

Global well-posedness for the critical Schrödinger-Debye system

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ABSTRACT. We establish global well-posedness results for the initial value problem associated to the Schrödinger-Debye system in dimension two, for data in $H^s(\mathbb{R}^2) \times L^2(\mathbb{R}^2)$, $2/3 < s \leq 1$ and for data in $H^1(\mathbb{R}^2) \times H^1(\mathbb{R}^2)$.

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1. Introduction

We consider the initial value problem (IVP) for the Schrödinger-Debye system

$$(1.1) \quad \begin{cases} iu_t + \frac{1}{2}\Delta u = uv, & t > 0, & x \in \mathbb{R}^n, \\ \mu v_t + v = \lambda|u|^2, & \mu > 0, & \lambda = \pm 1, \\ u(0) = u_0, & v(0) = v_0, \end{cases}$$

where $u = u(x, t)$ is a complex-valued function, $v = v(x, t)$ is a real-valued function and Δ is the Laplacian operator in the spatial variable. This model describes the propagation of an electromagnetic wave through a non-resonant medium whose material response time is relevant. See Newel and Moloney [10] for a more complete

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discussion of this model. In the absence of the delay ($\mu = 0$), the system (1.1) reduces to the cubic nonlinear Schrödinger equation (NLS)

$$iu_t + \frac{1}{2}\Delta u = \lambda u|u|^2,$$

which is focusing or defocusing for $\lambda = -1$ and 1 , respectively. Similarly, the sign of the parameter λ provides an analogous classification for (1.1). For sufficiently regular data, the mass of the solution u of the system (1.1) is invariant. More precisely,

$$(1.2) \quad \int_{\mathbb{R}^n} |u(x, t)|^2 dx = \int_{\mathbb{R}^n} |u_0(x)|^2 dx.$$

The system (1.1) has the following pseudo-Hamiltonian structure

$$(1.3) \quad \frac{d}{dt} E(u(t), v(t)) = 2\lambda\mu \int_{\mathbb{R}^n} |v_t|^2 dx,$$

where

$$(1.4) \quad \begin{aligned} E(u(t), v(t)) &:= E(u, v) = \int_{\mathbb{R}^n} (|\nabla u|^2 + 2v|u|^2 - \lambda v^2) dx \\ &= \int_{\mathbb{R}^n} (|\nabla u|^2 + \lambda|u|^4 - \lambda\mu^2|v_t|^2) dx. \end{aligned}$$

Also, the system (1.1) is equivalent to the following integral form

$$(1.5) \quad u(t) = S(t)u_0 + i \int_0^t S(t - \tau)u(\tau)v(\tau) d\tau,$$

and

$$(1.6) \quad v(t) = e^{-\frac{t}{\mu}}v_0 + \frac{\lambda}{\mu} \int_0^t e^{-\frac{(t-\tau)}{\mu}} |u(\tau)|^2 d\tau,$$

where $S(t) = e^{it\Delta/2}$ is the unitary Schrödinger group.

The well-posedness of the system (1.1) has been studied by different authors. In 2000, Bidégaray ([4] and [5]) studied the local well-posedness (LWP) of the system (1.1) in dimensions $n = 1, 2, 3$ and for data in $H^s(\mathbb{R}^n) \times H^s(\mathbb{R}^n)$, $s > n/2$, $s = 0$ and $s = 1$. In 2004, Corcho and Linares ([7]) obtained the best LWP result in dimension one. Later, in 2009, Corcho and Matheus ([8]) obtained a refined LWP and global-well posedness (GWP) result in the framework of Bourgain spaces, also in dimension one. Recently, Corcho, Oliveira e Silva (see [9]), also in the framework of Bourgain spaces $X_{s,b}$ (see definition in Section 2), obtained the following LWP result in dimensions two and three.

THEOREM 1.1. *Let $n = 2, 3$. For any $(u_0, v_0) \in H^{s_1}(\mathbb{R}^n) \times H^{s_2}(\mathbb{R}^n)$, with s_1 and s_2 satisfying the conditions*

$$(1.7) \quad \max\{0, s_1 - 1\} \leq s_2 \leq \min\{2s_1, s_1 + 1\},$$

there exists a positive time

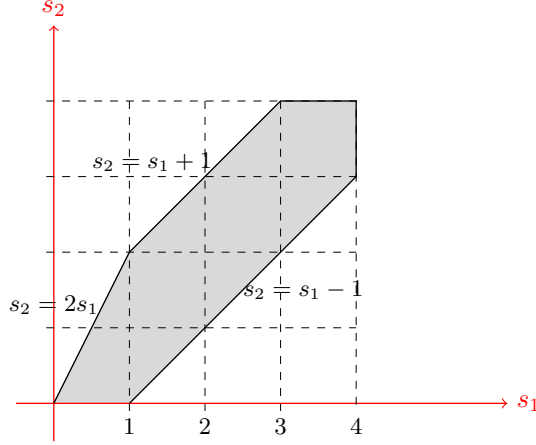
$$T = T(\|u_0\|_{H^{s_1}}, \|v_0\|_{H^{s_2}})$$

and a unique solution $(u(t), v(t))$ of the IVP (1.1) on the time interval $[0, T]$, such that

$$(i) \quad (\psi_T u(t), \psi_T v(t)) \in X_{s_1, b} \times H_{s_2, c};$$

(ii) $(u, v) \in C([0, T]; H^{s_1}(\mathbb{R}^n) \times H^{s_2}(\mathbb{R}^n))$
 for suitable b and c close to $1/2+$ (ψ_T denotes, as usual, a cutoff function for the time interval $[0, T]$). Moreover, the map $(u_0, v_0) \rightarrow (u(t), v(t))$ is locally Lipschitz from $H^{s_1}(\mathbb{R}^n) \times H^{s_2}(\mathbb{R}^n)$ into $C([0, T]; H^{s_1}(\mathbb{R}^n) \times H^{s_2}(\mathbb{R}^n))$.

In the following graphic we resume the LWP in the above theorem:



The shaded region represents the region where the LWP result exists. They also proved the following GWP result in dimension two.

THEOREM 1.2. *Let $(u_0, v_0) \in H^1(\mathbb{R}^2) \times L^2(\mathbb{R}^2)$. Then, for all $T > 0$, there exists a unique solution*

$$(u, v) \in C([0, T]; H^1(\mathbb{R}^2) \times L^2(\mathbb{R}^2))$$

to the initial value problem (1.1), such that

$$(1.8) \quad \|\nabla u(\cdot, t)\|_{L_x^2}^2 + \|v(\cdot, t)\|_{L_x^2}^2 \leq \alpha_0 e^{\alpha_1 T}, \quad t \in [0, T], \quad T > 0,$$

where $\alpha_0 = \alpha_0(E(u_0, v_0), \|v_0\|_{L_x^2}^2, \|u_0\|_{L_x^2}^2)$ and $\alpha_1 = \alpha_1(\|u_0\|_{L_x^2}^2)$.

In this work we prove the following GWP result in the space $H^1(\mathbb{R}^2) \times H^1(\mathbb{R}^2)$.

THEOREM 1.3. *Let $(u_0, v_0) \in H^1(\mathbb{R}^2) \times H^1(\mathbb{R}^2)$. Then, for all $T > 0$, there exists a unique solution*

$$(u, v) \in C([0, T]; H^1(\mathbb{R}^2) \times H^1(\mathbb{R}^2))$$

to the initial value problem (1.1).

This solves a problem left open in Corcho, Oliveira and Silva [9]; see Remark 4.3 therein. The proof of this theorem is simple; it uses basic properties of the unitary Schrödinger group and the Gronwell inequality.

Next we prove a result of GWP (small data) below energy space for the Schrödinger solution u . We will use the method of Bourgain ([3]) on high and low frequencies together with the framework of dispersive Sobolev spaces.

THEOREM 1.4. *Let $(u_0, v_0) \in H^s(\mathbb{R}^2) \times L^2(\mathbb{R}^2)$, with $2/3 < s \leq 1$, such that*

$$2c_0^2 \|u_0\|_{L_x^2} < 1,$$

where c_0 is the best constant in the Gagliardo-Nirenberg inequality

$$\|u\|_{L_x^4(\mathbb{R}^2)} \leq c_0 \|u\|_{L_x^2(\mathbb{R}^2)}^{1/2} \|u\|_{H_x^1(\mathbb{R}^2)}^{1/2}.$$

Then, for all $T > 0$, there exists a unique solution

$$(u, v) \in C([0, T]; H^s(\mathbb{R}^2) \times L^2(\mathbb{R}^2))$$

to the initial value problem (1.1) with $\lambda = -1$.

The difficulty in the proof of these theorems lies in the absence of conserved quantities in the energy space.

2. notation and basic properties

Let Y_x be a normed space on \mathbb{R}^n . We denote by $L_t^p Y_x(J \times \mathbb{R}^n) := L_t^p(J; Y_x(\mathbb{R}^n))$, where $J \subseteq \mathbb{R}$ is an interval, the completion of the space of Schwartz functions $f(x, t)$, with the norm

$$\|f\|_{L_t^p Y_x(J \times \mathbb{R}^n)} = \left(\int_J \|f(t)\|_{Y_x}^p dt \right)^{1/2}.$$

When \mathbb{R}^n and J are implicit, we denote this norm by $\|f\|_{L_t^p Y_x}$. Let h be a continuous function on \mathbb{R} . We define the space $X_{s,b}$ by

$$X_{s,b} = X_{s,b}^h(\mathbb{R}^n \times \mathbb{R}) := \overline{S_{x,t}(\mathbb{R}^n \times \mathbb{R})}^{\|\cdot\|_{X_{s,b}}},$$

where $S_{x,t}$ is the Schwartz space, with norm

$$\|u\|_{X_{s,b}^{h(\xi)}} := \|\langle \xi \rangle^s \langle \tau - h(\xi) \rangle^b \tilde{u}(\xi, \tau)\|_{L_\tau^2 L_x^2},$$

with the notation $\langle \cdot \rangle := \{1 + |\cdot|^2\}^{1/2} \sim 1 + |\cdot|$, recalling that the Fourier transform (space-time) is given by

$$\tilde{u}(\xi, \tau) = \int_{\mathbb{R}} \int_{\mathbb{R}^n} e^{-i(x \cdot \xi + t\tau)} u(x, t) dx dt.$$

The space $X_{s,b}$ is called Bourgain space or dispersive Sobolev space. Similarly we define the space $H_{l,c}$ with norm

$$\|u\|_{H_{l,c}} := \|\langle \xi \rangle^l \langle \tau \rangle^c \tilde{u}(\xi, \tau)\|_{L_\tau^2 L_x^2}.$$

Let $I \subset \mathbb{R}$ be any interval of time. We define

$$X_{s,b}(\mathbb{R}^n \times I) := \{u; u = v|_{\mathbb{R}^n \times I}, v \in X_{s,b}^h(\mathbb{R}^n \times \mathbb{R})\}$$

with norm

$$\|u\|_{X_{s,b}(\mathbb{R}^n \times I)} := \inf\{\|v\|_{X_{s,b}}; v|_{\mathbb{R}^n \times I} = u\}.$$

We have that

$$X_{s,0} = L_t^2 H_x^s \quad \text{and} \quad (X_{s,b}^{h(\xi)})' = X_{-s,-b}^{-h(-\xi)},$$

and moreover,

$$X_{s,b} \hookrightarrow C(\mathbb{R}; H^s(\mathbb{R}^n)) \quad \text{and} \quad H_{l,c} \hookrightarrow C(\mathbb{R}; H^l(\mathbb{R}^n)),$$

for any $b, c > 1/2$ and $s, l \in \mathbb{R}$.

PROPOSITION 2.1. *We have the following known embedding*

(1) if $2 \leq p < \infty$ and $b \geq \frac{1}{2} - \frac{1}{p}$ then $\|u\|_{L_t^p H_x^s} \leq c \|u\|_{X_{s,b}}$,

(2) if $2 \leq p, q < \infty$ and $b \geq \frac{1}{2} - \frac{1}{p}$ and $s \geq \frac{n}{2} - \frac{1}{q}$ then $\|u\|_{L_t^p L_x^q} \leq c \|u\|_{X_{s,b}}$,

(3) if $1 < p \leq 2$ and $b \leq \frac{1}{2} - \frac{1}{p}$ then $\|u\|_{X_{s,b}} \leq c \|u\|_{L_t^p H_x^s}$,

(4)

$$\|\mathcal{D}_x^{1/2^-}(u_1 u_2)\|_{L_t^2 L_x^2} \leq c \|u_1\|_{X_{1/2^-, 1/2^+}} \|u_2\|_{X_{0, 1/2^-}},$$

and

$$\|\mathcal{D}_x^{1/2^+}(u_1 u_2)\|_{L_t^2 L_x^2} \leq c \|u_1\|_{X_{1/2^+, 1/2^+}} \|u_2\|_{X_{0^+, 1/2^+}},$$

where $(\widehat{\mathcal{D}_x^s f})(\xi) = \langle \xi \rangle^s \widehat{f}(\xi)$.

PROOF. For the items (1), (2) and (3) see [11] and for the item (4) see [3]. \square

We define an operator \mathcal{L} by $\widehat{\mathcal{L}u}(\xi) = -ih(\xi)\widehat{u}(\xi)$.

PROPOSITION 2.2. Let $f \in H_x^s(\mathbb{R}^d)$ for some $s \in \mathbb{R}$ and for some polynomial $h : \mathbb{R}^d \rightarrow \mathbb{R}$. Then

i) For any Schwartz time cutoff $\eta \in S_x(\mathbb{R})$, we have

$$\|\eta(t)e^{t\mathcal{L}}f\|_{X_{s,b}^{h(\xi)}(\mathbb{R} \times \mathbb{R}^d)} \lesssim_{\eta,b} \|f\|_{H_x^s(\mathbb{R}^d)}.$$

ii) Let Y be a Banach space of functions on $\mathbb{R} \times \mathbb{R}^d$ with the property that

$$\|e^{it\tau_0}e^{t\mathcal{L}}f\|_Y \lesssim \|f\|_{H_x^s(\mathbb{R}^d)},$$

for all $f \in H_x^s(\mathbb{R}^d)$. Then we have the embedding

$$\|u\|_Y \lesssim_b \|u\|_{X_{s,b}^{h(\xi)}(\mathbb{R} \times \mathbb{R}^d)},$$

for $b > 1/2$.

PROOF. See [11]. \square

Now define the function $\vartheta : \mathbb{R} \rightarrow \mathbb{R}$ by

$$(2.1) \quad \vartheta(t) := \begin{cases} 1 & |t| < 1, \\ 0 & |t| \geq 2, \end{cases}$$

and $\vartheta \in C_0^\infty(\mathbb{R})$. Moreover, define $\vartheta_T(t) := \vartheta(\frac{t}{T})$.

Consider the IVP

$$(2.2) \quad \begin{cases} u_t - \mathcal{L}u = F, \\ u(0) = u_0. \end{cases}$$

The solution of this IVP is given by

$$(2.3) \quad u(t) = e^{it\mathcal{L}}u_0 + \int_0^t e^{i(t-s)\mathcal{L}}F(s)ds.$$

THEOREM 2.1. *Let u be the solution of the integral equation (2.3). If $s \in \mathbb{R}$ and $b > \frac{1}{2}$, then*

$$\|\vartheta_T u\|_{X_{s,b}^{h(\epsilon)}} \leq C \|u_0\|_{H_x^s} + C \|F\|_{X_{s,b-1}^{h(\epsilon)}}.$$

PROOF. See [11]. □

We establish the properties of the group $\{e^{it\Delta}\}_{t=-\infty}^\infty$ in the $L^p(\mathbb{R}^n)$ – space.

PROPOSITION 2.3. *If $t \neq 0$, $\frac{1}{p} + \frac{1}{p'} = 1$ and $p' \in [1, 2]$, then we have $e^{it\Delta} : L_x^{p'}(\mathbb{R}^n) \rightarrow L_x^p(\mathbb{R}^n)$ is continuous and*

$$(2.4) \quad \|e^{it\Delta} f\|_{L_x^p(\mathbb{R}^n)} \leq c|t|^{-n/2(1/p'-1/p)} \|f\|_{L_x^{p'}(\mathbb{R}^n)}.$$

For a proof of this proposition we refer to [11].

We proceed with the notion of admissible pair.

DEFINITION 2.1. *We say that the exponent pair (q, r) is admissible if*

$$\frac{2}{q} = n\left(\frac{1}{2} - \frac{1}{r}\right),$$

where

- $2 \leq r \leq \frac{2n}{n-2}$ if $n > 2$,
- $2 \leq r < \infty$ if $n = 2$,
- $2 \leq r \leq \infty$ if $n = 1$.

PROPOSITION 2.4 (Strichartz estimates). *If $n \geq 1$, $s \in \mathbb{R}$, (q_1, r_1) and (q_2, r_2) are admissible and $\frac{1}{q_2} + \frac{1}{q_2'} = 1$, $\frac{1}{r_2} + \frac{1}{r_2'} = 1$, then we have the homogeneous Strichartz estimate*

$$(2.5) \quad \|e^{it\Delta} u_0\|_{L_t^{q_1} L_x^{r_1}(\mathbb{R} \times \mathbb{R}^n)} \leq C(n, q, r) \|u_0\|_{L_x^2(\mathbb{R}^n)},$$

the dual homogeneous Strichartz estimate

$$(2.6) \quad \left\| \int_0^t e^{is\Delta} F(s) ds \right\|_{L_x^2(\mathbb{R}^n)} \leq C(n, q_2, r_2) \|F\|_{L_t^{q_2'} L_x^{r_2'}(\mathbb{R} \times \mathbb{R}^n)},$$

and the inhomogeneous Strichartz estimate

$$(2.7) \quad \left\| \int_0^t e^{i(t-s)\Delta} F(s) ds \right\|_{L_t^q L_x^r(\mathbb{R} \times \mathbb{R}^n)} \leq C(n, q_1, r_1, q_2, r_2) \|F\|_{L_t^{q_2'} L_x^{r_2'}(\mathbb{R} \times \mathbb{R}^n)}.$$

PROOF. See [11]. □

3. Global Well-posedness in $H^1(\mathbb{R}^2) \times H^1(\mathbb{R}^2)$

In this paper we give a negative answer to the question of the existence of blow-up solutions for the initial data in $H^1(\mathbb{R}^2) \times H^1(\mathbb{R}^2)$ in Corcho, Oliveira and Silva [9]; see Remark 4.3 therein.

In order to prove a global theory in $H^1(\mathbb{R}^2) \times H^1(\mathbb{R}^2)$ we need an estimate $\nabla v(\cdot)$. To achieve this, we apply the gradient in the equation (1.6) and we obtain

$$(3.1) \quad \nabla v(t) = \nabla v_0 e^{-\frac{t}{\mu}} + \frac{\lambda}{\mu} \int_0^t e^{-(t-\tau)/\mu} \nabla |u(\tau)|^2 d\tau.$$

Observe that,

$$\nabla|u|^2 = \nabla\{u.\bar{u}\} = 2\operatorname{Re}(u.\nabla\bar{u}).$$

Replacing this expression in (3.1), we get

$$(3.2) \quad \|\nabla v(t)\|_{L_x^2} \leq \|\nabla v_0\|_{L_x^2} + 2\left|\frac{\lambda}{\mu}\right| \int_0^t e^{-(t-\tau)/\mu} \|u(\tau).\nabla u(\tau)\|_{L_x^2} d\tau.$$

Therefore

$$(3.3) \quad \begin{aligned} \|\nabla v(t)\|_{L_x^2} &\leq \|\nabla v_0\|_{L_x^2} + 2\left|\frac{\lambda}{\mu}\right| \int_0^t e^{-(t-\tau)/\mu} \|u(\tau)\|_{L_x^4} \|\nabla u(\tau)\|_{L_x^4} d\tau \\ &\leq \|\nabla v_0\|_{L_x^2} + \sqrt{\frac{2}{\mu}} \|u(t)\|_{L_t^4 L_x^4} \|\nabla u(t)\|_{L_t^4 L_x^4}. \end{aligned}$$

Moreover, applying the gradient in the equation (1.5), we have

$$(3.4) \quad \nabla u = S(t)\nabla u_0 + i(I_1 + I_2),$$

where

$$I_1 = \int_0^t S(t-\tau)\nabla u(\tau).v(\tau) d\tau, \quad I_2 = \int_0^t S(t-\tau)u(\tau).\nabla v(\tau) d\tau.$$

Since (4, 4) is an admissible pair, we conclude that

$$(3.5) \quad \|\nabla u\|_{L_t^4 L_x^4} \leq c\|\nabla u_0\|_{L_x^2} + \|I_1\|_{L_t^4 L_x^4} + \|I_2\|_{L_t^4 L_x^4},$$

and using the inhomogeneous Strichartz estimate (2.7), we have

$$(3.6) \quad \|I_1\|_{L_t^4 L_x^4} \leq c\|\nabla u(\tau).v(\tau)\|_{L_t^{4/3} L_x^{4/3}} \leq c\|\nabla u\|_{L_t^2 L_x^2} \|v\|_{L_t^4 L_x^4}.$$

By immersion

$$(3.7) \quad \|v\|_{L_x^4} \leq c_0 \|v\|_{L_x^2}^{1/2} \|\nabla v\|_{L_x^2}^{1/2},$$

and from (3.6) and (3.7), we have

$$(3.8) \quad \begin{aligned} \|I_1\|_{L_t^4 L_x^4} &\leq c\|\nabla u\|_{L_t^2 L_x^2} \left(\int_0^t \|v\|_{L_x^2}^2 \|\nabla v\|_{L_x^2}^2 d\tau \right)^{1/4} \\ &\leq c\|\nabla u\|_{L_t^2 L_x^2} \|v\|_{L_t^\infty L_x^2}^{1/2} \left(\int_0^t \|\nabla v\|_{L_x^2}^2 d\tau \right)^{1/4} \\ &\leq c\|\nabla u\|_{L_t^2 L_x^2}^2 \|v\|_{L_t^\infty L_x^2} + \left(\int_0^t \|\nabla v\|_{L_x^2}^2 d\tau \right)^{1/2}. \end{aligned}$$

Now we estimate I_2 :

$$(3.9) \quad \begin{aligned} \|I_2\|_{L_t^4 L_x^4} &\leq \|u(\tau).\nabla v(\tau)\|_{L_t^{4/3} L_x^{4/3}} \\ &\leq c\|u_0\|_{L_x^2}^{1/2} \|\nabla u\|_{L_t^2 L_x^2}^{1/2} \left(\int_0^t \|\nabla v\|_{L_x^2}^2 d\tau \right)^{1/2}. \end{aligned}$$

Here we used (3.7) and the conserved quantity (1.2). Replacing (3.9) and (3.8) in (3.5), we get

$$(3.10) \quad \begin{aligned} \|\nabla u\|_{L_t^4 L_x^4} &\leq c\|\nabla u_0\|_{L_x^2} + c\|\nabla u\|_{L_t^2 L_x^2}^2 \|v\|_{L_t^\infty L_x^2} + \left(\int_0^t \|\nabla v\|_{L_x^2}^2 d\tau\right)^{1/2} \\ &\quad + c\|u_0\|_{L_x^2}^{1/2} \|\nabla u\|_{L_t^2 L_x^2}^{1/2} \left(\int_0^t \|\nabla v\|_{L_x^2}^2 d\tau\right)^{1/2}. \end{aligned}$$

Using (3.7), it follows that

$$(3.11) \quad \|u\|_{L_t^4 L_x^4} \leq T_{max}^{1/4} \|u\|_{L_t^\infty L_x^4} \leq c_0 T_{max}^{1/4} \|u_0\|_{L_x^2}^{1/2} \|\nabla u\|_{L_t^\infty L_x^2}^{1/2}.$$

Combining (3.3), (3.10) and (3.11), we can show that

$$(3.12) \quad \begin{aligned} \|\nabla v(t)\|_{L_x^2} &\leq \|\nabla v_0\|_{L_x^2} + \mathcal{K}_T \|\nabla u(t)\|_{L_t^4 L_x^4} \\ &\leq \|\nabla v_0\|_{L_x^2} + \mathcal{K}_T \left\{ \|\nabla u_0\|_{L_x^2} + \|\nabla u\|_{L_t^2 L_x^2}^2 \|v\|_{L_t^\infty L_x^2} + \left(\int_0^t \|\nabla v\|_{L_x^2}^2 d\tau\right)^{1/2} \right\}, \\ &\quad + \mathcal{K}_T \|u_0\|_{L_x^2}^{1/2} \|\nabla u\|_{L_T^\infty L_x^2}^{1/2} \left(\int_0^t \|\nabla v\|_{L_x^2}^2 d\tau\right)^{1/2} \end{aligned}$$

where

$$\mathcal{K}_T = c_0 \sqrt{\frac{2}{\mu}} T_{max}^{1/4} \|u_0\|_{L_x^2}^{1/2} \|\nabla u\|_{L_t^\infty L_x^2}^{1/2}.$$

Observe that

$$(3.13) \quad \|\nabla v(t)\|_{L_x^2} \leq \Phi + G(T) \left(\int_0^t \|\nabla v(\tau)\|_{L_x^2}^2 d\tau\right)^{1/2},$$

where

$$\begin{aligned} G(T) &= \mathcal{K}_T (1 + \|u_0\|_{L_x^2}^{1/2} \|\nabla u\|_{L_T^\infty L_x^2}^{1/2}) \\ \Phi &= \|\nabla v_0\|_{L_x^2} + \mathcal{K}_T \|\nabla u_0\|_{L_x^2} + \mathcal{K}_T \|\nabla u\|_{L_T^\infty L_x^2}^2 \|v\|_{L_T^\infty L_x^2}, \end{aligned}$$

for all $0 \leq T \leq T_{max}$.

Hence, by Gronwall inequality, we obtain

$$(3.14) \quad \|\nabla v(t)\|_{L_x^2}^2 \leq 2\Phi^2 e^{2G(T)^2 t}, \quad t \in [0, T_{max}].$$

The estimate (3.14) proves Theorem 1.3.

We also have the following.

REMARK 3.1. *Let $\mu > 0$ and let $u_\mu(t), v_\mu(t)$ be the solutions of the initial value problem (1.1), with $u_\mu \in C([0, T_{max}], L_x^2)$. If the initial data $v_0 \in L^1$, then we get*

$$\limsup_{\mu \rightarrow \infty} \|v_\mu(t)\|_{L^1} \leq \|u_0\|_{L^2}, \quad t \in [0, T_{max}],$$

and

$$\left| \int_{\mathbb{R}^n} v_\mu(t, x) dx - \lambda \|u_0\|_{L^2}^2 \right| \xrightarrow{\mu \rightarrow \infty} 0.$$

PROOF. Observe that

$$(3.15) \quad |v_\mu(t)| \leq e^{-\frac{t}{\mu}} |v_0| + \frac{|\lambda|}{\mu} \int_0^t e^{-\frac{t-s}{\mu}} |u_\mu(s)|^2 ds.$$

Integrating (3.15) and using the conserved quantity (1.2), we have

$$(3.16) \quad \int_{\mathbb{R}^n} |v_\mu(x, t)| dx \leq e^{-\frac{t}{\mu}} (\|v_0\|_1 - |\lambda| \|u_0\|_{L^2}^2) + |\lambda| \|u_0\|_{L^2}^2, \quad \forall t \in [0, T_{\max}].$$

Similarly, we obtain

$$\int_{\mathbb{R}^n} v(x, t) dx = e^{-\frac{t}{\mu}} \left(\int_{\mathbb{R}^n} v_0(x) dx - \lambda \|u_0\|_{L^2}^2 \right) + \lambda \|u_0\|_{L^2}^2, \quad \forall t \in [0, T_{\max}].$$

Note that this equality proves the remark.

LEMMA 3.1. *Let $T > 0$, $1 \leq p < q \leq \infty$ and $A, B \geq 0$. It follows that there exists $\Gamma = \Gamma(B, p, q, T)$ such that if $f \in L^q_{(0, T)}$ satisfies*

$$(3.17) \quad \|f\|_{L^q_{(0, t)}} \leq A + B \|f\|_{L^p_{(0, t)}},$$

for all $0 < t < T$, then

$$\|f\|_{L^q_{(0, t)}} \leq A\Gamma.$$

PROOF. See [6]. □

REMARK 3.2. *In order to estimate $\|u\|_{L^4_x L^4_{(0, t)}}$, we also could have used the Lemma 3.1.*

In fact, let $0 \leq t \leq T \leq T_{max}$ and let $v_0 \in L^2$, $u_0 \in L^2$. Since (4, 4) is an admissible pair in \mathbb{R}^2 , using the integral equation for u and the global well-posedness result in $L^2 \times L^2$, we have

$$\begin{aligned} \|u\|_{L^4_x L^4_{(0, t)}} &\leq C \|u_0\|_{L^2} + \left\| \int_0^t S(t-t') u(t') v(t') dt' \right\|_{L^4_x L^4_{(0, t)}} \\ &\leq C \|u_0\|_{L^2} + C \|u v\|_{L^{4/3}_x L^{4/3}_{(0, t)}} \\ &\leq C \|u_0\|_{L^2} + C \|u\|_{L^{4/3}_{(0, t)} L^4_x} \|v\|_{L^\infty_{(0, T)} L^2_x}. \end{aligned}$$

Now by Lemma 3.1 we conclude that

$$\|u\|_{L^4_x L^4_{(0, t)}} \leq \|u_0\|_{L^2} \Gamma(\|v\|_{L^\infty_{(0, T)} L^2_x}, T).$$

4. Global well-posedness in $H^s(\mathbb{R}^2) \times L^2(\mathbb{R}^2)$, $2/3 < s \leq 1$

In this section we will prove Theorem 1.4.

4.1. A priori Estimates. If one takes $\lambda = -1$ in (1.3), then the energy of the system is decreasing, i.e.,

$$(4.1) \quad E(u, v) = E(u(t), v(t)) \leq E(u_0, v_0), \quad \forall t \geq 0,$$

where

$$(4.2) \quad E(u, v) = \int_{\mathbb{R}^n} (|\nabla u|^2 + 2v|u|^2 + v^2) dx = \int_{\mathbb{R}^n} (|\nabla u|^2 - |u|^4 + \mu^2 |v_t|^2) dx.$$

From (1.2), (4.2) and the immersion

$$(4.3) \quad \|u(t)\|_{L_x^4} \leq c_0 \|u(t)\|_{L_x^2}^{1/2} \|\nabla u(t)\|_{L_x^2}^{1/2},$$

we get

$$\begin{aligned} \int_{\mathbb{R}^2} |\nabla u|^2 dx + \int_{\mathbb{R}^2} |v|^2 dx &\leq E(u_0, v_0) - 2 \int_{\mathbb{R}^2} |u|^2 v dx \\ &\leq E(u_0, v_0) + 2 \|v\|_{L_x^2} \|u\|_{L_x^4}^2 \\ &\leq E(u_0, v_0) + 2c_0^2 \|v\|_{L_x^2} \|u_0\|_{L_x^2} \|\nabla u(t)\|_{L_x^2}. \end{aligned}$$

Thus

$$(4.4) \quad \frac{1}{2} \int_{\mathbb{R}^2} |\nabla u|^2 dx + (1 - 2c_0^4 \|u_0\|_{L_x^2}^2) \int_{\mathbb{R}^2} |v|^2 dx \leq E(u_0, v_0),$$

which gives

$$(4.5) \quad \int_{\mathbb{R}^2} |v|^2 dx \leq 2E(u_0, v_0), \quad \int_{\mathbb{R}^2} |\nabla u|^2 dx \leq 2E(u_0, v_0), \quad \forall t \geq 0,$$

since

$$(4.6) \quad 2c_0^2 \|u_0\|_{L_x^2} \leq 1.$$

In a similar way, by (4.2), the immersion (4.3) and (4.5), we get

$$\begin{aligned} \int_{\mathbb{R}^2} \mu^2 |v_t|^2 dx &\leq E(u_0, v_0) + \int_{\mathbb{R}^2} |u|^4 dx \\ &\leq E(u_0, v_0) + c_0^4 \|u_0\|_{L_x^2}^2 \|\nabla u(t)\|_{L_x^2}^2 \\ (4.7) \quad &\leq \frac{3}{2} E(u_0, v_0), \quad \text{if } 2c_0^2 \|u_0\|_{L_x^2} \leq 1. \end{aligned}$$

Moreover, by (4.3), (4.4) and (4.6), also is not difficult to see that

$$(4.8) \quad \int_{\mathbb{R}^2} |u|^4 dx \leq \frac{1}{2} E(u_0, v_0), \quad \text{if } 2c_0^2 \|u_0\|_{L_x^2} \leq 1.$$

REMARK 4.1. **1)** As a consequence of the immersion (4.3) it follows that if $c_0^4 \|u_0\|_{L_x^2}^2 \leq 1$, then

$$E(u, v) \geq 0.$$

2) If $4c_0^2 \|u_0\|_{L_x^2} \leq 1$, then

$$(4.9) \quad \int_{\mathbb{R}^2} |\nabla u|^2 dx + \int_{\mathbb{R}^2} |v|^2 dx \leq \frac{5}{3} \int_{\mathbb{R}^2} |\nabla u_0|^2 dx + \frac{5}{3} \int_{\mathbb{R}^2} |v_0|^2 dx.$$

In fact, by (4.2) and (4.1), we have

$$\begin{aligned} \int_{\mathbb{R}^2} |\nabla u|^2 dx + \int_{\mathbb{R}^2} |v|^2 dx &\leq \int_{\mathbb{R}^2} |\nabla u_0|^2 dx + \int_{\mathbb{R}^2} |v_0|^2 dx \\ &\quad + 2c_0^2 \|v\|_{L_x^2} \|u_0\|_{L_x^2} \|\nabla u(t)\|_{L_x^2} + 2c_0^2 \|v_0\|_{L_x^2} \|u_0\|_{L_x^2} \|\nabla u_0\|_{L_x^2}, \end{aligned}$$

and using the Young inequality we deduce (4.9).

3) The integral representation (1.6) of v and the Cauchy-Schwartz inequality give

$$(4.10) \quad \|v(t)\|_{L_x^2} \leq e^{-t/\mu} \|v_0\|_{L^2} + \sqrt{\frac{1}{2\mu}} (1 - e^{-2t/\mu})^{1/2} \|u\|_{L_{[0,t]}^4 L_x^4}, \quad t \geq 0,$$

and the estimate (4.8) shows that if $2c_0^2\|u_0\|_{L^2} \leq 1$, then

$$(4.11) \quad \|v(t)\|_{L_x^2} \leq e^{-t/\mu}\|v_0\|_{L^2} + \sqrt{\frac{1}{\mu}}E(u_0, v_0)^{1/2}(1 - e^{-2t/\mu})^{1/2}t^{1/2}, \quad t \geq 0.$$

4.2. Iteration. Now let $v_0 \in L^2$ and $u_0 \in H^s$, $2/3 < s < 1$, be the initial data of the IVP (1.1), with the small condition (4.6), i.e.,

$$(4.12) \quad 2c_0^2\|u_0\|_{L_x^2} < 1.$$

Fix a large time T and let $N = N(T)$ be a cutoff (to be specified). Write

$$(4.13) \quad u_0 := \omega_0 + \eta_0 \quad \text{with} \quad \omega_0 = \int_{|\xi| < N} e^{ix \cdot \xi} \widehat{u_0}(\xi) d\xi, \quad \eta_0 = \int_{|\xi| \geq N} e^{ix \cdot \xi} \widehat{u_0}(\xi) d\xi.$$

We observe that

$$(4.14) \quad \|\omega_0\|_{\dot{H}^\theta} = \left\{ \int_{|\xi| < N} |\xi|^{2\theta} |\widehat{u_0}(\xi)|^2 d\xi \right\}^{1/2} < \|u_0\|_{\dot{H}^s} N^{\theta-s}, \quad \text{for any } \theta \geq s.$$

Similarly, we get

$$(4.15) \quad \|\eta_0\|_{\dot{H}^\sigma} = \left\{ \int_{|\xi| \geq N} |\xi|^{2\sigma} |\widehat{u_0}(\xi)|^2 d\xi \right\}^{1/2} \leq \|u_0\|_{\dot{H}^s} N^{\sigma-s}, \quad \text{if } \sigma \leq s,$$

and

$$(4.16) \quad \|\omega_0\|_{L_x^2} \leq \|u_0\|_{L_x^2},$$

In particular, these estimates proves that $w_0 \in H^\infty$ and $\eta_0 \in H^s$. Now we consider the IVP

$$(4.17) \quad \begin{cases} i\dot{u}_t + \frac{1}{2}\Delta \dot{u} = \dot{u}\dot{v}, \\ \mu\dot{v}_t + \dot{v} = -|\dot{u}|^2, \\ \dot{u}(0) = \omega_0 \quad \text{and} \quad \dot{v}(0) = v_0, \end{cases}$$

where w_0 is defined in (4.13) and verifies (4.14) and (4.16). We know by the Theorem 1.2 that if the initial data $(u_0, v_0) \in H^1(\mathbb{R}^2) \times L^2(\mathbb{R}^2)$, then there exists a unique solution $(u, v) \in C([0, T]; H^1(\mathbb{R}^2) \times L^2(\mathbb{R}^2))$ of IVP (1.1). Shortly problem (4.17) is globally well-posedness in $H^1(\mathbb{R}^2) \times L^2(\mathbb{R}^2)$.

We write the solution (u, v) of the system (1.1), as

$$(4.18) \quad u = \dot{u} + \dot{\eta} \quad \text{and} \quad v = \dot{v} + \dot{z},$$

where $(\dot{\eta}, \dot{z})$ satisfies the IVP

$$(4.19) \quad \begin{cases} i\dot{\eta}_t + \frac{1}{2}\Delta \dot{\eta} = (\dot{u} + \dot{\eta})(\dot{v} + \dot{z}) - \dot{u}\dot{v}, \\ \mu\dot{z}_t + \dot{z} = -|\dot{u} + \dot{\eta}|^2 + |\dot{u}|^2, \\ \dot{\eta}(0) = \eta_0, \quad \dot{z}(0) = z_0 = 0, \end{cases}$$

where η_0 is defined in (4.13) and verifies (4.15).

Consider the IVP (4.19) in the following integral form

$$(4.20) \quad \overset{0}{\eta}(t) = e^{it\Delta/2}\eta_0 + \overset{0}{\psi}(t),$$

where

$$(4.21) \quad \overset{0}{\psi}(t) = i \int_0^t e^{i\Delta(t-\tau)/2} (\overset{0}{u} \overset{0}{z} + \overset{0}{\eta} \overset{0}{v} + \overset{0}{\eta} \overset{0}{z}) d\tau,$$

and

$$(4.22) \quad \overset{0}{z} = \frac{1}{\mu} \int_0^t e^{-(t-\tau)/\mu} \left(-|\overset{0}{\eta}|^2 - 2\text{Re}(\overset{0}{u} \overline{\overset{0}{\eta}}) \right) d\tau.$$

We have the following LWP result

THEOREM 4.1. *Let $(\tilde{\omega}_0, \tilde{v}_0) \in H^{s_1}(\mathbb{R}^2) \times H^{s_2}(\mathbb{R}^2)$, where (s_1, s_2) satisfy the condition (1.7). Then, for all T such that*

$$(4.23) \quad 0 < T \leq \frac{c}{\|\tilde{\omega}_0\|_{H^{s_1}}^2 + \|\tilde{v}_0\|_{H^{s_2}}^2},$$

there exists a unique solution

$$(\overset{0}{u}, \overset{0}{v}) \in C([0, T]; H^{s_1}(\mathbb{R}^2) \times H^{s_2}(\mathbb{R}^2))$$

to the initial value problem (4.17) with initial data $(\tilde{\omega}_0, \tilde{v}_0)$ such that:

$$(4.24) \quad \|\vartheta_T(t)\overset{0}{u}(\cdot, t)\|_{X_{s_1, b_1}} + \|\vartheta_T(t)\overset{0}{v}(\cdot, t)\|_{H_{s_2, b_2}} \leq c \|\tilde{\omega}_0\|_{H^{s_1}} + c \|\tilde{v}_0\|_{H^{s_2}},$$

where ϑ_T is defined in (2.1) and for some suitable $b_1 > 1/2$ and $b_2 > 1/2$.

PROOF. It follows immediately from the proof of Theorem 1.1 in Corcho, Oliveira and Silva [9]. □

We also have the following.

THEOREM 4.2. *Let u_0 be in $H^s(\mathbb{R}^2)$ and let η_0 be as defined in (4.13). Then, there exists $t_1 > 0$ such that*

$$(4.25) \quad t_1 = N^{-2(1-s)^-},$$

and there exists a unique solution

$$(\overset{0}{\eta}, \overset{0}{z}) \in C([0, t_1]; H^s(\mathbb{R}^2) \times L^2(\mathbb{R}^2))$$

to the initial value problem (4.19) with initial data η_0 and $z_0 = 0$, such that

$$(4.26) \quad \|\vartheta_{t_1}(t)\overset{0}{\eta}(\cdot, t)\|_{X_{s, b}} \leq c \|\eta_0\|_{H^s},$$

for some suitable $b > 1/2$.

PROOF. The proof is very similar to the proof of Theorem 1.1 and to the proof of Theorem 4.1 in Corcho, Oliveira and Silva [9]. □

REMARK 4.2. *If in Theorem 4.1 we take $s_1 = 1, s_2 = 0$ and if also we consider $\tilde{\omega}_0 = \omega_0$ as defined in (4.13) and $\tilde{v}_0 = v_0$, then we can take $T = t_1 = N^{-2(1-s)^-}$ and thus obtain the same existence interval $[0, t_1]$, for the systems (4.17) and (4.19).*

Note that by (4.14), we have $\|\nabla\omega_0\|_{L_x^2(\mathbb{R}^2)} \lesssim \|\omega_0\|_{\dot{H}^1(\mathbb{R}^2)} \leq \|u_0\|_{\dot{H}^s} N^{1-s}$. From (4.16) it follows that

$$\begin{aligned} E(\omega_0, v_0) &= \int_{\mathbb{R}^2} (|\nabla\omega_0|^2 + 2|\omega_0|^2 v_0 + |v_0|^2) dx \leq \|\nabla\omega_0\|_{L_x^2(\mathbb{R}^2)}^2 + \|v_0\|_{L_x^2(\mathbb{R}^2)}^2 \\ &\quad + c_0 \|\omega_0\|_{L_x^2(\mathbb{R}^2)} \|\nabla\omega_0\|_{L_x^2(\mathbb{R}^2)} \|v_0\|_{L_x^2(\mathbb{R}^2)} \\ (4.27) \quad &\lesssim_s N^{2(1-s)}. \end{aligned}$$

Integrating the inequality (4.8), considering the time t_1 as in (4.25), we have

$$(4.28) \quad \int_0^{t_1} \int_{\mathbb{R}^2} |u|^4 dx dt \leq \frac{1}{2} E(u_0, v_0) t_1 \leq 1.$$

By Remark (4.2) in the time interval $[0, t_1]$ we have local existence for both systems (4.17) and (4.19).

Now from (4.16), $2c_0^2 \|\omega_0\|_{L_x^2} \leq 2c_0^2 \|u_0\|_{L_x^2} < 1$, thus by (4.5) we have

$$(4.29) \quad \|\nabla u^0(t_1)\|_{L_x^2(\mathbb{R}^2)} \lesssim_s N^{(1-s)},$$

and

$$(4.30) \quad \|\dot{v}^0(t_1)\|_{L_x^2(\mathbb{R}^2)} \lesssim_s N^{(1-s)}.$$

Furthermore, the immersion (4.3), conservation law (1.2), inequalities (4.29) above and (4.36) below imply that

$$(4.31) \quad \|\psi^0(t_1)\|_{L_x^4} \lesssim N^{(-3s/4)+} \quad \text{and} \quad \|u^0(t_1)\|_{L_x^4} \lesssim_s N^{(1-s)/2}.$$

For $t \in [0, t_1]$, by (4.18) and (4.20), we have

$$(4.32) \quad u(t) = \dot{u}^0(t) + \psi^0(t) + e^{it\Delta/2} \eta_0 \quad \text{and} \quad v(t) = \dot{v}^0(t) + \dot{z}^0(t),$$

where (\dot{u}^0, \dot{v}^0) is the solution of (4.17) and $(\dot{\eta}^0, \dot{z}^0) = (\psi + e^{it\Delta/2} \eta_0, \dot{z})$ is the solution of (4.19). Now we define the new initial data for the second iteration

$$(4.33) \quad \begin{aligned} \omega_1 &= \dot{u}^0(t_1) + \psi^0(t_1) \quad \text{and} \quad v_1 = \dot{v}^0(t_1) + \dot{z}^0(t_1), \\ \eta_1 &= e^{it_1\Delta/2} \eta_0 \quad \text{and} \quad z_1 = 0. \end{aligned}$$

In each iteration we consider the decomposition of the initial data as in (4.33). Therefore $\eta_1, \dots, \eta_k = e^{ik t\Delta/2} \eta_0$ have the same properties of η_0 with $\|\eta_k\|_{H^s} = \|\eta_0\|_{H^s}$ and $z_1 = \dots = z_k = 0$. We hope that $\omega_1, \dots, \omega_k$ and v_1, \dots, v_k also have the same properties of ω_0 and v_0 respectively in order to ensure the same existence interval $[0, t_1]$ in each iteration and attach the existence interval $[0, T]$, extending the solution of the systems (4.17) and (4.19). This fact is proved by induction. Here we will prove only the case $k = 1$ and note that a similar argument works in the general case.

From (4.1), we have

$$E(\dot{u}^0(t_1), \dot{v}^0(t_1)) \leq E(\dot{u}^0(0), \dot{v}^0(0)) = E(\omega_0, v_0).$$

Thus we get

$$(4.34) \quad E(\omega_1, v_1) \leq E(\omega_0, v_0) + \left[E(\omega_1, v_1) - E(\overset{0}{u}(t_1), \overset{0}{v}(t_1)) \right].$$

Using the immersion (4.3) and (4.33), we obtain

$$(4.35) \quad \begin{aligned} |E(\omega_1, v_1) - E(\overset{0}{u}(t_1), \overset{0}{v}(t_1))| &= |E(\overset{0}{u}(t_1) + \overset{0}{\psi}(t_1), \overset{0}{v}(t_1) + \overset{0}{z}(t_1)) - E(\overset{0}{u}(t_1), \overset{0}{v}(t_1))| \\ &\leq \|\nabla \overset{0}{\psi}(t_1)\|_{L_x^2}^2 + 2\|\nabla \overset{0}{\psi}(t_1)\|_{L_x^2} \|\nabla \overset{0}{u}(t_1)\|_{L_x^2} + 2\|\overset{0}{v}(t_1)\|_{L_x^2(\mathbb{R}^2)} \|\overset{0}{\psi}(t_1)\|_{L_x^4}^2 \\ &\quad + 4\|\overset{0}{u}(t_1)\|_{L_x^4} \|\overset{0}{\psi}(t_1)\|_{L_x^4} \left(\|\overset{0}{v}(t_1)\|_{L_x^2} + \|\overset{0}{z}(t_1)\|_{L_x^2} \right) + 2\|\overset{0}{z}(t_1)\|_{L_x^2} \|\overset{0}{\psi}(t_1)\|_{L_x^4}^2 \\ &\quad + \|\overset{0}{z}(t_1)\|_{L_x^2}^2 + 2\|\overset{0}{v}(t_1)\|_{L_x^2} \|\overset{0}{z}(t_1)\|_{L_x^2} + 2\|\overset{0}{z}(t_1)\|_{L_x^2} \|\overset{0}{u}(t_1)\|_{L_x^4}. \end{aligned}$$

In order to estimate (4.35), initially we will assume the following result, which be will proved later.

LEMMA 4.1. *Let $\overset{0}{\eta}(t)$ be a solution of the IVP (4.19), and let $\overset{0}{\psi}(t)$ be the forcing term as defined in (4.20) and (4.21), then we have the following estimates*

$$(4.36) \quad \|\overset{0}{\psi}(t)\|_{L_x^2(\mathbb{R}^2)} \leq cN^{-s} \quad \text{and} \quad \|\nabla \overset{0}{\psi}(t)\|_{L_x^2(\mathbb{R}^2)} \leq cN^{(-s/2)^+},$$

and also that

$$(4.37) \quad \|\overset{0}{\eta}\|_{L_t^4 L_x^4([0, t_1] \times \mathbb{R}^2)} \leq cN^{-s}.$$

Using (4.22), the Minkowsky and the Cauchy-Schwartz inequalities, together with (4.28) and (4.37), for any $t \in [0, t_1]$ we have

$$(4.38) \quad \begin{aligned} \|\overset{0}{z}(t)\|_{L_x^2(\mathbb{R}^2)} &\leq \frac{1}{\mu} \int_0^t e^{-(t-\tau)/\mu} \left(\|\overset{0}{\eta}(\tau)\|_{L_x^4(\mathbb{R}^2)}^2 + 2\|\overset{0}{\eta}(\tau)\|_{L_x^4(\mathbb{R}^2)} \|\overset{0}{u}(\tau)\|_{L_x^4(\mathbb{R}^2)} \right) \\ &\lesssim_{\mu} \|\overset{0}{\eta}\|_{L_{[0, t_1]}^4 L_x^4}^2 + \|\overset{0}{\eta}\|_{L_{[0, t_1]}^4 L_x^4} \|\overset{0}{u}\|_{L_{[0, t_1]}^4 L_x^4} \\ &\lesssim_{\mu} \|\overset{0}{\eta}\|_{L_{[0, t_1]}^4 L_x^4}^2 + \|\overset{0}{\eta}\|_{L_{[0, t_1]}^4 L_x^4} \end{aligned}$$

$$(4.39) \quad \lesssim_{\mu} N^{-s}.$$

From (4.29), (4.30), (4.31), (4.35), (4.36) and (4.39) we obtain

$$(4.40) \quad \begin{aligned} |E(\omega_1, v_1) - E(\overset{0}{u}(t_1), \overset{0}{v}(t_1))| &\lesssim N^{(-s)^+} + N^{1-s} N^{(-s/2)^+} \\ &\quad + N^{(1-s)} N^{(-3s/2)^+} + N^{(1-s)/2} N^{(-3s/4)^+} \left(N^{(1-s)} + N^{-s} \right) \\ &\quad + N^{-s} N^{(-3s/2)^+} + N^{-2s} + N^{1-s} N^{-s} + N^{1-s} N^{-s} \lesssim N^{((2-3s)/2)^+}. \end{aligned}$$

Combining (4.27), (4.34) and (4.40), we get that

$$(4.41) \quad E(\omega_1, v_1) \leq E(\omega_0, v_0) + cN^{((2-3s)/2)^+}.$$

Also, observe that by conservation quantity (1.2) and Lemma 4.1, we have

$$(4.42) \quad \|\omega_1\|_{L^2} \leq \|\overset{0}{u}(t_1)\|_{L^2} + \|\overset{0}{\psi}(t_1)\|_{L^2} \leq \|\omega_0\|_{L^2} + cN^{-s} \leq \|u_0\|_{L^2} + cN^{-s}.$$

Thus, the small condition (4.12) remains valid in the second iteration if

$$(4.43) \quad 2c_0^2 \|\omega_1\|_{L^2} \leq 2c_0^2 (\|u_0\|_{L^2} + cN^{-s}) < 1,$$

i.e., if $2c_0^2 cN^{-s} < 1 - 2c_0^2 \|u_0\|_{L^2}$, which happens indeed if N is very large. Also from (4.5), it follows that

$$(4.44) \quad \|v_1\|_{L^2} \leq \|\overset{0}{v}(t_1)\|_{L^2} + \|\overset{0}{z}(t_1)\|_{L^2} \leq \sqrt{2E(\omega_0, v_0)} + cN^{-s} \leq cN^{1-s}.$$

The number of steps in the iteration is

$$\frac{T}{t_1} \sim TN^{2(1-s)^+}.$$

Thus, by (4.27), we need that

$$TN^{2(1-s)^+} N^{((2-3s)/2)^+} < E(\omega_0, v_0) \sim N^{2(1-s)},$$

which is possible if $s > 2/3$ and

$$N = N(T) = T^{2^+/(3s-2)}, \quad \text{or equivalently} \quad T = N^{(3s-2)/(2^+)}.$$

Observe also that the small condition remains valid in each iteration since, in similar way as in (4.43), we have

$$(4.45) \quad TN^{2(1-s)^+} 2c_0^2 cN^{-s} = N^{(3s-2)/(2^+)} N^{2(1-s)^+} 2c_0^2 cN^{-s} = 2c_0^2 cN^{(2-3s)/2} < 1 - 2c_0^2 \|u_0\|_{L^2},$$

and similarly as in (4.44)

$$(4.46) \quad \sqrt{2E(\omega_0, v_0)} + TN^{2(1-s)^+} cN^{-s} \leq \sqrt{2E(\omega_0, v_0)} + N^{(3s-2)/2^+} N^{2(1-s)^+} cN^{-s} \leq cN^{1-s},$$

and the inequalities (4.45), (4.46) are true if N is very large and $s > 2/3$.

4.3. Proof of Lemma 4.1. First we will prove the inequality (4.37). Since (4, 4) is an admissible pair of the group $\{e^{it\Delta/2}\}$, by (4.20) and Proposition 2.4, it follows that

$$(4.47) \quad \|\overset{0}{\eta}\|_{L^4_{[0,t_1]} L^4_x} \leq c \|\eta_0\|_{L^2_x} + \|\overset{0}{\psi}\|_{L^4_t L^4_x([0,t_1] \times \mathbb{R}^2)}.$$

Moreover, Proposition 2.4, the equality (4.21) and the Hölder inequality show that

$$(4.48) \quad \begin{aligned} \|\overset{0}{\psi}\|_{L^4_t L^4_x([0,t_1] \times \mathbb{R}^2)} &\lesssim \|\overset{0}{u} \overset{0}{z} + \overset{0}{\eta} \overset{0}{v} + \overset{0}{\eta} \overset{0}{z}\|_{L^{4/3}_{[0,t_1]} L^{4/3}_x} \\ &\lesssim \|\overset{0}{u}\|_{L^4_{[0,t_1]} L^4_x} \|\overset{0}{z}\|_{L^2_{[0,t_1]} L^2_x} + \|\overset{0}{\eta}\|_{L^4_{[0,t_1]} L^4_x} \|\overset{0}{v}\|_{L^2_{[0,t_1]} L^2_x} \\ &\quad + \|\overset{0}{\eta}\|_{L^4_{[0,t_1]} L^4_x} \|\overset{0}{z}\|_{L^2_{[0,t_1]} L^2_x}. \end{aligned}$$

By estimates (4.5) and (4.27), we get

$$(4.49) \quad \|\overset{0}{v}(t)\|_{L^2_x} \leq \sqrt{2E(\omega_0, v_0)} \lesssim N^{1-s}, \quad t \geq 0.$$

Therefore, combining (4.28), (4.38), (4.48) and (4.49), we obtain

$$\begin{aligned}
& \|\overset{0}{\psi}\|_{L^4_{[0,t_1]}L^4_x} \lesssim t_1^{1/2} \|\overset{0}{z}\|_{L^\infty_{[0,t_1]}L^2_x} + t_1^{1/2} \|\overset{0}{\eta}\|_{L^4_{[0,t_1]}L^4_x} \|\overset{0}{v}\|_{L^\infty_{[0,t_1]}L^2_x} \\
& \quad + t_1^{1/2} \|\overset{0}{\eta}\|_{L^4_{[0,t_1]}L^4_x} \|\overset{0}{z}\|_{L^\infty_{[0,t_1]}L^2_x} \\
(4.50) \quad & \lesssim t_1^{1/2} \|\overset{0}{\eta}\|_{L^4_{[0,t_1]}L^4_x} \left(N^{1-s} + \|\overset{0}{\eta}\|_{L^4_{[0,t_1]}L^4_x} + \|\overset{0}{\eta}\|_{L^4_{[0,t_1]}L^4_x}^2 \right).
\end{aligned}$$

Note that $t_1^{1/2} = N^{-(1-s)^-}$, $1 \ll N$. Thus, it follows from (4.47) and (4.50) that

$$\|\overset{0}{\eta}\|_{L^4_{[0,t_1]}L^4_x} \leq c \|\eta_0\|_{L^2_x} + c t_1^{1/2} \|\overset{0}{\eta}\|_{L^4_{[0,t_1]}L^4_x}^2 \left(1 + \|\overset{0}{\eta}\|_{L^4_{[0,t_1]}L^4_x} \right),$$

and from a standard continuity argument it follows that

$$(4.51) \quad \|\overset{0}{\eta}\|_{L^4_t L^4_x([0,t_1] \times \mathbb{R}^2)} \leq 2c \|\eta_0\|_{L^2_x} \lesssim N^{-s}.$$

Now we will prove the first inequality in (4.36). Since $(\infty, 2)$ and $(4, 4)$ are admissible pairs of the group $\{e^{it\Delta/2}\}$, using (4.21), (4.48)-(4.51) it follows that

$$\begin{aligned}
\|\overset{0}{\psi}\|_{L^2_x} & \lesssim \|\overset{0}{u} \overset{0}{z} + \overset{0}{\eta} \overset{0}{v} + \overset{0}{\eta} \overset{0}{z}\|_{L^{4/3}_{[0,t_1]}L^{4/3}_x} \\
& \lesssim t_1^{1/2} \|\overset{0}{\eta}\|_{L^4_{[0,t_1]}L^4_x} \left(N^{1-s} + \|\overset{0}{\eta}\|_{L^4_{[0,t_1]}L^4_x} + \|\overset{0}{\eta}\|_{L^4_{[0,t_1]}L^4_x}^2 \right) \\
(4.52) \quad & \lesssim N^{-s}.
\end{aligned}$$

Finally we will prove the second inequality in (4.36). By Theorem 2.1, it follows that

$$\begin{aligned}
\|\nabla \overset{0}{\psi}(t)\|_{L^2_x(\mathbb{R}^2)} & \leq c \|\overset{0}{\psi}\|_{X_{1,b}} \\
& \leq c \|\overset{0}{u} \overset{0}{z} + \overset{0}{\eta} \overset{0}{v} + \overset{0}{\eta} \overset{0}{z}\|_{X_{1,b-1}} \\
& = c \sup_{\|W\|_{X_{\{-1,1-b\}} \leq 1}} \left| \int_{\mathbb{R}_t \times \mathbb{R}_x^2} (\overset{0}{u} \overset{0}{z} + \overset{0}{\eta} \overset{0}{v} + \overset{0}{\eta} \overset{0}{z}) \overline{W} \, dx \, dt \right| \\
(4.53) \quad & = c \sup_{\|W\|_{X_{\{0,1-b\}} \leq 1}} \left| \int_{\mathbb{R}_t \times \mathbb{R}_x^2} \mathcal{D}_x^1 (\overset{0}{u} \overset{0}{z} + \overset{0}{\eta} \overset{0}{v} + \overset{0}{\eta} \overset{0}{z}) \overline{W} \, dx \, dt \right|,
\end{aligned}$$

where $(\widehat{\mathcal{D}_x^s f})(\xi) = \langle \xi \rangle^s \widehat{f}(\xi)$, $\mathcal{D}_x^1 := \mathcal{D}_x$. Without loss of generality we only consider the term with $\overset{0}{u} \overset{0}{z}$, because the estimates on the other terms in (4.53) are similar or better. Using the Plancherel equality and the Hölder inequality, we have

$$\begin{aligned}
& \left| \int_{\mathbb{R}_t \times \mathbb{R}_x^2} \mathcal{D}_x (\overset{0}{u} \overset{0}{z}) \overline{W} \, dx \, dt \right| \leq \left| \int_{\mathbb{R}_t \times \mathbb{R}_x^2} \mathcal{D}_x (\overset{0}{u}) \overset{0}{z} \overline{W} \, dx \, dt \right| + \left| \int_{\mathbb{R}_t \times \mathbb{R}_x^2} \overset{0}{u} \mathcal{D}_x (\overset{0}{z}) \overline{W} \, dx \, dt \right| \\
& \leq \|\overset{0}{z}\|_{L^{2^+}_{t_1} L^{2^+}_x} \|\mathcal{D}_x (\overset{0}{u})\|_{L^4_{t_1} L^4_x} \|W\|_{L^{4^-}_{t_1} L^{4^-}_x} + \left| \int_{\mathbb{R}_t \times \mathbb{R}_x^2} \mathcal{D}_x^{1/2^+} (\overset{0}{z}) \mathcal{D}_x^{1/2^-} (\overline{W} \overset{0}{u}) \, dx \, dt \right| \\
& \leq \|\overset{0}{z}\|_{L^{2^+}_{t_1} L^{2^+}_x} \|\mathcal{D}_x (\overset{0}{u})\|_{L^4_{t_1} L^4_x} \|W\|_{L^{4^-}_{t_1} L^{4^-}_x} + \|\mathcal{D}_x^{1/2^+} (\overset{0}{z})\|_{L^2_{t_1} L^2_x} \|\mathcal{D}_x^{1/2^-} (\overline{W} \overset{0}{u})\|_{L^2_{t_1} L^2_x} \\
(4.54) \quad & := I_1 I_2 I_3 + I_4 I_5.
\end{aligned}$$

Now we will estimate all the terms in (4.54):

(1) Estimate of I_5 :

Observe that the local well-posedness theory in Theorem 4.1, the conservation quantity (1.2) and a priori estimates (4.5) imply that

$$(4.55) \quad \|\vartheta_{t_1}^0 u(\cdot, t)\|_{X_{1,b}} \leq c \|\omega_0\|_{H^1} + c \|v_0\|_{L^2} \leq c N^{1-s},$$

and

$$(4.56) \quad \|\vartheta_{t_1}^0 u(\cdot, t)\|_{X_{0,b}} \leq c \|\omega_0\|_{L^2} + c \|v_0\|_{L^2} \leq c.$$

Using the Proposition 2.1 item (4), interpolation, (4.55) and (4.56), we have

$$\begin{aligned} I_5 &= \|\mathcal{D}_x^{1/2^-} (\overline{W}^0 u)\|_{L_{t_1}^2 L_x^2} \leq c \|u\|_{X_{1/2,1/2^+}}^0 \|\overline{W}\|_{X_{0,1/2^-}} \\ &\leq c \|u\|_{X_{1,1/2^+}}^{1/2} \|u\|_{X_{0,1/2^+}}^{1/2} \\ &\leq c N^{(1-s)/2}. \end{aligned}$$

(2) Estimate of I_3 :

Using the Proposition 2.2 item ii) (because (4, 4) is an admissible pair) we have

$$\|W\|_{L_{x,t}^4} \leq c \|W(\cdot, t)\|_{X_{0,1/2^+}}.$$

Interpolating this inequality with

$$\|W\|_{L_{x,t}^2} \leq c \|W(\cdot, t)\|_{X_{0,0}},$$

we obtain

$$(4.57) \quad I_3 = \|W\|_{L_{t_1}^{4^-} L_x^{4^-}} \leq c \|W(\cdot, t)\|_{X_{0,1/2^-}} \leq c.$$

(3) Estimate of I_2 :

To estimate I_2 we also will use the Proposition 2.2 item ii). Thus

$$I_2 = \|\mathcal{D}_x^0 u\|_{L_{t_1}^4 L_x^4} \leq c \|\vartheta_{t_1}^0 u\|_{X_{1,1/2^+}} \leq c N^{1-s}.$$

(4) Estimate of I_1 and I_4 :

By Theorem 4.2 and the bilinear estimates in Proposition 2.4 of Corcho, Oliveira and Silva [9] for $s > 1/2^+$ we have that

$$(4.58) \quad \|\vartheta_{t_1}^0 z\|_{H_{s,1/2^+}} \leq c \|\eta_0\|_{H^s} \leq c \|u_0\|_{H^s}.$$

From Gagliardo-Nirenberg inequality and (4.39), we get

$$I_1 = \|z\|_{L_{t_1}^{2^+} L_x^{2^+}}^0 = t_1^{1/2^+} \|z\|_{L_{t_1}^\infty L_x^{2^+}}^0 \leq t_1^{1/2} \|z\|_{L_{t_1}^\infty L_x^2}^{1-\theta} \|\mathcal{D}_x^s z\|_{L_{t_1}^\infty L_x^2}^\theta \lesssim_\mu N^{-1^+},$$

where $2\theta = 0^+/(2^+)$. Finally, using interpolating and (4.58), we obtain

$$\begin{aligned} I_4 &= \|\mathcal{D}_x^{1/2^+} (z)\|_{L_{t_1}^2 L_x^2} \leq t_1^{1/2} \|\mathcal{D}_x^s z\|_{L_{t_1}^\infty L_x^2}^{\theta_1} \|z\|_{L_{t_1}^\infty L_x^2}^{1-\theta_1} \leq c t_1^{1/2} \|z\|_{L_{t_1}^\infty L_x^2}^{1-\theta_1} \\ &\lesssim_\mu N^{(-1/2)^+}, \end{aligned}$$

where $s\theta_1 = (1/2)^+$.

Combining (4.53), (4.54) and the estimates on I_1, \dots, I_5 , we have

$$\|\nabla \psi(t)\|_{L_x^2(\mathbb{R}^2)} \leq c N^{(-s/2)^+}.$$

REMARK 4.3. 1) By (4.4), the condition (4.12) can be replaced by the weaker condition:

$$\sqrt{2} c_0^2 \|u_0\|_{L_x^2} < 1.$$

2) The inequality (4.35) shows that a better estimate for $\|\nabla^0 \psi(\cdot)\|_{L_x^2}$ implies a better GWP result.

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