lambda)}{4b(\nu)} \right) \frac{1}{\sqrt{b(\lambda)}},$$

where $\mathcal{F} : (-\infty, 1) \rightarrow \mathbb{R}^+$ is the analytic function

$$\mathcal{F}(x) = \begin{cases} \frac{\arcsin(\sqrt{x})}{\sqrt{x}} & \text{if } x > 0 \\ \frac{\operatorname{arcsinh}(\sqrt{|x|})}{\sqrt{|x|}} & \text{if } x < 0 \end{cases}$$

Proof. Using the first hypothesis

$$\sum_{n \geq 0} \frac{1}{(n+1)!} \operatorname{ad}_\xi^n = \operatorname{Id} + A_1(a(\xi))\operatorname{ad}_\xi + A_2(a(\xi))\operatorname{ad}_\xi^2$$

where A_1 and A_2 are analytic scalar functions. Then

$$\begin{aligned} \langle \mu + \nu, M(\lambda) \cdot \dot{\lambda} \rangle &= \langle \mu + \nu, \dot{\lambda} \rangle + A_1(a(\lambda))\langle \mu + \nu, [\lambda, \dot{\lambda}] \rangle \\ &= \langle \mu, \dot{\lambda} \rangle + A_1(a(\lambda))b(\nu)\langle \mu, [\lambda, \dot{\lambda}] \rangle, \end{aligned}$$

that is $h(\lambda, \nu) = A_1(a(\lambda))b(\nu)$. It can be checked that $A_1(x) = \frac{1 - \cos(\sqrt{x})}{x}$. If we assume $a(\lambda) > 0$ then using simple formal manipulations (35) is equivalent to $m_1\sqrt{a(\lambda)} = \arccos\left(1 - \frac{a(\lambda)}{2b(\nu)}\right)$ and as $2 \arcsin x = \arccos(1 - 2x^2)$ then $m_1(\lambda, \nu) = \mathcal{F}\left(\frac{a(\lambda)}{4b(\nu)}\right) \frac{1}{\sqrt{b(\lambda)}}$. If $a(\lambda) \leq 0$ a similar computation gives the same result. \square

7.1. $SO(3)$ simple tube

Under the hat map the Lie algebra $\mathfrak{g} = \mathfrak{so}(3)$ can be identified with \mathbb{R}^3 equipped with the cross product. The standard inner product $\langle \cdot, \cdot \rangle$ on $\mathbb{R}^3 \cong \mathfrak{g}$ will correspond to the dual pairing between \mathfrak{g} and \mathfrak{g}^* , identifying them.

Fix an element $\mu \in \mathfrak{g}^*$. We have two different possibilities:

- $\mu = 0$. In this case the G -tube is trivial (see Remark 4.4).
- $\mu \neq 0$. In this case \mathfrak{g}_μ is the subspace generated by μ and we will define \mathfrak{q} as the orthogonal complement to \mathfrak{g}_μ . The subspace \mathfrak{g}_μ^* being the annihilator of \mathfrak{q} is also identified with the subspace generated by μ .

The vector identity $\mathbf{a} \times (\mathbf{a} \times \mathbf{c}) = \langle \mathbf{a}, \mathbf{c} \rangle \mathbf{a} - \langle \mathbf{a}, \mathbf{a} \rangle \mathbf{c}$ implies that both conditions of Lemma 7.3 hold for $\mathfrak{so}(3)$ with $a(\lambda) = \|\lambda\|^2$. Therefore

the map

$$(37) \quad \begin{aligned} G \times \mathfrak{g}_\mu^* \times \mathfrak{q} &\longrightarrow SO(3) \times \mathbb{R}^3 \cong T^*SO(3) \\ (g, \nu, \lambda) &\longmapsto (gE(\nu, \lambda), E(\nu, \lambda) \cdot (\nu + \mu)) \end{aligned}$$

with $E(\nu, \lambda) = \exp\left(2 \frac{\arcsin\left(\frac{\sqrt{\frac{\mu}{\mu+\nu}} \|\lambda\|}{2}\right) \widehat{\lambda}}{\|\lambda\|}\right)$ is a simple $SO(3)$ -tube at $(e, \mu) \in T^*SO(3)$.

Note that this expression is exactly the same as the one obtained in Theorem 3 of [29]. In fact, this map was known in celestial mechanics as *regularized Serret-Andoyer-Deprit coordinates* (see [3] and references therein).

7.2. $SL(2, \mathbb{R})$ simple tube

On the Lie algebra $\mathfrak{g} = \mathfrak{sl}(2, \mathbb{R})$ the bilinear form $\langle A, B \rangle = -2\text{Tr}(AB)$ is non-degenerate and we will use it to identify \mathfrak{g} and \mathfrak{g}^* . If $\xi, \eta \in \mathfrak{sl}(2, \mathbb{R})$ it can be checked that $\text{ad}_\xi \text{ad}_\eta = \langle \xi, \eta \rangle \xi - \langle \xi, \xi \rangle \eta$ and then for any $\xi \in \mathfrak{g}$ we have $\text{ad}_\xi^3 + \|\xi\|^2 \text{ad}_\xi = 0$.

Fix an element $\mu \in \mathfrak{g}^*$. We now have three different cases:

- $\mu = 0$. In this case the G -tube is trivial (see Remark 4.4)
- $\|\mu\|^2 := \langle \mu, \mu \rangle \neq 0$. Then \mathfrak{g}_μ is one dimensional and is the space generated by μ . We will define \mathfrak{q} to be the orthogonal space to μ with respect to the pairing. Since the norm of μ is non-zero $\mathfrak{g} = \mathfrak{g}_\mu \oplus \mathfrak{q}$. As before $\mathfrak{g}_\mu^* = \mathfrak{q}^\circ = \mathfrak{g}_\mu$ so we can apply Lemma 7.3 obtaining that

$$\begin{aligned} G \times \mathfrak{g}_\mu^* \times \mathfrak{q} &\longrightarrow T^*SL(2, \mathbb{R}) \\ (g, \nu, \lambda) &\longmapsto (gE(\nu, \lambda), \text{Ad}_{E(\nu, \lambda)}^*(\nu + \mu)) \end{aligned}$$

with $E(\nu, \lambda) = \exp\left(\mathcal{F}\left(\frac{\|\lambda\|^2 \mu}{4(\mu+\nu)}\right) \sqrt{\frac{\mu}{\mu+\nu}} \lambda\right)$ is a simple $SL(2, \mathbb{R})$ -tube at $(e, \mu) \in T^*SL(2, \mathbb{R})$.

- $\|\mu\|^2 = 0$ and $\mu \neq 0$. In this case, using basic linear algebra, it can be shown that there is $k \in SL(2, \mathbb{R})$ such that $\mu = k \begin{bmatrix} 0 & s \\ 0 & 0 \end{bmatrix} k^{-1}$ with $s = 1$ or $s = -1$.

Also in this case \mathfrak{g}_μ is the subspace generated by μ , and we will define \mathfrak{q} as the subspace generated by $k \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} k^{-1}$ and $k \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix} k^{-1}$.

A generic element in \mathfrak{q} will be represented as $k \begin{bmatrix} a & 0 \\ b & -a \end{bmatrix} k^{-1}$. It can be checked that $\mathfrak{g}_\mu^* = \mathfrak{q}^\circ$ is the subspace generated by $k \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix} k^{-1}$.

A simple computation shows that \mathfrak{q} is a subalgebra of \mathfrak{g} so we can apply Lemma 7.1, obtaining that the map

$$G \times \mathfrak{g}_\mu^* \times \mathfrak{q} \longrightarrow T^*SL(2, \mathbb{R})$$

$$(g, \nu\mu, \lambda) \longmapsto \left(gE(\lambda), \text{Ad}_{E(\lambda)}^*((\nu + 1)\mu) \right),$$

where $\lambda = k \begin{bmatrix} a & 0 \\ b & -a \end{bmatrix} k^{-1}$, $\nu \in \mathbb{R}$ and $E(\lambda) = \exp(\mathcal{E}(2a)\lambda)$, is a simple $SL(2, \mathbb{R})$ -tube at $(e, \mu) \in T^*SL(2, \mathbb{R})$. Note that for this tube the domain is the whole space $G \times \mathfrak{g}_\mu^* \times \mathfrak{q}$. There are no restrictions on ν or λ but the map is not onto.

7.3. A $SO(3)$ restricted tube

Let H be a compact non-discrete proper subgroup of $SO(3)$. Note that H must be one-dimensional. We will denote by $\xi_{\mathfrak{h}} \in \mathbb{R}^3$ the generator of \mathfrak{h} with unit norm. In this setting the adapted splitting of Proposition 3.1 reduces to $\mathfrak{l} = \mathbb{R} \cdot \xi_{\mathfrak{h}}$, $\mathfrak{p} = \mathbb{R} \cdot \mu$ and $\mathfrak{n} = \mathbb{R} \cdot \xi_{\mathfrak{h}} \times \mu$. To obtain the restricted tube we will use (19), so we need to find $\zeta \in \mathfrak{n}$ satisfying the condition

$$(38) \quad \mathbf{J}_R(\Theta(g, \nu, \zeta))|_{\mathfrak{l}} = -\varepsilon$$

as a function of ν and ε . Using the notation of (16), Θ can be written as

$$\Theta(g, \nu, \lambda) = (gE(\nu, \lambda), \text{Ad}_{E(\nu, \lambda)}^*(\nu + \mu)).$$

Using Proposition 4.1 we can rewrite (38) as

$$(39) \quad \text{Ad}_{E(\nu, \zeta)}^*(\nu + \mu)|_{\mathfrak{l}} = \varepsilon.$$

Applying the explicit expression (37) for the $SO(3)$ simple tube we have $E(\nu, \zeta) = \exp(\rho(\nu, \zeta) \frac{\xi_{\mathfrak{h}} \times \mu}{\|\xi_{\mathfrak{h}} \times \mu\|})$. Then, solving (39) is equivalent to finding the real parameter ρ as a function of ν and ε that satisfies

$$\left\langle \exp \left(-\rho \frac{\xi_{\mathfrak{h}} \times \mu}{\|\xi_{\mathfrak{h}} \times \mu\|} \right) \cdot (\nu + \mu), \xi_{\mathfrak{h}} \right\rangle = \langle \varepsilon, \xi_{\mathfrak{h}} \rangle.$$

Since $\{\xi_{\mathfrak{h}}, \frac{\mu}{\|\mu\|}, \frac{\xi_{\mathfrak{h}} \times \mu}{\|\xi_{\mathfrak{h}} \times \mu\|}\}$ is an orthogonal basis this last equation is equivalent to

$$\langle \sin(\rho)(\nu + \mu), \frac{\mu}{\|\mu\|} \rangle = \langle \varepsilon, \xi_{\mathfrak{h}} \rangle.$$

Therefore, if we denote by r the expression $\arcsin \frac{(\varepsilon \cdot \xi_{\mathfrak{h}})\|\mu\|}{(\nu + \mu) \cdot \mu}$, the equation

$$(40) \quad \Phi(g, \nu; \varepsilon) = \left(g \exp \left(r \frac{\xi_{\mathfrak{h}} \times \mu}{\|\xi_{\mathfrak{h}} \times \mu\|} \right), \exp \left(-r \frac{\xi_{\mathfrak{h}} \times \mu}{\|\xi_{\mathfrak{h}} \times \mu\|} \right) \cdot (\nu + \mu) \right) \in SO(3) \times \mathbb{R}^3$$

defines a a restricted $SO(3)$ -tube.

With (40) and (37) we could apply Theorem 5.6 to obtain an explicit Hamiltonian tube for any cotangent lifted $SO(3)$ -action, however in the next example we will compute an explicit Hamiltonian tube not only for $SO(3)$ acting on $T^*\mathbb{R}^3$ but for $SO(n)$ acting on $T^*\mathbb{R}^n$.

7.4. Hamiltonian tube for $SO(n)$ acting on $T^*\mathbb{R}^n$

Consider the natural action of $SO(n)$ on \mathbb{R}^n and fix a point $z = (q, p) \in T^*\mathbb{R}^n$. Although we don't have an explicit expression even for a simple $SO(n)$ -tube we will see that we don't need it to compute an explicit tube for $T^*\mathbb{R}^n$. We can identify the Lie algebra of the orthogonal group $SO(n)$ with the second skew-symmetric power $\Lambda^2(\mathbb{R}^n)$, moreover the euclidean metric of \mathbb{R}^n induces an identification of $\mathfrak{so}(n)$ with its dual.

If $\mu = \mathbf{J}(z) = q \wedge p = 0$ then [28] provides an explicit computation of the Hamiltonian tube centered at (q, p) . Therefore, we will assume $\mu = q \wedge p \neq 0$, and in particular $q \neq 0$ so the isotropy $H := G_q = SO(\langle q \rangle_{\mathbb{R}}^{\perp})$ is the group of rotations of the hyperplane $\langle q \rangle_{\mathbb{R}}^{\perp}$. The linear slice $S = (\mathfrak{g} \cdot q)^{\perp}$ is the subspace generated by q , and note that this subspace is fixed by H and in this setting $\alpha = z|_S = \frac{p \cdot q}{\|q\|^2} q \in S^*$. Note also that $G_z = H_{\mu} = SO(\langle p, q \rangle_{\mathbb{R}}^{\perp}) \cong SO(n - 2)$.

The linear splitting of Proposition 3.1 becomes

$$\begin{aligned} \mathfrak{g} = & \underbrace{\{v \wedge w \mid v, w \in \langle q, p \rangle_{\mathbb{R}}^{\perp}\}}_{\mathfrak{h}_{\mu}} \oplus \underbrace{\mathbb{R} \cdot \mu}_{\mathfrak{p}} \oplus \underbrace{\{p \wedge v \mid v \in \langle q, p \rangle_{\mathbb{R}}^{\perp}\}}_{\mathfrak{t}} \\ & \oplus \underbrace{\{q \wedge v \mid v \in \langle q, p \rangle_{\mathbb{R}}^{\perp}\}}_{\mathfrak{n}}. \end{aligned}$$

Let $S_r := \{a \in \mathbb{R} \cdot q \mid \|a\| < \|q\|\} \subset S$ and $\mathfrak{p}_r^* := \{\nu \in \mathbb{R} \cdot \mu \mid \|\nu\| < \|\mu\|\}$. If Φ is a restricted tube adapted to the previous splitting the map (26) given by Theorem 5.6 becomes

$$\begin{aligned} \mathcal{T} : G \times_{H_\mu} (\mathfrak{p}_r^* \times S_r \times S^*) &\longrightarrow T^*(G \times_H S) \\ [g, \nu, a, b]_{H_\mu} &\longmapsto \varphi(\Phi(g, \nu, 0, 0); a, b + \alpha) \end{aligned}$$

because H fixes S and $\mathfrak{s} = 0$. But using (20) $\Phi(g, \nu, 0, 0) = (g, \nu + \mu)$. It can be easily checked that the map $\mathfrak{s} : G \times_H S_r \rightarrow \mathbb{R}^n$ defined by $[g, a]_H \mapsto g \cdot (q + a)$ is a Palais' tube around q . Using this tube, after some easy manipulations the previous Hamiltonian tube at (q, p) can be written as

$$\begin{aligned} \mathcal{T} : SO(n) \times_{H_\mu} (\mathfrak{p}_r^* \times S_r \times S^*) &\longrightarrow T^*\mathbb{R}^n \\ [g, \nu, a, b]_{H_\mu} &\longmapsto \left(g \cdot (q + a), g \cdot \left(\widehat{\nu + \mu} \frac{q + a}{\|q + a\|^2} + b + \frac{q \cdot p}{\|q\|^2} q \right) \right) \end{aligned}$$

where $\widehat{\xi}v$ represents the Lie algebra action of $\xi \in \mathfrak{g} \cong \mathfrak{g}^*$ at the point $v \in \mathbb{R}^n$. Alternatively, this map can be written as

$$\mathcal{T}([g, \nu, a, b]_{H_\mu}) = \left(g \cdot (q + a), g \cdot \left(\frac{\mu + \nu}{\mu} \frac{q}{q + a} \left(p - \frac{q \cdot p}{\|q\|^2} q \right) + b + \frac{q \cdot p}{\|q\|^2} q \right) \right)$$

where $\frac{\mu + \nu}{\mu}$ is the real scalar ρ such that $\rho\mu = \mu + \nu$ and similarly for $\frac{q}{q+a}$.

The same ideas and computations can be used to obtain a cotangent Hamiltonian tube for the cotangent bundle of the sphere T^*S^{n-1} endowed with the natural $SO(n)$ action. More specially if we identify T^*S^{n-1} with the subset $T^*S^{n-1} = \{(q, p) \mid \|q\| = 1, \langle q, p \rangle = 0\} \subset T^*\mathbb{R}^n \cong \mathbb{R}^n \times \mathbb{R}^n$ and we fix a point $(q, p) \in T^*S^{n-1}$ then

$$\begin{aligned} \mathcal{T} : SO(n) \times_{H_\mu} \mathfrak{p}_r^* &\longrightarrow T^*S^{n-1} \\ [g, \nu]_{H_\mu} &\longmapsto \left(g \cdot q, g \cdot \left(\frac{\mu + \nu}{\mu} p \right) \right) \end{aligned}$$

is a Hamiltonian tube at $(q, p) \in T^*S^{n-1}$.

References

- [1] R. Abraham and J. Marsden, *Foundations of mechanics*, Benjamin-Cummings, New York, 1978. Second edition.
- [2] L. Bates and E. Lerman, *Proper group actions and symplectic stratified spaces*, Pacific J. Math. **181** (1997), no. 2, 201–229.

- [3] G. Benettin, F. Fassò, and M. Guzzo, *Long term stability of proper rotations and local chaotic motions in the perturbed Euler rigid body*, Regul. Chaotic Dyn. **11** (2006), no. 1, 1–17.
- [4] R. M. Corless, G. H. Gonnet, D. E. G. Hare, D. J. Jeffrey, and D. E. Knuth, *On the Lambert W function*, Adv. Comput. Math. **5** (1996), no. 4, 329–359.
- [5] J. J. Duistermaat and J. A. Kolk, *Lie groups*, Springer, Berlin, 2000.
- [6] V. Guillemin and S. Sternberg, *A normal form for the moment map*, in: Differential geometric methods in mathematical physics (Jerusalem, 1982), volume 6 of Math. Phys. Stud., pages 161–175. Reidel, Dordrecht, 1984.
- [7] D. D. Holm, J. E. Marsden, and T. S. Ratiu, *The Euler-Poincaré equations and semidirect products with applications to continuum theories*, Adv. Math. **137** (1998), no. 1, 1–81.
- [8] E. Lerman and S. F. Singer, *Stability and persistence of relative equilibria at singular values of the moment map*, Nonlinearity **11** (1998), no. 6, 1637–1649.
- [9] E. Lerman and C. Willett, *The topological structure of contact and symplectic quotients*, Internat. Math. Res. Notices (2001), no. 1, 33–52.
- [10] D. Lewis, *Lagrangian block diagonalization*, J. Dynam. Differential Equations **4** (1992), no. 1, 1–41.
- [11] C.-M. Marle, *Modèle d'action hamiltonienne d'un groupe de Lie sur une variété symplectique*, Rend. Sem. Mat. Univ. Politec. Torino **43** (1985), no. 2, 227–251.
- [12] J. Marsden and A. Weinstein, *Reduction of symplectic manifolds with symmetry*, Rep. Mathematical Phys. **5** (1974), no. 1, 121–130.
- [13] J. Montaldi, *Persistence and stability of relative equilibria*, Nonlinearity **10** (1997), 449–466.
- [14] J. Montaldi and R. M. Roberts, *Relative equilibria of molecules*, J. Nonlinear Sci. **9** (1999), no. 1, 53–88.
- [15] J. Montaldi, R. M. Roberts, and I. N. Stewart, *Periodic solutions near equilibria of symmetric Hamiltonian systems*, Philos. Trans. Roy. Soc. London Ser. A **325** (1988), no. 1584, 237–293.

- [16] J. Montaldi and T. Tokieda, *Openness of momentum maps and persistence of extremal relative equilibria*, *Topology* **42** (2003), no. 4, 833–844.
- [17] J.-P. Ortega and T. S. Ratiu, *Persistence and smoothness of critical relative elements in Hamiltonian systems with symmetry*, *C. R. Acad. Sci. Paris Sér. I Math.* **325** (1997), no. 10, 1107–1111.
- [18] J.-P. Ortega and T. S. Ratiu, *Stability of Hamiltonian relative equilibria*, *Nonlinearity* **12** (1999), no. 3, 693–720.
- [19] J.-P. Ortega and T. S. Ratiu, *Momentum maps and Hamiltonian reduction*, Birkhäuser, 2004.
- [20] R. S. Palais, *On the existence of slices for actions of non-compact Lie groups*, *Ann. of Math.* **73** (1961), 295–323.
- [21] G. Patrick, *Relative equilibria of hamiltonian systems with symmetry: linearization, smoothness, and drift*, *J. Nonlinear Sci.* **5** (1995), no. 5, 373–418.
- [22] M. Perlmutter, M. Rodríguez-Olmos, and M. E. Sousa-Dias, *On the geometry of reduced cotangent bundles at zero momentum*, *J. Geom. Phys.* **57** (2007), no. 2, 571–596.
- [23] M. Perlmutter, M. Rodríguez-Olmos, and M. E. Sousa-Dias, *The symplectic normal space of a cotangent-lifted action*, *Differential Geom. Appl.* **26** (2008), no. 3, 277–297.
- [24] M. Roberts, C. Wulff, and J. S. W. Lamb, *Hamiltonian systems near relative equilibria*, *J. Differential Equations* **179** (2002), no. 2, 562–604.
- [25] R. M. Roberts and M. E. Sousa-Dias, *Bifurcations from relative equilibria of Hamiltonian systems*, *Nonlinearity* **10** (1997), no. 6, 1719–1738.
- [26] W. J. Satzer, Jr., *Canonical reduction of mechanical systems invariant under abelian group actions with an application to celestial mechanics*, *Indiana Univ. Math. J.* **26** (1977), no 5, 951–976.
- [27] T. Schmah, *Symmetries of cotangent bundles*. PhD thesis, Ecole Polytechnique Fdrale de Lausanne, 2001.
- [28] T. Schmah, *A cotangent bundle slice theorem*, *Differential Geom. Appl.* **25** (2007), no. 1, 101–124.
- [29] T. Schmah and C. Stoica, *Normal forms for Lie symmetric cotangent bundle systems with free and proper actions*, preprint, [arXiv:1311.7447](https://arxiv.org/abs/1311.7447), 2013.

- [30] J. Simo, D. Lewis, and J. Marsden, *Stability of relative equilibria. part i: The reduced energy-momentum method*, Arch. Ration. Mech. An. **115** (1991), no. 1, 15–59.
- [31] R. Sjamaar and E. Lerman, *Stratified symplectic spaces and reduction*, Ann. of Math. **134** (1991), no. 2, 375–422.

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