

A SCALING LIMIT FROM THE WAVE MAP TO THE HEAT FLOW INTO \mathbb{S}^{2*}

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Abstract. In this paper we study a limit connecting a scaled wave map with the heat flow into the unit sphere \mathbb{S}^2 . We show quantitatively how the two equations are connected by means of an initial layer correction. This limit is motivated as a first step into understanding the limit of zero inertia for the hyperbolic-parabolic Ericksen-Leslie’s liquid crystal model.

Keywords. Ericksen-Leslie; Wave map; Heat flow; Initial layer.

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1. Introduction

1.1. Wave map and heat flow. We consider a hyperbolic system for functions $d: \mathbb{R}^+ \times \mathbb{R}^3 \rightarrow \mathbb{S}^2$:

$$\partial_t d = -\square d + (|\nabla d|^2 - |\partial_t d|^2)d, \tag{1.1}$$

subject to initial data: for any $x \in \mathbb{R}^3$,

$$d|_{t=0} = d^0(x) \in \mathbb{S}^2, \quad \partial_t d|_{t=0} = \tilde{d}^0(x) \in \mathbb{R}^3, \quad d^0(x) \cdot \tilde{d}^0(x) = 0, \tag{1.2}$$

where $\square = \partial_{tt} - \Delta$ is the standard wave operator, and the compatibility condition $d^0 \cdot \tilde{d}^0 = 0$ on the initial data is due to the fact that $|d^0| \equiv 1$.

The system (1.1) is a wave map from \mathbb{R}^3 to the unit sphere \mathbb{S}^2 , with a damping term $\partial_t d$. One way of interpreting this system is as follows: setting the right-hand side of (1.1) equal to 0, we obtain $\square d = (|\nabla d|^2 - |\partial_t d|^2)d$. This is the well-known *wave map*, which can be characterized variationally as a critical point of the functional

$$\mathcal{A}(d) = \frac{1}{2} \iint (|\nabla d|^2 - |\partial_t d|^2) dx dt, \tag{1.3}$$

among maps d satisfying the target constraint, $d: \mathbb{R}^+ \times \mathbb{R}^3 \rightarrow \mathbb{S}^2$. Thus the full system (1.1) can be viewed as a “gradient flow” of the functional (1.3).

Another gradient flow can be obtained by formally dropping some terms out of the previous system, and obtaining the *heat flow*

$$\partial_t d = \Delta d + |\nabla d|^2 d. \tag{1.4}$$

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Similarly as before setting the right-hand side equal to zero we obtain the equations for the harmonic map from \mathbb{R}^3 to the unit sphere \mathbb{S}^2 namely

$$\Delta d + |\nabla d|^2 d = 0, \tag{1.5}$$

which is a critical point of the energy functional

$$E(d) = \frac{1}{2} \int |\nabla d|^2 dx. \tag{1.6}$$

There exist deep relations between the two systems, (1.1) and (1.4) and one way to see this is by considering the following parabolic scaling:

$$d^\varepsilon(t, x) := d\left(\frac{t}{\varepsilon}, \frac{x}{\sqrt{\varepsilon}}\right), \tag{1.7}$$

Then, d^ε satisfies the following scaled wave map:

$$\partial_t d^\varepsilon = -(\varepsilon \partial_{tt} - \Delta) d^\varepsilon + (|\nabla d^\varepsilon|^2 - \varepsilon |\partial_t d^\varepsilon|^2) d^\varepsilon \tag{1.8}$$

on $\mathbb{R}^+ \times \mathbb{R}^3$. For this scaled system we take the initial values independent of ε , namely:

$$d^\varepsilon|_{t=0} = d^{in}(x) \in \mathbb{S}^2, \quad \partial_t d^\varepsilon|_{t=0} = \tilde{d}^{in}(x) \in \mathbb{R}^3. \tag{1.9}$$

The finite-time behaviour of the limit $\varepsilon \rightarrow 0$ for the system (1.8) with initial data (1.9) is the focus of this paper. It is easy to see that letting $\varepsilon = 0$ in (1.8) will formally give the heat flow (1.4). However, a refined analysis and the introduction of an *initial layer* is needed in order to overcome the singular character of this limit and understand the relationship between the system (1.8) and its formal limit, as it will be seen in the Theorem 1.1 below.

Notations and conventions: Throughout this paper, we use the following standard notations:

$$|d|_{L^p}^p = \int_{\mathbb{R}^3} |d|^p dx, \quad |d|_{H^k} = \sum_{\gamma \leq k} |\nabla^\gamma d|_{L^2}, \quad |d|_{\dot{H}^k} = \sum_{1 \leq \gamma \leq k} |\nabla^\gamma d|_{L^2}.$$

Additionally, for the Hilbert space $L^2 \equiv L^2(dx, \mathbb{R}^3)$, we use the following notation to denote the standard inner product:

$$\langle f, g \rangle = \int_{\mathbb{R}^3} f g dx.$$

Furthermore, if there is a generic constant $C > 0$ such that the inequality $f(t) \leq Cg(t)$ holds for all $t \geq 0$, we denote this inequality by

$$f \lesssim g.$$

1.2. Ericksen-Leslie’s hyperbolic liquid crystal model. Our motivation for considering the previously mentioned limit comes from the hydrodynamic theory of nematic liquid crystals.

The most widely accepted equations of nematics were proposed by Ericksen [5–7] and Leslie [12, 13] in the 1960’s (see for more details Section 5.1 of [14]). The general hyperbolic-parabolic Ericksen-Leslie system consists of an equation for velocity u of the

centers of mass of the rod-like molecules, coupled with an equation for the direction d of these molecules. More specifically we have the following equations (in non-dimensional form):

$$\begin{cases} \partial_t u + u \cdot \nabla u - \frac{1}{2} \mu_4 \Delta u + \nabla p = -\operatorname{div}(\nabla d \odot \nabla d) + \operatorname{div} \tilde{\sigma}, \\ \operatorname{div} u = 0, \\ \varepsilon D_u^2 d = \Delta d + \gamma d + \lambda_1 (D_u d - B d) + \lambda_2 A d \end{cases} \tag{1.10}$$

on $\mathbb{R}^+ \times \mathbb{R}^3$ with constraint $|d|=1$, where $A = \frac{1}{2}(\nabla u + \nabla u^\top)$ and $B = \frac{1}{2}(\nabla u - \nabla u^\top)$, $D_u f = \partial_t f + u \cdot \nabla f$ is the material derivative of f with respect to the vector u , $D_u^2 d = \partial_t D_u d + u \cdot \nabla D_u d$. The Lagrangian γ that enforces the unit-length constraint $|d|=1$ is given by

$$\gamma \equiv \gamma(u, d, D_u d) = -\varepsilon |D_u d|^2 + |\nabla d|^2 - \lambda_2 d^\top A d.$$

The stress tensor appearing in the equation for u is given by:

$$\begin{aligned} \tilde{\sigma}_{ij} \equiv (\tilde{\sigma}(u, d, D_u d))_{ij} &= \mu_1 d_k d_p A_{kp} d_i d_j + \mu_2 d_j ((D_u d)_i + B_{ki} d_k) \\ &\quad + \mu_3 d_i ((D_u d)_j + B_{kj} d_k) + \mu_5 d_j d_k A_{ki} + \mu_6 d_i d_k A_{kj}. \end{aligned}$$

The constant $\varepsilon > 0$ measures the inertial effects. The constants μ_i ($1 \leq i \leq 6$) are known as Leslie coefficients and one has $\mu_4 > 0$. Furthermore, we have:

$$\lambda_1 = \mu_2 - \mu_3, \lambda_2 = \mu_5 - \mu_6, \mu_2 + \mu_3 = \mu_6 - \mu_5,$$

where the last relation is called *Parodi* relation. For more background and derivation of (1.10), see [12] and [9].

For any fixed $\varepsilon > 0$, in [9] the first two authors of the current paper proved the local well-posedness of the system (1.10) under assumptions on the Leslie coefficients which ensure the dissipativity of the basic energy law, and global well-posedness with small initial data under further damping effect, i.e. $\lambda_1 < 0$. In [10], the first three authors also proved similar results for the corresponding compressible case. We also mention that recently in [3, 8], the fourth author of the current paper with collaborators considered the Qian-Sheng model. These works are among the first rigorous analytic treatment of the liquid crystals system involving second order material derivatives.

As noted in the ‘‘Conclusion’’ section of [9], the inertial constant $\varepsilon > 0$ is, physically, in most common non-dimensionalisations and materials, very small. Formally, letting $\varepsilon = 0$ will give the parabolic Ericksen-Leslie system which is basically a coupling of Navier-Stokes equations and an extension of the heat flow to the unit sphere. However it is a very challenging task to obtaining estimates uniform in ε for the full system (1.10), in order to understand the limit $\varepsilon \rightarrow 0$. In [11], the first three authors justified rigorously the zero initial limit $\varepsilon \rightarrow 0$ from hyperbolic to parabolic system when the bulk velocity u is given. In the current paper, the problem we consider appears as a simple instance of this general problem, namely the case where the bulk velocity $u = \mathbf{0}$ and the coefficient $\lambda_1 = -1$ in (1.10). For this case, the system (1.10) is reduced to the scaled wave map (1.8), i.e. the wave map (1.8) with a damping can be regarded as an Ericksen-Leslie’s liquid crystal flow unaffected by the fluid velocity.

1.3. Initial layer and the main result. As mentioned in the previous two subsections the formal limit of the Equation (1.8), obtained by setting $\varepsilon = 0$ is provided by the heat flow for functions with values into \mathbb{S}^2 :

$$\partial_t d_0 = \Delta d_0 + |\nabla d_0|^2 d_0, d_0 \in \mathbb{S}^2, \tag{1.11}$$

The limit we considered is a singular limit, as the character of the equations changes, from a hyperbolic-type system for $\varepsilon > 0$ to a parabolic system for $\varepsilon = 0$. An immediate manifestation of the difference between the two types of equations is related to the initial conditions, which for the limit equation take the form:

$$d_0|_{t=0} = d^{in}(x) \in \mathbb{S}^2. \tag{1.12}$$

Thus, we note that the wave map is a system of hyperbolic equations with two initial conditions, while the heat flow is a parabolic system with only one initial condition. Usually the solution of the heat flow does not satisfy the second initial condition in (1.9). This disparity between the initial conditions of the wave map (1.8) and of the heat flow (1.11) indicates that one should expect an ‘‘initial layer’’ in time, appearing in the limiting process $\varepsilon \rightarrow 0$. A formal derivation (postponed for later, in Section 2) indicates that this should be of the form:

$$\begin{aligned} d_0^I\left(\frac{t}{\varepsilon}, x\right) &= -\varepsilon(\tilde{d}^{in}(x) - \partial_t d_0(0, x)) \exp\left(-\frac{t}{\varepsilon}\right) \\ &= -\varepsilon(\tilde{d}^{in}(x) - \Delta d^{in}(x) - |\nabla d^{in}(x)|^2 d^{in}(x)) \exp\left(-\frac{t}{\varepsilon}\right) \\ &\equiv -\varepsilon D(x) \exp\left(-\frac{t}{\varepsilon}\right), \end{aligned} \tag{1.13}$$

where $D(x)$ is defined as

$$D(x) \equiv \tilde{d}^{in}(x) - \Delta d^{in}(x) - |\nabla d^{in}(x)|^2 d^{in}(x).$$

Our study of the limit from the wave map (1.8) to the heat flow (1.5) is inspired by the classical approach of Caffisch on the compressible Euler limit of the Boltzmann equation [1]. This approach is based on the Hilbert expansion in which the leading term is given by solutions of the limit equation. The Caffisch’s approach assumes that a solution of the limiting equation (which in our case is the heat flow (1.5)) is known beforehand. Then the solution to the original equation (which in our case is the wave map (1.8)) can be constructed around the limiting equation with perturbations as expansions in powers of ε . Based on the arguments above and the formal analysis in Section 2, in the expansions, besides the heat flow, the leading term should also include an initial layer. More specifically, we take the following ansatz of the solution d^ε to the system (1.8):

$$d^\varepsilon(t, x) = d_0(t, x) + d_0^I\left(\frac{t}{\varepsilon}, x\right) + \sqrt{\varepsilon} d_R^\varepsilon(t, x), \tag{1.14}$$

where $d_0(t, x)$ obeys the heat flow (1.11) and the initial layer $d_0^I\left(\frac{t}{\varepsilon}, x\right)$ is defined in (1.13). Plugging (1.14) into the system (1.8), the remainder term $d_R^\varepsilon(t, x)$ must satisfy the system

$$\partial_{tt} d_R^\varepsilon + \frac{1}{\varepsilon} \partial_t d_R^\varepsilon - \frac{1}{\varepsilon} \Delta d_R^\varepsilon = \mathcal{S}(d_R^\varepsilon) + \mathcal{R}(d_R^\varepsilon) \tag{1.15}$$

with the initial conditions

$$d_R^\varepsilon(0, x) = \sqrt{\varepsilon} D(x), \quad \partial_t d_R^\varepsilon(0, x) = 0, \tag{1.16}$$

where the singular term $\mathcal{S}(d_R^\varepsilon)$ is

$$\begin{aligned} \mathcal{S}(d_R^\varepsilon) &= -\frac{1}{\sqrt{\varepsilon}} \partial_{tt} d_0 - \frac{1}{\sqrt{\varepsilon}} \Delta D(x) \exp\left(-\frac{t}{\varepsilon}\right) + \frac{1}{\varepsilon} |\nabla d_0|^2 d_R^\varepsilon + \frac{1}{\sqrt{\varepsilon}} |\nabla d_R^\varepsilon|^2 d_0 \\ &\quad - \frac{1}{\sqrt{\varepsilon}} |\partial_t d_0 + D(x) \exp\left(-\frac{t}{\varepsilon}\right)|^2 d_0 - \frac{1}{\sqrt{\varepsilon}} |\nabla d_0|^2 D(x) \exp\left(-\frac{t}{\varepsilon}\right) \\ &\quad + \frac{2}{\sqrt{\varepsilon}} (\nabla d_0 \cdot \nabla d_R^\varepsilon) d_R^\varepsilon - \frac{2}{\sqrt{\varepsilon}} (\nabla d_0 \cdot \nabla D(x)) \exp\left(-\frac{t}{\varepsilon}\right) d_0 + \frac{2}{\varepsilon} (\nabla d_0 \cdot \nabla d_R^\varepsilon) d_0, \end{aligned}$$

and the regular term $\mathcal{R}(d_R^\varepsilon)$ is

$$\begin{aligned} & \mathcal{R}(d_R^\varepsilon) \\ &= - \left| \partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon}) + \sqrt{\varepsilon} \partial_t d_R^\varepsilon \right|^2 \left[-\sqrt{\varepsilon} D(x) \exp(-\frac{t}{\varepsilon}) + d_R^\varepsilon \right] \\ & \quad - \left[2(\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})) \cdot \partial_t d_R^\varepsilon + \sqrt{\varepsilon} |\partial_t d_R^\varepsilon|^2 \right] d_0 - 2(\nabla D(x) \cdot \nabla d_R^\varepsilon) \exp(-\frac{t}{\varepsilon}) d_0 \\ & \quad + \sqrt{\varepsilon} |\nabla D(x)|^2 \exp(-\frac{2t}{\varepsilon}) d_0 + |\nabla d_R^\varepsilon|^2 d_R^\varepsilon - 2(\nabla d_0 \cdot \nabla D(x)) \exp(-\frac{t}{\varepsilon}) d_R^\varepsilon \\ & \quad - 2(\nabla d_0 \cdot \nabla d_R^\varepsilon) D(x) \exp(-\frac{t}{\varepsilon}) - 2\sqrt{\varepsilon} (\nabla D(x) \cdot \nabla d_R^\varepsilon) \exp(-\frac{t}{\varepsilon}) d_R^\varepsilon \\ & \quad - \sqrt{\varepsilon} |\nabla d_R^\varepsilon|^2 D(x) \exp(-\frac{t}{\varepsilon}) + \varepsilon |\nabla D(x)|^2 \exp(-\frac{2t}{\varepsilon}) d_R^\varepsilon - \varepsilon^{\frac{3}{2}} |\nabla D(x)|^2 \exp(-\frac{3t}{\varepsilon}) D(x) \\ & \quad + 2\sqrt{\varepsilon} (\nabla d_0 \cdot \nabla D(x)) \exp(-\frac{2t}{\varepsilon}) D(x) + 2\varepsilon (\nabla D(x) \cdot \nabla d_R^\varepsilon) \exp(-\frac{2t}{\varepsilon}) D(x). \end{aligned}$$

According to Eells-Sampson’s classical result in [4], for the heat flow (1.11) on the unit sphere \mathbb{S}^2 , one can have the following results of local well-posedness:

PROPOSITION 1.1. *For any given $d^{in} \in \mathbb{S}^2$ satisfying $d^{in} \in \dot{H}^k(\mathbb{R}^3)$ for any integer $k > 2$, there exists a time $T = T(|d^{in}|_{\dot{H}^k}) > 0$ such that (1.11) admits a unique classical solution $d_0 \in L^\infty(0, T; \dot{H}^k) \cap L^2(0, T; \dot{H}^{k+1})$. Moreover, there is a constant $C^* = C^*(|d^{in}|_{\dot{H}^k}, T) > 0$ such that the solution d_0 satisfies*

$$|d_0|_{L^\infty(0, T; \dot{H}^k)}^2 + |\nabla d_0|_{L^2(0, T; H^k)}^2 \leq C^*.$$

The proof can be found in Chapter 5 in the book [15].

Now we state the main result of this paper:

THEOREM 1.1. *We consider vector fields $d^{in} : \mathbb{R}^3 \rightarrow \mathbb{S}^2$ and $\tilde{d}^{in} : \mathbb{R}^3 \rightarrow \mathbb{R}^3$, satisfying the compatibility condition $d^{in} \cdot \tilde{d}^{in} \equiv 0$. Assume that $\nabla d^{in} \in H^6$, $\tilde{d}^{in} \in H^5$, and let $T > 0$ be the time interval of existence of the solution of the heat flow (1.11) with initial condition d^{in} , determined in Proposition 1.1.*

Then, there exists an $\varepsilon_0 \equiv \varepsilon_0(|\nabla d^{in}|_{H^6}, |\tilde{d}^{in}|_{H^5}, T) \in (0, \frac{1}{2})$ such that for all $\varepsilon \in (0, \varepsilon_0)$ we have that on the interval $[0, T]$ the wave map Equation (1.8) with the initial conditions (1.9) admits a unique solution with the form (1.14), i.e.,

$$d^\varepsilon(t, x) = d_0(t, x) + d_0^I(\frac{t}{\varepsilon}, x) + \sqrt{\varepsilon} d_R^\varepsilon(t, x),$$

where d_0 is the solution of the heat flow (1.11) with initial condition d^{in} and $d_0^I(\frac{t}{\varepsilon}, x)$ is the initial layer (1.13). Moreover, there exists a positive constant $C_0 = C_0(d^{in}, \tilde{d}^{in}, T) > 0$, such that the remainder term d_R^ε satisfies the bound

$$|\partial_t d_R^\varepsilon|_{L^\infty(0, T; H^2)}^2 + \frac{1}{\varepsilon} |d_R^\varepsilon|_{L^\infty(0, T; H^3)}^2 \leq C_0 \tag{1.17}$$

for all $\varepsilon \in (0, \varepsilon_0)$.

REMARK 1.1. The rate of convergence we obtain is optimal. Indeed, in order to see this, it suffices to note that the limit we study contains as a particular case the linear scalar case of the singular limit of the damped wave equation to the heat equation.

Indeed, let us consider a solution of the scalar damped wave equation:

$$\varepsilon \partial_{tt} \theta^\varepsilon + \partial_t \theta^\varepsilon = \Delta \theta^\varepsilon \tag{1.18}$$

for $\theta^\varepsilon : \mathbb{R} \times \mathbb{R}^3 \rightarrow \mathbb{R}$ with initial data:

$$\theta^\varepsilon(0, x) = \theta_0(x), \partial_t \theta^\varepsilon(0, x) = \theta_1(x) \tag{1.19}$$

all smooth functions.

Also consider the solution of the heat equation:

$$\partial_t \theta^0 = \Delta \theta^0 \tag{1.20}$$

for $\theta^0 : \mathbb{R} \times \mathbb{R}^3 \rightarrow \mathbb{R}$ with initial data:

$$\theta^0(0, x) = \theta_0(x), \tag{1.21}$$

all smooth functions.

Denoting $n^0(t, x) := (\cos \theta^0(t, x), \sin \theta^0(t, x), 0)$ and $n^\varepsilon(t, x) := (\cos \theta^\varepsilon(t, x), \sin \theta^\varepsilon(t, x), 0)$ we have that n^0 is a solution of the heat-flow:

$$\partial_t n^0 = \Delta n^0 + n^0 |\nabla n^0|^2 \tag{1.22}$$

with initial data

$$n^0(0, x) = n_0(x) := (\cos \theta_0, \sin \theta_0, 0) \tag{1.23}$$

while n^ε is a solution of the wave-map flow:

$$\varepsilon(\partial_{tt} n^\varepsilon + n^\varepsilon |\partial_t n^\varepsilon|^2) + \partial_t n^\varepsilon = \Delta n^\varepsilon + n^\varepsilon |\nabla n^\varepsilon|^2 \tag{1.24}$$

with initial data:

$$\begin{aligned} n^\varepsilon(0, x) &= n_0(x) := (\cos \theta_0(x), \sin \theta_0(x), 0), \\ \partial_t n^\varepsilon(0, x) &= n_1(x) = (-\sin \theta_0(x), \cos \theta_0(x), 0) \theta_1(x). \end{aligned} \tag{1.25}$$

Taking $\theta_1 = \Delta \theta_0$, the claimed optimality of the rate of convergence is shown in [2].

A rigorous justification of the formal expansion (1.14) in the context of classical solutions is provided in this paper. For the original wave map (1.8) with a damping the energy bounds of d^ε uniform in small $\varepsilon > 0$ do not seem available. By taking the expansion (1.14) of the solutions d^ε to the system (1.8) with the initial conditions (1.9), one can yield a remainder system (1.15)-(1.16) of d_R^ε . Although the remainder system (1.15) with the initial data conditions (1.16) is still nonlinear and singular (with singular terms of the type $\frac{1}{\varepsilon}$), it has weaker nonlinearities than the original system (1.8). More precisely, after using the expansion (1.14), the nonlinear term $(-\varepsilon |\partial_t d^\varepsilon|^2 + |\nabla d^\varepsilon|^2) d^\varepsilon$ is replaced by either linear terms (the unknown d^ε is superseded by the known d_0) or a nonlinear term with the same form but with some higher order powers of ε in front. So, by standard energy estimates, we can get uniform bounds in small $\varepsilon > 0$ of the remainder system (1.15)-(1.16).

The organization of this paper is as follows: in next section, we give the formal analysis for the asymptotic behavior of the wave map (1.8) with a damping and initial conditions (1.9) as the inertia density $\varepsilon \rightarrow 0$ by constructing the initial layer $d_0^I(\frac{t}{\varepsilon}, x)$ to deal with the compatibility of the original initial conditions (1.9) and the initial condition of the limit system (1.11). In Section 3, we estimate the uniform energy bounds on small $\varepsilon > 0$ of the remainder system (1.15)-(1.16). Finally, based on the uniform energy estimates in the previous section, Theorem 1.1 of this paper is proved in Section 4.

2. Formal analysis

In this section we present the formal analysis of the limit $\varepsilon \rightarrow 0$ for the damped wave map (1.8) with the initial conditions (1.9). Out of the Equation (1.8) we note that the formal limit, obtaining by setting $\varepsilon = 0$, is the heat flow system (1.11) for functions taking values into \mathbb{S}^2 . We can then naturally take the ansatz

$$d^\varepsilon(t, x) = d_0(t, x) + \hat{d}_R^\varepsilon(t, x), \tag{2.1}$$

where $d_0(t, x)$ is a solution of the heat flow system (1.11) and $\hat{d}_R^\varepsilon(t, x)$ satisfies a hyperbolic system, formally similar to (1.8) but without the geometric constraint of taking values into \mathbb{S}^2 .

If the ansatz (2.1) were reasonable then $\hat{d}_R^\varepsilon(t, x) = O(\varepsilon^\alpha)$ in some sense for some $\alpha > 0$ as $\varepsilon > 0$ is small enough. However, by the second initial condition in (1.9) and the heat flow system (1.11), we know that

$$\begin{aligned} \partial_t \hat{d}_R^\varepsilon(0, x) &= \partial_t d^\varepsilon(0, x) - \partial_t d_0(0, x) \\ &= \tilde{d}^{in}(x) - \partial_t d_0(0, x), \end{aligned}$$

which will not go to 0 as $\varepsilon \rightarrow 0$ for arbitrarily given vectors $\tilde{d}^{in}(x)$ and $d^{in}(x)$. As a consequence, $\hat{d}_R^\varepsilon(t, x) \neq O(\varepsilon^\alpha)$ uniformly in time for any $\alpha > 0$, and then the ansatz (2.1) is not satisfactory.

Therefore, in order to compensate the effect of the initial data, we need to introduce a correction term $d_0^I(\frac{t}{\varepsilon^\beta}, x)$ for some $\beta > 0$ to be determined, called *initial layer*, such that the second initial condition in (1.9) is satisfied by $d_0(t, x) + d_0^I(\frac{t}{\varepsilon^\beta}, x)$ without disturbing too much the first initial condition in (1.9), namely

$$d_0^I(\frac{0}{\varepsilon^\beta}, x) = O(\varepsilon^\alpha)$$

for some $\alpha > 0$ as $\varepsilon \rightarrow 0$. Thus we take the alternative ansatz

$$d^\varepsilon(t, x) = d_0(t, x) + d_0^I(\frac{t}{\varepsilon^\beta}, x) + \sqrt{\varepsilon} d_R^\varepsilon(t, x), \tag{2.2}$$

where the power $\sqrt{\varepsilon}$ in front of the remainder term is motivated by the scaling we chose. This measures the rate of convergence, as it will be shown in the proof of Theorem 1.1.

Recalling that d_0 is a solution of the heat flow (1.11), we plug (2.2) into the system (1.8) and obtain:

$$\begin{aligned} &\varepsilon \partial_{tt}(d_0 + \sqrt{\varepsilon} d_R^\varepsilon) + \sqrt{\varepsilon} \partial_t d_R^\varepsilon + (\varepsilon \partial_{tt} d_0^I + \partial_t d_0^I) - \Delta(d_0 + \sqrt{\varepsilon} d_R^\varepsilon) \\ &= -\varepsilon |\partial_t(d_0 + d_0^I + \sqrt{\varepsilon} d_R^\varepsilon)|^2 (d_0 + d_0^I + \sqrt{\varepsilon} d_R^\varepsilon) \\ &\quad + |\nabla(d_0 + d_0^I + \sqrt{\varepsilon} d_R^\varepsilon)|^2 (d_0^I + \sqrt{\varepsilon} d_R^\varepsilon) + [2\nabla d_0 \cdot \nabla(d_0^I + \sqrt{\varepsilon} d_R^\varepsilon) + |\nabla(d_0^I + \sqrt{\varepsilon} d_R^\varepsilon)|^2] d_0. \end{aligned}$$

Then we construct the initial layer in order to cancel certain time-dependent terms in the previous equations and to accommodate the discrepancy in the initial data, namely we take d_0^I satisfying the x -dependent ODE and the initial-data condition:

$$\begin{cases} \varepsilon \partial_{tt} d_0^I(\frac{t}{\varepsilon^\beta}, x) + \partial_t d_0^I(\frac{t}{\varepsilon^\beta}, x) = 0, \\ \partial_t d_0^I(\frac{0}{\varepsilon^\beta}, x) = \tilde{d}^{in}(x) - \partial_t d_0(0, x). \end{cases}$$

Furthermore, since d_0^I is an initial layer, the following condition at infinity is required:

$$d_0^I(\infty, x) = \lim_{y \rightarrow \infty} d_0^I(y, x) = 0. \tag{2.3}$$

By solving the ODE system with the given boundary conditions we have

$$d_0^I\left(\frac{t}{\varepsilon^\beta}, x\right) = -\varepsilon(\tilde{d}^{in}(x) - \partial_t d_0(0, x)) \exp\left(-\frac{t}{\varepsilon}\right). \tag{2.4}$$

We remark that the initial layer $d_0^I\left(\frac{t}{\varepsilon^\beta}, x\right)$ in (2.4) is, in fact, independent of $\beta > 0$ and $d_0^I\left(\frac{0}{\varepsilon^\beta}, x\right) = -\varepsilon(\tilde{d}^{in}(x) - \partial_t d_0(0, x)) \rightarrow 0$ for any given $\tilde{d}^{in}(x)$ and $d^{in}(x)$ as $\varepsilon \rightarrow 0$. Consequently, the ansatz (2.2) is reasonable.

Without loss of generality, we take $\beta=1$ in the ansatz (2.2). Thus, by substituting (1.11) and (1.14) into the system (1.8), we derive the equation satisfied by the remainder $d_R^\varepsilon(t, x)$ as follows:

$$\begin{aligned} & \varepsilon^{\frac{3}{2}} \partial_{tt} d_R^\varepsilon + \varepsilon \partial_{tt} d_0 + \sqrt{\varepsilon} \partial_t d_R^\varepsilon - \sqrt{\varepsilon} \Delta d_R^\varepsilon + \varepsilon \Delta D(x) \exp\left(-\frac{t}{\varepsilon}\right) \\ &= -\varepsilon \left| \partial_t \left(d_0 - \varepsilon D(x) \exp\left(-\frac{t}{\varepsilon}\right) + \sqrt{\varepsilon} d_R^\varepsilon \right) \right|^2 \left(d_0 - \varepsilon D(x) \exp\left(-\frac{t}{\varepsilon}\right) + \sqrt{\varepsilon} d_R^\varepsilon \right) \\ & \quad + \left| \nabla \left(d_0 - \varepsilon D(x) \exp\left(-\frac{t}{\varepsilon}\right) + \sqrt{\varepsilon} d_R^\varepsilon \right) \right|^2 \left(-\varepsilon D(x) \exp\left(-\frac{t}{\varepsilon}\right) + \sqrt{\varepsilon} d_R^\varepsilon \right) \\ & \quad + \left[2 \nabla d_0 \cdot \nabla \left(-\varepsilon D(x) \exp\left(-\frac{t}{\varepsilon}\right) + \sqrt{\varepsilon} d_R^\varepsilon \right) + \left| \nabla \left(-\varepsilon D(x) \exp\left(-\frac{t}{\varepsilon}\right) + \sqrt{\varepsilon} d_R^\varepsilon \right) \right|^2 \right] d_0, \end{aligned}$$

which, after multiplication by $\varepsilon^{-\frac{3}{2}}$ is the Equation (1.15) we used before.

3. Uniform energy estimates

In this section, we will provide, by energy methods, bounds that are uniform with respect to small inertia constant $\varepsilon > 0$, for the remainder system (1.15)-(1.16). By Proposition 1.1, the d_0 , which obeys the heat flow (1.11) into the unit sphere S^2 is regarded as a known quantity in the remainder system (1.15)-(1.16).

To conveniently state our results, we need to introduce the following energy functionals:

$$\begin{aligned} E_\varepsilon(t) &= |\partial_t d_R^\varepsilon|_{H^2}^2 + \left(\frac{1}{\varepsilon} - 1\right) |d_R^\varepsilon|_{H^2}^2 + \frac{2}{\varepsilon} |\nabla d_R^\varepsilon|_{H^2}^2 + |\partial_t d_R^\varepsilon + d_R^\varepsilon|_{H^2}^2, \\ F_\varepsilon(t) &= \left(\frac{1}{\varepsilon} - \frac{1}{2}\right) |\partial_t d_R^\varepsilon|_{H^2}^2 + \frac{1}{2\varepsilon} |\nabla d_R^\varepsilon|_{H^2}^2. \end{aligned}$$

The following lemma provides the claimed uniform energy estimates:

LEMMA 3.1. *Let $d^{in} \in \dot{H}^7(\mathbb{R}^3; S^2)$ and $[0, T]$ be the interval of existence of the solution of the heat flow with initial data d^{in} , as provided in Proposition 1.1.*

For $\tilde{d}^{in} \in H^5$ assume that there exists a $\varepsilon_0 \equiv \varepsilon_0(|\nabla d^{in}|_{H^6}, |\tilde{d}^{in}|_{H^5}, T) \in (0, \frac{1}{2})$ such that for all $\varepsilon \in (0, \varepsilon_0)$ we deduce that $d_R^\varepsilon \in L^\infty([0, T]; H^3)$ and $\partial_t d^\varepsilon \in L^\infty([0, T]; H^2)$ is a solution to the remainder system (1.15)-(1.16). Then there exists a positive constant $C = C(|\nabla d^{in}|_{H^6}, |\tilde{d}^{in}|_{H^5}, T) > 0$ such that the inequality

$$\frac{d}{dt} E_\varepsilon(t) + 3F_\varepsilon(t) \leq C[1 + E_\varepsilon(t)][1 + \varepsilon E_\varepsilon(t)] \tag{3.1}$$

holds for all $\varepsilon \in (0, \frac{1}{2})$ and $t \in [0, T]$.

We remark that the condition $0 < \varepsilon < \frac{1}{2}$ guarantees the relation

$$\frac{1}{2\varepsilon} < \frac{1}{\varepsilon} - 1 < \frac{1}{\varepsilon},$$

which makes the energy functionals $E_\varepsilon(t)$ and $F_\varepsilon(t)$ non-negative. Since our goal is to rigorously analyze the asymptotic behavior as $\varepsilon \rightarrow 0$ for the wave map (1.8)-(1.9), the condition $0 < \varepsilon < \frac{1}{2}$ is sufficient.

Proof. For the convenience of notations, we rewrite the singular terms of the remainder system (1.15) as

$$\begin{aligned}
 \mathcal{S}(d_R^\varepsilon) = & \left. \begin{aligned} & -\frac{1}{\sqrt{\varepsilon}} \left(\partial_{tt} d_0 + \Delta D(x) \exp(-\frac{t}{\varepsilon}) + |\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})|^2 d_0 \right. \\ & \left. + |\nabla d_0|^2 D(x) \exp(-\frac{t}{\varepsilon}) + 2\nabla d_0 \cdot \nabla D(x) \exp(-\frac{t}{\varepsilon}) d_0 \right) \end{aligned} \right\} \mathcal{S}_1 \\
 & \left. \begin{aligned} & + \frac{1}{\sqrt{\varepsilon}} (2(\nabla d_0 \cdot \nabla d_R^\varepsilon) d_R^\varepsilon + |\nabla d_R^\varepsilon|^2 d_0) \end{aligned} \right\} \mathcal{S}_2 \\
 & \left. \begin{aligned} & + \frac{1}{\varepsilon} [|\nabla d_0|^2 d_R^\varepsilon + 2(\nabla d_0 \cdot \nabla d_R^\varepsilon) d_0] \end{aligned} \right\} \mathcal{S}_3 \\
 \triangleq & \mathcal{S}_1 + \mathcal{S}_2 + \mathcal{S}_3
 \end{aligned} \tag{3.2}$$

and the regular terms as

$$\begin{aligned}
 \mathcal{R}(d_R^\varepsilon) = & \left. \begin{aligned} & \sqrt{\varepsilon} |\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})|^2 D(x) \exp(-\frac{t}{\varepsilon}) + \sqrt{\varepsilon} |\nabla D(x)|^2 \exp(-\frac{2t}{\varepsilon}) d_0 \\ & - \varepsilon^{\frac{3}{2}} |\nabla D(x)|^2 \exp(-\frac{3t}{\varepsilon}) D(x) + 2\sqrt{\varepsilon} (\nabla d_0 \cdot \nabla D(x)) \exp(-\frac{2t}{\varepsilon}) D(x) \end{aligned} \right\} \mathcal{R}_1 \\
 & \left. \begin{aligned} & - 2(\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})) \cdot \partial_t d_R^\varepsilon d_0 - |\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})|^2 d_R^\varepsilon \\ & + 2\varepsilon (\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})) \cdot \partial_t d_R^\varepsilon D(x) \exp(-\frac{t}{\varepsilon}) + \varepsilon |\nabla D|^2 \exp(-\frac{2t}{\varepsilon}) d_R^\varepsilon \\ & - 2(\nabla d_0 \cdot \nabla D(x)) \exp(-\frac{t}{\varepsilon}) d_R^\varepsilon - 2(\nabla d_0 \cdot \nabla d_R^\varepsilon) D(x) \exp(-\frac{t}{\varepsilon}) \\ & - 2(\nabla D(x) \cdot \nabla d_R^\varepsilon) \exp(-\frac{t}{\varepsilon}) d_0 + 2\varepsilon (\nabla D(x) \cdot \nabla d_R^\varepsilon) \exp(-\frac{2t}{\varepsilon}) D(x) \end{aligned} \right\} \mathcal{R}_2 \\
 & \left. \begin{aligned} & + \varepsilon^{\frac{3}{2}} D(x) \exp(-\frac{t}{\varepsilon}) |\partial_t d_R^\varepsilon|^2 - \sqrt{\varepsilon} |\partial_t d_R^\varepsilon|^2 d_0 \\ & - 2\sqrt{\varepsilon} (\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})) \cdot \partial_t d_R^\varepsilon d_R^\varepsilon \\ & - 2\sqrt{\varepsilon} (\nabla D(x) \cdot \nabla d_R^\varepsilon) \exp(-\frac{t}{\varepsilon}) d_R^\varepsilon - \sqrt{\varepsilon} |\nabla d_R^\varepsilon|^2 D(x) \exp(-\frac{t}{\varepsilon}) \end{aligned} \right\} \mathcal{R}_3 \\
 & \left. \begin{aligned} & - \varepsilon |\partial_t d_R^\varepsilon|^2 d_R^\varepsilon + |\nabla d_R^\varepsilon|^2 d_R^\varepsilon \end{aligned} \right\} \mathcal{R}_4 \\
 \triangleq & \mathcal{R}_1 + \mathcal{R}_2 + \mathcal{R}_3 + \mathcal{R}_4.
 \end{aligned} \tag{3.3}$$

Step 1. L^2 -estimates. Multiplying the remainder Equation (1.15) by $\partial_t d_R^\varepsilon$, integrating over \mathbb{R}^3 and by parts, we obtain the following equation:

$$\frac{1}{2} \frac{d}{dt} (|\partial_t d_R^\varepsilon|_{L^2}^2 + \frac{1}{\varepsilon} |\nabla d_R^\varepsilon|_{L^2}^2) + \frac{1}{\varepsilon} |\partial_t d_R^\varepsilon|_{L^2}^2 = \langle \mathcal{S}(d_R^\varepsilon), \partial_t d_R^\varepsilon \rangle + \langle \mathcal{R}(d_R^\varepsilon), \partial_t d_R^\varepsilon \rangle. \tag{3.4}$$

(I) *Estimates for the singular terms $\langle \mathcal{S}(d_R^\varepsilon), \partial_t d_R^\varepsilon \rangle$:*

For estimating $\langle \mathcal{S}_1, \partial_t d_R^\varepsilon \rangle$, we use the Hölder inequality, the Sobolev embedding theorems, the facts that $\exp(-\frac{t}{\varepsilon}) \leq 1$ and $|d_0| = 1$ to obtain:

$$\frac{1}{\sqrt{\varepsilon}} \langle \partial_{tt} d_0 + \Delta D(x) \exp(-\frac{t}{\varepsilon}), \partial_t d_R^\varepsilon \rangle \lesssim (|\partial_{tt} d_0|_{L_t^\infty L_x^2} + |\Delta D(x)|_{L^2}) \left| \frac{\partial_t d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2}$$

and

$$\begin{aligned}
 & \frac{1}{\sqrt{\varepsilon}} \langle (|\nabla d_0|^2 D(x) \exp(-\frac{t}{\varepsilon}) + 2\nabla d_0 \cdot \nabla D(x) \exp(-\frac{t}{\varepsilon}) d_0), \partial_t d_R^\varepsilon \rangle \\
 & \lesssim (|\nabla d_0|_{L^\infty}^2 |D(x)|_{L^2} + |\nabla d_0|_{L^\infty} |\nabla D(x)|_{L^2}) \left| \frac{\partial_t d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2} \\
 & \lesssim (|\nabla d_0|_{L_t^\infty H_x^2}^2 + |\nabla d_0|_{L_t^\infty H_x^2}) |D(x)|_{H^1} \left| \frac{\partial_t d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2}.
 \end{aligned}$$

Similarly:

$$\begin{aligned}
 \frac{1}{\sqrt{\varepsilon}} \langle |\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})|^2 d_0, \partial_t d_R^\varepsilon \rangle & \lesssim |\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})|_{L^4}^2 \left| \frac{\partial_t d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2} \\
 & \lesssim (|\partial_t d_0|_{L_t^\infty L_x^4}^2 + |D(x)|_{L^4}^2) \left| \frac{\partial_t d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2} \\
 & \lesssim (|\partial_t d_0|_{L_t^\infty H_x^1}^2 + |D(x)|_{H^1}^2) \left| \frac{\partial_t d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2}.
 \end{aligned}$$

Summarizing, we estimate $\langle \mathcal{S}_1, \partial_t d_R^\varepsilon \rangle$ as follows:

$$\begin{aligned} \langle \mathcal{S}_1, \partial_t d_R^\varepsilon \rangle &\lesssim \left[|\partial_{tt} d_0|_{L_t^\infty L_x^2} + |\partial_t d_0|_{L_t^\infty H_x^1}^2 + |D(x)|_{H^1}^2 \right. \\ &\quad \left. + (|\nabla d_0|_{L_t^\infty H_x^2}^2 + 1) |D(x)|_{H^2} \right] \left| \frac{\partial_t d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L_x^2}. \end{aligned} \quad (3.5)$$

It is easy to derive the estimates of $\langle \mathcal{S}_2, \partial_t d_R^\varepsilon \rangle$ and $\langle \mathcal{S}_3, \partial_t d_R^\varepsilon \rangle$ as follows:

$$\begin{aligned} \langle \mathcal{S}_2, \partial_t d_R^\varepsilon \rangle &= \frac{1}{\sqrt{\varepsilon}} \langle 2(\nabla d_0 \cdot \nabla d_R^\varepsilon) d_R^\varepsilon + |\nabla d_R^\varepsilon|^2 d_0, \partial_t d_R^\varepsilon \rangle \\ &\lesssim (|\nabla d_0|_{L^\infty} |d_R^\varepsilon|_{L^\infty} |\nabla d_R^\varepsilon|_{L^2} + |\nabla d_R^\varepsilon|_{L^4}^2) \left| \frac{\partial_t d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2} \\ &\lesssim (|\nabla d_0|_{L_t^\infty H_x^2} + 1) |d_R^\varepsilon|_{H^2}^2 \left| \frac{\partial_t d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2} \end{aligned} \quad (3.6)$$

and

$$\begin{aligned} \langle \mathcal{S}_3, \partial_t d_R^\varepsilon \rangle &= \frac{1}{\varepsilon} \langle |\nabla d_0|^2 d_R^\varepsilon + 2(\nabla d_0 \cdot \nabla d_R^\varepsilon) d_0, \partial_t d_R^\varepsilon \rangle \\ &\lesssim (|\nabla d_0|_{L^\infty}^2 \left| \frac{d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2} + |\nabla d_0|_{L^\infty} \left| \frac{\nabla d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2}) \left| \frac{\partial_t d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2} \\ &\lesssim (|\nabla d_0|_{L_t^\infty H_x^2} + |\nabla d_0|_{L_t^\infty H_x^2}^2) \left(\left| \frac{d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2} + \left| \frac{\nabla d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2} \right) \left| \frac{\partial_t d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2}. \end{aligned} \quad (3.7)$$

Hence we have the estimate of $\langle \mathcal{S}(d_R^\varepsilon), \partial_t d_R^\varepsilon \rangle$ by combining the inequalities (3.5), (3.6) and (3.7):

$$\langle \mathcal{S}(d_R^\varepsilon), \partial_t d_R^\varepsilon \rangle \leq C_{11} \left(1 + |d_R^\varepsilon|_{H^2}^2 + \left| \frac{d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2} + \left| \frac{\nabla d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2} \right) \left| \frac{\partial_t d_R^\varepsilon}{\sqrt{\varepsilon}} \right|_{L^2}, \quad (3.8)$$

where the constant

$$C_{11} = C \left[|\partial_{tt} d_0|_{L_t^\infty L_x^2} + |\partial_t d_0|_{L_t^\infty H_x^1}^2 + |D(x)|_{H^1}^2 + (|\nabla d_0|_{L_t^\infty H_x^2}^2 + 1) (|D(x)|_{H^2} + 1) \right] > 0$$

for some computable positive constant $C > 0$.

(II) *Estimates for the regular terms $\langle \mathcal{R}(d_R^\varepsilon), \partial_t d_R^\varepsilon \rangle$:*

We have divided the regular terms $\mathcal{R}(d_R^\varepsilon)$ into four parts, which we will estimate separately.

The estimate of $\langle \mathcal{R}_1, \partial_t d_R^\varepsilon \rangle$. By the Hölder inequality and the Sobolev embedding theorems we have

$$\begin{aligned} &\sqrt{\varepsilon} \langle |\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})|^2 D(x) \exp(-\frac{t}{\varepsilon}), \partial_t d_R^\varepsilon \rangle \\ &\lesssim \sqrt{\varepsilon} |\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})|_{L^\infty}^2 |D(x)|_{L^2} |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim \sqrt{\varepsilon} (|\partial_t d_0|_{L_t^\infty H_x^2}^2 + |D(x)|_{H^2}^2) |D(x)|_{L^2} |\partial_t d_R^\varepsilon|_{L^2} \end{aligned}$$

and

$$\begin{aligned} \sqrt{\varepsilon} \langle (\nabla d_0 \nabla D(x)) \exp(-\frac{2t}{\varepsilon}) D(x), \partial_t d_R^\varepsilon \rangle &\lesssim \sqrt{\varepsilon} |\nabla d_0|_{L^\infty} |D(x)|_{L^\infty} |\nabla D(x)|_{L^2} |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim \sqrt{\varepsilon} |\nabla d_0|_{L_t^\infty H_x^2} |D(x)|_{H^2}^2 |\partial_t d_R^\varepsilon|_{L^2}, \end{aligned}$$

where we have used the fact that $\exp(-\frac{t}{\varepsilon}) \leq 1$.

The other two terms in $\langle \mathcal{R}_1, \partial_t d_R^\varepsilon \rangle$ are similarly estimated, as follows:

$$\begin{aligned} \sqrt{\varepsilon} \langle |\nabla D(x)|^2 \exp(-\frac{2t}{\varepsilon}) d_0, \partial_t d_R^\varepsilon \rangle &\lesssim \sqrt{\varepsilon} |\nabla D(x)|_{L^4}^2 |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim \sqrt{\varepsilon} |D(x)|_{H^2}^2 |\partial_t d_R^\varepsilon|_{L^2}, \\ \varepsilon^{\frac{3}{2}} \langle |\nabla D(x)|^2 \exp(-\frac{3t}{\varepsilon}) D(x), \partial_t d_R^\varepsilon \rangle &\lesssim \varepsilon^{\frac{3}{2}} |D(x)|_{L^\infty} |\nabla D(x)|_{L^4}^2 |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim \varepsilon^{\frac{3}{2}} |D(x)|_{H^2}^3 |\partial_t d_R^\varepsilon|_{L^2}. \end{aligned}$$

For any small enough ε such that $\varepsilon \in (0, \frac{1}{2}]$, we have $\varepsilon^{\frac{3}{2}} \leq \varepsilon^{\frac{1}{2}}$. So from the above inequalities, we obtain the following estimate:

$$\langle \mathcal{R}_1, \partial_t d_R^\varepsilon \rangle \lesssim \sqrt{\varepsilon} [(|\partial_t d_0|_{L_t^\infty H_x^2}^2 + |\nabla d_0|_{L_t^\infty H_x^2}^2 + |D(x)|_{H^2}^2) |D(x)|_{H^2} + |D(x)|_{H^2}^2] |\partial_t d_R^\varepsilon|_{L^2}. \tag{3.9}$$

The estimate of $\langle \mathcal{R}_2, \partial_t d_R^\varepsilon \rangle$. We have the following estimates for the first three terms in $\langle \mathcal{R}_2, \partial_t d_R^\varepsilon \rangle$:

$$\begin{aligned} (|\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})|^2 d_R^\varepsilon, \partial_t d_R^\varepsilon) &\lesssim |\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})|_{L^\infty}^2 |d_R^\varepsilon|_{L^2} |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim (|\partial_t d_0|_{L_t^\infty H_x^2}^2 + |D(x)|_{H^2}^2) |d_R^\varepsilon|_{L^2} |\partial_t d_R^\varepsilon|_{L^2}, \\ \langle (\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})) \cdot \partial_t d_R^\varepsilon d_0, \partial_t d_R^\varepsilon \rangle &\lesssim (|\partial_t d_0|_{L^\infty} + |D(x)|_{L^\infty}) |\partial_t d_R^\varepsilon|_{L^2}^2 \\ &\lesssim (|\partial_t d_0|_{L_t^\infty H_x^2} + |D(x)|_{H^2}) |\partial_t d_R^\varepsilon|_{L^2}^2 \end{aligned}$$

and

$$\begin{aligned} \varepsilon \langle (\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})) \cdot D(x) \exp(-\frac{t}{\varepsilon}) \partial_t d_R^\varepsilon, \partial_t d_R^\varepsilon \rangle \\ \lesssim \varepsilon |\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})|_{L^\infty} |D(x)|_{L^\infty} |\partial_t d_R^\varepsilon|_{L^2}^2 \\ \lesssim \varepsilon (|\partial_t d_0|_{L_t^\infty H_x^2} + |D(x)|_{H^2}) |D(x)|_{H^2} |\partial_t d_R^\varepsilon|_{L^2}^2, \end{aligned}$$

where we have used the Hölder inequality, the Sobolev embedding theorems and the fact that $|d_0| = 1$.

As for the following three terms, one can easily obtain:

$$\begin{aligned} \varepsilon \langle |\nabla D(x)|^2 \exp(-\frac{2t}{\varepsilon}) d_R^\varepsilon, \partial_t d_R^\varepsilon \rangle &\lesssim \varepsilon |\nabla D(x)|_{L^\infty}^2 |d_R^\varepsilon|_{L^2} |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim \varepsilon |D(x)|_{H^3}^2 |d_R^\varepsilon|_{L^2} |\partial_t d_R^\varepsilon|_{L^2}, \\ \langle (\nabla d_0 \cdot \nabla D(x)) \exp(-\frac{t}{\varepsilon}) d_R^\varepsilon, \partial_t d_R^\varepsilon \rangle &\lesssim |\nabla d_0|_{L^\infty} |\nabla D(x)|_{L^\infty} |d_R^\varepsilon|_{L^2} |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim |\nabla d_0|_{L_t^\infty H_x^2} |D(x)|_{H^3} |d_R^\varepsilon|_{L^2} |\partial_t d_R^\varepsilon|_{L^2}, \\ \langle (\nabla d_0 \cdot \nabla d_R^\varepsilon) D(x) \exp(-\frac{t}{\varepsilon}), \partial_t d_R^\varepsilon \rangle &\lesssim |\nabla d_0|_{L^\infty} |D(x)|_{L^\infty} |\nabla d_R^\varepsilon|_{L^2} |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim |\nabla d_0|_{L_t^\infty H_x^2} |D(x)|_{H^2} |\nabla d_R^\varepsilon|_{L^2} |\partial_t d_R^\varepsilon|_{L^2}, \end{aligned}$$

where we have used the Hölder inequality, the Sobolev embedding theorems and the bound $\exp(-\frac{t}{\varepsilon}) \leq 1$.

Similarly as before, we estimate the last two terms, as follows:

$$\begin{aligned} \langle (\nabla D(x) \cdot \nabla d_R^\varepsilon) \exp(-\frac{t}{\varepsilon}) d_0, \partial_t d_R^\varepsilon \rangle &\lesssim |\nabla D(x)|_{L^\infty} |\nabla d_R^\varepsilon|_{L^2} |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim |D(x)|_{H^3} |\nabla d_R^\varepsilon|_{L^2} |\partial_t d_R^\varepsilon|_{L^2}, \\ \varepsilon \langle (\nabla D(x) \cdot \nabla d_R^\varepsilon) \exp(-\frac{2t}{\varepsilon}) D(x), \partial_t d_R^\varepsilon \rangle &\lesssim \varepsilon |\nabla D(x)|_{L^\infty} |D(x)|_{L^\infty} |\nabla d_R^\varepsilon|_{L^2} |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim \varepsilon |D(x)|_{H^3}^2 |\nabla d_R^\varepsilon|_{L^2} |\partial_t d_R^\varepsilon|_{L^2}. \end{aligned}$$

Combining the above estimates and using that $\varepsilon \in (0, \frac{1}{2}]$, we get

$$\begin{aligned} \langle \mathcal{R}_2, \partial_t d_R^\varepsilon \rangle &\lesssim (1 + |\partial_t d_0|_{L_t^\infty H_x^2} + |\nabla d_0|_{L_t^\infty H_x^2} + |D(x)|_{H^3}) (|\partial_t d_0|_{L_t^\infty H_x^2} + |D(x)|_{H^3}) \\ &\quad \times (|d_R^\varepsilon|_{L^2} + |\nabla d_R^\varepsilon|_{L^2} + |\partial_t d_R^\varepsilon|_{L^2}) |\partial_t d_R^\varepsilon|_{L^2}. \end{aligned} \tag{3.10}$$

The estimate of $\langle \mathcal{R}_3, \partial_t d_R^\varepsilon \rangle$. One can easily derive the following estimates

$$\begin{aligned} \sqrt{\varepsilon} \langle (\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})) \cdot \partial_t d_R^\varepsilon d_R^\varepsilon, \partial_t d_R^\varepsilon \rangle &\lesssim \sqrt{\varepsilon} (|\partial_t d_0|_{L_t^\infty H_x^2} + |D(x)|_{H^2}) |d_R^\varepsilon|_{H^2} |\partial_t d_R^\varepsilon|_{L^2}^2, \\ \sqrt{\varepsilon} \langle (\nabla D(x) \cdot \nabla d_R^\varepsilon) \exp(-\frac{t}{\varepsilon}) d_R^\varepsilon, \partial_t d_R^\varepsilon \rangle &\lesssim \sqrt{\varepsilon} |D(x)|_{H^3} |d_R^\varepsilon|_{H^2} |\nabla d_R^\varepsilon|_{L^2} |\partial_t d_R^\varepsilon|_{L^2}, \end{aligned}$$

by using the Hölder inequality and the Sobolev embedding theorems. Recalling that $|d_0|=1$ and using the Sobolev embeddings $H^1 \hookrightarrow L^4$ and $H^2 \hookrightarrow L^\infty$ we get

$$\begin{aligned} \varepsilon^{\frac{3}{2}} \langle D(x) \exp(-\frac{t}{\varepsilon}) |\partial_t d_R^\varepsilon|^2, \partial_t d_R^\varepsilon \rangle &\lesssim \varepsilon^{\frac{3}{2}} |D(x)|_{L^\infty} |\partial_t d_R^\varepsilon|_{L^4}^2 |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim \varepsilon^{\frac{3}{2}} |D(x)|_{H^2} |\partial_t d_R^\varepsilon|_{H^1}^2 |\partial_t d_R^\varepsilon|_{L^2}, \\ \sqrt{\varepsilon} \langle |\partial_t d_R^\varepsilon|^2 d_0, \partial_t d_R^\varepsilon \rangle &\lesssim \sqrt{\varepsilon} |\partial_t d_R^\varepsilon|_{L^4}^2 |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim \sqrt{\varepsilon} |\partial_t d_R^\varepsilon|_{H^1}^2 |\partial_t d_R^\varepsilon|_{L^2}, \\ \sqrt{\varepsilon} \langle |\nabla d_R^\varepsilon|^2 D(x) \exp(-\frac{t}{\varepsilon}), \partial_t d_R^\varepsilon \rangle &\lesssim \sqrt{\varepsilon} |D(x)|_{L^\infty} |\nabla d_R^\varepsilon|_{L^4}^2 |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim \sqrt{\varepsilon} |D(x)|_{H^2} |\nabla d_R^\varepsilon|_{H^1}^2 |\partial_t d_R^\varepsilon|_{L^2}, \end{aligned}$$

The above estimates immediately give the bound on $\langle \mathcal{R}_3, \partial_t d_R^\varepsilon \rangle$:

$$\begin{aligned} \langle \mathcal{R}_3, \partial_t d_R^\varepsilon \rangle &\lesssim \sqrt{\varepsilon} (1 + |\partial_t d_0|_{L_t^\infty H_x^2} + |D(x)|_{H^3}) (|\partial_t d_R^\varepsilon|_{H^1}^2 + |\nabla d_R^\varepsilon|_{H^1}^2 \\ &\quad + |d_R^\varepsilon|_{H^2} |\nabla d_R^\varepsilon|_{L^2} + |d_R^\varepsilon|_{H^2} |\partial_t d_R^\varepsilon|_{L^2}) |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim \sqrt{\varepsilon} (1 + |\partial_t d_0|_{L_t^\infty H_x^2} + |D(x)|_{H^3}) (|\partial_t d_R^\varepsilon|_{H^1}^2 + |d_R^\varepsilon|_{H^2}^2) |\partial_t d_R^\varepsilon|_{L^2}. \end{aligned} \tag{3.11}$$

The estimate of $\langle \mathcal{R}_4, \partial_t d_R^\varepsilon \rangle$. The first term in $\langle \mathcal{R}_4, \partial_t d_R^\varepsilon \rangle$ can be bounded as

$$\begin{aligned} \varepsilon \langle |\partial_t d_R^\varepsilon|^2 d_R^\varepsilon, \partial_t d_R^\varepsilon \rangle &\leq \varepsilon |d_R^\varepsilon|_{L^\infty} |\partial_t d_R^\varepsilon|_{L^4}^2 |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim \varepsilon |d_R^\varepsilon|_{H^2} |\partial_t d_R^\varepsilon|_{H^1}^2 |\partial_t d_R^\varepsilon|_{L^2} \end{aligned}$$

by using the Hölder inequality and the Sobolev embedding theorems. The other term can be bounded in a similar way:

$$\begin{aligned} \langle |\nabla d_R^\varepsilon|^2 d_R^\varepsilon, \partial_t d_R^\varepsilon \rangle &\lesssim |d_R^\varepsilon|_{L^\infty} |\nabla d_R^\varepsilon|_{L^4}^2 |\partial_t d_R^\varepsilon|_{L^2} \\ &\lesssim |d_R^\varepsilon|_{H^2} |\nabla d_R^\varepsilon|_{H^1}^2 |\partial_t d_R^\varepsilon|_{L^2}. \end{aligned}$$

Hence we obtain the estimate of $\langle \mathcal{R}_4, \partial_t d_R^\varepsilon \rangle$ as follows:

$$\langle \mathcal{R}_4, \partial_t d_R^\varepsilon \rangle \lesssim (\varepsilon |\partial_t d_R^\varepsilon|_{H^1}^2 + |\nabla d_R^\varepsilon|_{H^1}^2) |d_R^\varepsilon|_{H^2} |\partial_t d_R^\varepsilon|_{L^2}. \tag{3.12}$$

Summing up the inequalities (3.9)-(3.12), we get

$$\begin{aligned} \langle \mathcal{R}(d_R^\varepsilon), \partial_t d_R^\varepsilon \rangle &\leq C_{12} \left[\sqrt{\varepsilon} + |d_R^\varepsilon|_{H^1} + |\partial_t d_R^\varepsilon|_{L^2} + \sqrt{\varepsilon} |\partial_t d_R^\varepsilon|_{H^1}^2 + \sqrt{\varepsilon} |d_R^\varepsilon|_{H^2}^2 \right. \\ &\quad \left. + |d_R^\varepsilon|_{H^2} (\varepsilon |\partial_t d_R^\varepsilon|_{H^1}^2 + |\nabla d_R^\varepsilon|_{H^1}^2) \right] |\partial_t d_R^\varepsilon|_{L^2}, \end{aligned} \tag{3.13}$$

where

$$C_{12} = C (1 + |\partial_t d_0|_{L_t^\infty H_x^2}^2 + |\nabla d_0|_{L_t^\infty H_x^2}^2 + |D(x)|_{H^3}^2) (1 + |\partial_t d_0|_{L_t^\infty H_x^2} + |D(x)|_{H^3}) > 0,$$

and C is a positive computable constant.

Therefore, plugging the estimates (3.8) and (3.13) into the equality (3.4), we have

$$\begin{aligned} &\frac{1}{2} \frac{d}{dt} (|\partial_t d_R^\varepsilon|_{L^2}^2 + \frac{1}{\varepsilon} |\nabla d_R^\varepsilon|_{L^2}^2) + \frac{1}{\varepsilon} |\partial_t d_R^\varepsilon|_{L^2}^2 \\ &\leq C_1 \left\{ (1 + |d_R^\varepsilon|_{H^2}^2 + |\frac{d_R^\varepsilon}{\sqrt{\varepsilon}}|_{L^2} + |\frac{\nabla d_R^\varepsilon}{\sqrt{\varepsilon}}|_{L^2}) |\frac{\partial_t d_R^\varepsilon}{\sqrt{\varepsilon}}|_{L^2} + \left[\sqrt{\varepsilon} + |d_R^\varepsilon|_{H^1} + |\partial_t d_R^\varepsilon|_{L^2} \right. \right. \\ &\quad \left. \left. + \sqrt{\varepsilon} |\partial_t d_R^\varepsilon|_{H^1}^2 + \sqrt{\varepsilon} |d_R^\varepsilon|_{H^2}^2 + |d_R^\varepsilon|_{H^2} (\varepsilon |\partial_t d_R^\varepsilon|_{H^1}^2 + |\nabla d_R^\varepsilon|_{H^1}^2) \right] |\partial_t d_R^\varepsilon|_{L^2} \right\}, \end{aligned} \tag{3.14}$$

where the constant

$$C_1 = C(1 + |D(x)|_{H^3}^3 + |\partial_{tt}d_0|_{L^\infty(0,T;L^2)} + |\partial_t d_0|_{L^\infty(0,T;H^2)}^3 + |\nabla d_0|_{L^\infty(0,T;H^2)}^4) > 0$$

for some computable positive constant C .

(III) *Estimates of the norm $|d_R^\epsilon|_{L^2}$:*

Observing that the norm $|d_R^\epsilon|_{L^2}$ appearing on the right-hand side of (3.14) is not yet controlled, we need additional work to estimate $|d_R^\epsilon|_{L^2}$. In order to do this it is natural to multiply the equation of the reminder term (1.15) by d_R^ϵ , integrate over \mathbb{R}^3 and by parts, and use the identity:

$$\begin{aligned} \langle \partial_{tt}d_R^\epsilon, d_R^\epsilon \rangle &= \frac{d}{dt} \langle \partial_t d_R^\epsilon, d_R^\epsilon \rangle - |\partial_t d_R^\epsilon|_{L^2}^2 \\ &= \frac{1}{2} \frac{d}{dt} (|\partial_t d_R^\epsilon + d_R^\epsilon|_{L^2}^2 - |\partial_t d_R^\epsilon|_{L^2}^2 - |d_R^\epsilon|_{L^2}^2) - |\partial_t d_R^\epsilon|_{L^2}^2, \end{aligned}$$

in order to get

$$\begin{aligned} &\frac{1}{2} \frac{d}{dt} \left[|\partial_t d_R^\epsilon + d_R^\epsilon|_{L^2}^2 + \left(\frac{1}{\epsilon} - 1\right) |d_R^\epsilon|_{L^2}^2 - |\partial_t d_R^\epsilon|_{L^2}^2 \right] - |\partial_t d_R^\epsilon|_{L^2}^2 + \frac{1}{\epsilon} |\nabla d_R^\epsilon|_{L^2}^2 \\ &= \langle \mathcal{S}(d_R^\epsilon), d_R^\epsilon \rangle + \langle \mathcal{R}(d_R^\epsilon), d_R^\epsilon \rangle. \end{aligned} \tag{3.15}$$

Using the estimates (3.8) and (3.13) previously derived for bounding the terms $\langle \mathcal{S}(d_R^\epsilon), \partial_t d_R^\epsilon \rangle$ and $\langle \mathcal{R}(d_R^\epsilon), \partial_t d_R^\epsilon \rangle$, we can analogously estimate the terms $\langle \mathcal{S}(d_R^\epsilon), d_R^\epsilon \rangle$ and $\langle \mathcal{R}(d_R^\epsilon), d_R^\epsilon \rangle$ as follows:

$$\langle \mathcal{S}(d_R^\epsilon), d_R^\epsilon \rangle \leq C_{11} \left(1 + |d_R^\epsilon|_{H^2}^2 + \left| \frac{d_R^\epsilon}{\sqrt{\epsilon}} \right|_{L^2} + \left| \frac{\nabla d_R^\epsilon}{\sqrt{\epsilon}} \right|_{L^2} \right) \left| \frac{d_R^\epsilon}{\sqrt{\epsilon}} \right|_{L^2}, \tag{3.16}$$

and

$$\begin{aligned} \langle \mathcal{R}(d_R^\epsilon), d_R^\epsilon \rangle &\leq C_{12} \left[\sqrt{\epsilon} + |d_R^\epsilon|_{H^1} + |\partial_t d_R^\epsilon|_{L^2} + \sqrt{\epsilon} |\partial_t d_R^\epsilon|_{H^1}^2 + \sqrt{\epsilon} |d_R^\epsilon|_{H^2}^2 \right. \\ &\quad \left. + |d_R^\epsilon|_{H^2} (\epsilon |\partial_t d_R^\epsilon|_{H^1}^2 + |\nabla d_R^\epsilon|_{H^1}^2) \right] |d_R^\epsilon|_{L^2}. \end{aligned} \tag{3.17}$$

So plugging (3.16) and (3.17) into (3.15) we obtain:

$$\begin{aligned} &\frac{1}{2} \frac{d}{dt} \left[|\partial_t d_R^\epsilon + d_R^\epsilon|_{L^2}^2 + \left(\frac{1}{\epsilon} - 1\right) |d_R^\epsilon|_{L^2}^2 - |\partial_t d_R^\epsilon|_{L^2}^2 \right] - |\partial_t d_R^\epsilon|_{L^2}^2 + \frac{1}{\epsilon} |\nabla d_R^\epsilon|_{L^2}^2 \\ &\leq C_1 \left\{ \left(1 + |d_R^\epsilon|_{H^2}^2 + \left| \frac{d_R^\epsilon}{\sqrt{\epsilon}} \right|_{L^2} + \left| \frac{\nabla d_R^\epsilon}{\sqrt{\epsilon}} \right|_{L^2} \right) \left| \frac{d_R^\epsilon}{\sqrt{\epsilon}} \right|_{L^2} + \left[\sqrt{\epsilon} + |d_R^\epsilon|_{H^1} + |\partial_t d_R^\epsilon|_{L^2} \right. \right. \\ &\quad \left. \left. + \sqrt{\epsilon} |\partial_t d_R^\epsilon|_{H^1}^2 + \sqrt{\epsilon} |d_R^\epsilon|_{H^2}^2 + |d_R^\epsilon|_{H^2} (\epsilon |\partial_t d_R^\epsilon|_{H^1}^2 + |\nabla d_R^\epsilon|_{H^1}^2) \right] |d_R^\epsilon|_{L^2} \right\}. \end{aligned} \tag{3.18}$$

Multiplying the inequality (3.18) by $\frac{1}{2}$ and then adding it to the inequality (3.14), we get the L^2 -energy estimate:

$$\begin{aligned} &\frac{1}{4} \frac{d}{dt} \left[|\partial_t d_R^\epsilon|_{L^2}^2 + |\partial_t d_R^\epsilon + d_R^\epsilon|_{L^2}^2 + \left(\frac{1}{\epsilon} - 1\right) |d_R^\epsilon|_{L^2}^2 \right. \\ &\quad \left. + \frac{2}{\epsilon} |\nabla d_R^\epsilon|_{L^2}^2 \right] + \left(\frac{1}{\epsilon} - \frac{1}{2}\right) |\partial_t d_R^\epsilon|_{L^2}^2 + \frac{1}{2\epsilon} |\nabla d_R^\epsilon|_{L^2}^2 \\ &\leq C_1 \left\{ \left(1 + |d_R^\epsilon|_{H^2}^2 + \left| \frac{d_R^\epsilon}{\sqrt{\epsilon}} \right|_{L^2} + \left| \frac{\nabla d_R^\epsilon}{\sqrt{\epsilon}} \right|_{L^2} \right) \left(\left| \frac{\partial_t d_R^\epsilon}{\sqrt{\epsilon}} \right|_{L^2} + \left| \frac{d_R^\epsilon}{\sqrt{\epsilon}} \right|_{L^2} \right) \right. \\ &\quad \left. + \left[\sqrt{\epsilon} + |d_R^\epsilon|_{H^1} + |\partial_t d_R^\epsilon|_{L^2} + \sqrt{\epsilon} |\partial_t d_R^\epsilon|_{H^1}^2 + \sqrt{\epsilon} |d_R^\epsilon|_{H^2}^2 \right. \right. \\ &\quad \left. \left. + |d_R^\epsilon|_{H^2} (\epsilon |\partial_t d_R^\epsilon|_{H^1}^2 + |\nabla d_R^\epsilon|_{H^1}^2) \right] \left(|\partial_t d_R^\epsilon|_{L^2} + |d_R^\epsilon|_{L^2} \right) \right\}, \end{aligned} \tag{3.19}$$

where the constant

$$C_1 = C(1 + |D(x)|_{H^3}^3 + |\partial_{tt}d_0|_{L^\infty(0,T;L^2)} + |\partial_t d_0|_{L^\infty(0,T;H^2)}^3 + |\nabla d_0|_{L^\infty(0,T;H^2)}^4) > 0$$

for some computable positive constant C .

Step 2. Higher order estimates.

In order to use the inequality (3.19) we also need a higher order estimate. To obtain this we take $\nabla^k (k=1,2)$ in the Equation (1.15), we multiply it by $\nabla^k \partial_t d_R^\epsilon$, integrate over \mathbb{R}^3 and by parts, thus obtaining the following equality

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} (|\partial_t \nabla^k d_R^\epsilon|_{L^2}^2 + \frac{1}{\epsilon} |\nabla^{k+1} d_R^\epsilon|_{L^2}^2) + \frac{1}{\epsilon} |\partial_t \nabla^k d_R^\epsilon|_{L^2}^2 \\ & = \langle \nabla^k \mathcal{S}(d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle + \langle \nabla^k \mathcal{R}(d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle. \end{aligned} \tag{3.20}$$

(I) *Estimates of the singular terms $\langle \nabla^k \mathcal{S}(d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle$:*

The singular terms can be divided into three parts: $\langle \nabla^k \mathcal{S}_i, \partial_t \nabla^k d_R^\epsilon \rangle (i=1,2,3)$ which we estimate separately.

For the term $\langle \nabla^k \mathcal{S}_1, \partial_t \nabla^k d_R^\epsilon \rangle$, by using the Hölder inequality and the Sobolev embedding theorems, we obtain:

$$\frac{1}{\sqrt{\epsilon}} \langle \nabla^k (\partial_{tt}d_0 + \Delta D(x) \exp(-\frac{t}{\epsilon})), \partial_t \nabla^k d_R^\epsilon \rangle \lesssim (|\partial_{tt}d_0|_{L_t^\infty H_x^2} + |D(x)|_{H^4}) \left| \frac{\partial_t d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2}$$

and

$$\begin{aligned} & \frac{1}{\sqrt{\epsilon}} \langle \nabla^k (|\partial_t d_0 + D(x) \exp(-\frac{t}{\epsilon})|^2 d_0), \partial_t \nabla^k d_R^\epsilon \rangle \\ & \lesssim \frac{1}{\sqrt{\epsilon}} \sum_{\substack{i+j+e=k \\ e \geq 1}} |\langle \nabla^i (\partial_t d_0 + D(x)) \nabla^j (\partial_t d_0 + D(x)) \nabla^e d_0, \partial_t \nabla^k d_R^\epsilon \rangle| \\ & \quad + \frac{1}{\sqrt{\epsilon}} \sum_{i+j=k} |\langle \nabla^i (\partial_t d_0 + D(x)) \nabla^j (\partial_t d_0 + D(x)) d_0, \partial_t \nabla^k d_R^\epsilon \rangle| \\ & \lesssim (1 + |\nabla d_0|_{L_t^\infty H_x^2}) (|\partial_t d_0|_{L_t^\infty H_x^2}^2 + |D(x)|_{H^2}^2) \left| \frac{\partial_t d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2}. \end{aligned}$$

Similarly as for estimating $\frac{1}{\sqrt{\epsilon}} \langle \nabla^k (|\partial_t d_0 + D(x) \exp(-\frac{t}{\epsilon})|^2 d_0), \partial_t \nabla^k d_R^\epsilon \rangle$, we can easily get the following estimates:

$$\begin{aligned} & \frac{1}{\sqrt{\epsilon}} \langle \nabla^k (|\nabla d_0|^2 D(x) \exp(-\frac{t}{\epsilon})), \partial_t \nabla^k d_R^\epsilon \rangle \\ & \lesssim \frac{1}{\sqrt{\epsilon}} \sum_{i+j+e=k} |\langle \nabla^{i+1} d_0 \nabla^{j+1} d_0 \nabla^e D(x), \partial_t \nabla^k d_R^\epsilon \rangle| \\ & \lesssim |\nabla d_0|_{L_t^\infty H_x^2}^2 |D(x)|_{H^3} \left| \frac{\partial_t d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2} \end{aligned}$$

and

$$\begin{aligned} & \frac{1}{\sqrt{\epsilon}} \langle \nabla^k (\nabla d_0 \cdot \nabla D(x) \exp(-\frac{t}{\epsilon}) d_0), \partial_t \nabla^k d_R^\epsilon \rangle \\ & \lesssim \frac{1}{\sqrt{\epsilon}} \sum_{\substack{i+j+e=k \\ e \geq 1}} |\langle \nabla^{i+1} d_0 \cdot \nabla^{j+1} D(x) \nabla^e d_0, \partial_t \nabla^k d_R^\epsilon \rangle| \\ & \quad + \frac{1}{\sqrt{\epsilon}} \sum_{i+j=k} |\langle \nabla^{i+1} d_0 \cdot \nabla^{j+1} D(x) d_0, \partial_t \nabla^k d_R^\epsilon \rangle| \\ & \lesssim (1 + |\nabla d_0|_{L_t^\infty H_x^2}) |\nabla d_0|_{L_t^\infty H_x^3} |D(x)|_{H^3} \left| \frac{\partial_t d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2}. \end{aligned}$$

Thus we have the estimate of $\langle \nabla^k \mathcal{S}_1, \partial_t \nabla^k d_R^\epsilon \rangle$ as follows:

$$\begin{aligned} \langle \nabla^k \mathcal{S}_1, \partial_t \nabla^k d_R^\epsilon \rangle &\lesssim \left\{ (1 + |\nabla d_0|_{L_t^\infty H_x^2}) (|\partial_t d_0|_{L_t^\infty H_x^3}^2 + |\nabla d_0|_{L_t^\infty H_x^3}^2 + |D(x)|_{H^3}^2) \right. \\ &\quad \left. + |\partial_{tt} d_0|_{L_t^\infty H_x^2} + |D(x)|_{H^4} \right\} \left| \frac{\partial_t d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2}. \end{aligned} \tag{3.21}$$

For the term $\langle \nabla^k \mathcal{S}_2, \partial_t \nabla^k d_R^\epsilon \rangle$, we can also use the Hölder inequality and Soblev embedding theorems to get

$$\begin{aligned} &\frac{1}{\sqrt{\epsilon}} \langle \nabla^k ((\nabla d_0 \cdot \nabla d_R^\epsilon) d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle \\ &\lesssim \frac{1}{\sqrt{\epsilon}} \sum_{\substack{i+j+e=k \\ e \geq 1}} |\langle \nabla^{i+1} d_0 \cdot \nabla^{j+1} d_R^\epsilon \nabla^e d_R^\epsilon, \partial_t \nabla^k d_R^\epsilon \rangle| \\ &\quad + \frac{1}{\sqrt{\epsilon}} \sum_{i+j=k} |\langle \nabla^{i+1} d_0 \cdot \nabla^{j+1} d_R^\epsilon d_R^\epsilon, \partial_t \nabla^k d_R^\epsilon \rangle| \\ &\lesssim |\nabla d_0|_{L_t^\infty H_x^4} |d_R^\epsilon|_{H^2} |\nabla d_R^\epsilon|_{H^2} \left| \frac{\partial_t d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2}, \end{aligned}$$

and

$$\begin{aligned} \frac{1}{\sqrt{\epsilon}} \langle \nabla^k (|\nabla d_R^\epsilon|^2 d_0), \partial_t \nabla^k d_R^\epsilon \rangle &\lesssim \frac{1}{\sqrt{\epsilon}} \sum_{\substack{i+j+e=k \\ e \geq 1}} |\langle \nabla^{i+1} d_R^\epsilon \nabla^{j+1} d_R^\epsilon \nabla^e d_0, \partial_t \nabla^k d_R^\epsilon \rangle| \\ &\quad + \frac{1}{\sqrt{\epsilon}} \sum_{i+j=k} |\langle \nabla^{i+1} d_R^\epsilon \nabla^{j+1} d_R^\epsilon d_0, \partial_t \nabla^k d_R^\epsilon \rangle| \\ &\lesssim (1 + |\nabla d_0|_{L_t^\infty H_x^3}) |\nabla d_R^\epsilon|_{H^2}^2 \left| \frac{\partial_t d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2}. \end{aligned}$$

Summarizing, we obtain

$$\langle \nabla^k \mathcal{S}_2, \partial_t \nabla^k d_R^\epsilon \rangle \lesssim (1 + |\nabla d_0|_{L_t^\infty H_x^4}) (|d_R^\epsilon|_{H^2} + |\nabla d_R^\epsilon|_{H^2}) |\nabla d_R^\epsilon|_{H^2} \left| \frac{\partial_t d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2}. \tag{3.22}$$

For the estimate of $\langle \nabla^k \mathcal{S}_3, \partial_t \nabla^k d_R^\epsilon \rangle$, we get the estimate of the first term by using again the Hölder inequality and the Sobolev embedding theorems:

$$\begin{aligned} \frac{1}{\epsilon} \langle \nabla^k (|\nabla d_0|^2 d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle &= \frac{1}{\epsilon} \sum_{i+j+e=k} \langle \nabla^{i+1} d_0 \nabla^{j+1} d_0 \nabla^e d_R^\epsilon, \partial_t \nabla^k d_R^\epsilon \rangle \\ &\lesssim |\nabla d_0|_{L_t^\infty H_x^4}^2 \left| \frac{d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2} \left| \frac{\partial_t d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2}. \end{aligned}$$

Recalling that $|d_0|=1$, one can easily estimate the second term

$$\begin{aligned} &\frac{1}{\epsilon} \langle \nabla^k ((\nabla d_0 \cdot \nabla d_R^\epsilon) d_0), \partial_t \nabla^k d_R^\epsilon \rangle \\ &\lesssim \frac{1}{\epsilon} \sum_{\substack{i+j+e=k \\ e \geq 1}} |\langle \nabla^{i+1} d_0 \cdot \nabla^{j+1} d_R^\epsilon \nabla^e d_0, \partial_t \nabla^k d_R^\epsilon \rangle| \\ &\quad + \frac{1}{\epsilon} \sum_{i+j=k} |\langle \nabla^{i+1} d_0 \cdot \nabla^{j+1} d_R^\epsilon d_0, \partial_t \nabla^k d_R^\epsilon \rangle| \\ &\lesssim (1 + |\nabla d_0|_{L_t^\infty H_x^3}) |\nabla d_0|_{L_t^\infty H_x^4} \left| \frac{\nabla d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2} \left| \frac{\partial_t d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2}. \end{aligned}$$

Thus by the above two estimates we have

$$\langle \nabla^k \mathcal{S}_3, \partial_t \nabla^k d_R^\epsilon \rangle \lesssim (1 + |\nabla d_0|_{L_t^\infty H_x^4}) |\nabla d_0|_{L_t^\infty H_x^4} \left(\left| \frac{d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2} + \left| \frac{\nabla d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2} \right) \left| \frac{\partial_t d_R^\epsilon}{\sqrt{\epsilon}} \right|_{H^2}. \tag{3.23}$$

Then the inequalities (3.21), (3.22) and (3.23) give the following estimate

$$\langle \nabla^k \mathcal{S}(d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle \leq C_{k1} (1 + |d_R^\epsilon|_{H^2}^2 + |\nabla d_R^\epsilon|_{H^2}^2 + |\frac{d_R^\epsilon}{\sqrt{\epsilon}}|_{H^2} + |\frac{\nabla d_R^\epsilon}{\sqrt{\epsilon}}|_{H^2}) |\frac{\partial_t d_R^\epsilon}{\sqrt{\epsilon}}|_{H^2}, \quad (3.24)$$

where the positive constant C_{k1} is

$$C_{k1} = C \left\{ (1 + |\nabla d_0|_{L_t^\infty H_x^4}) (1 + |\partial_t d_0|_{L_t^\infty H_x^3}^2 + |\nabla d_0|_{L_t^\infty H_x^4}^2 + |D(x)|_{H^3}^2) + |\partial_{tt} d_0|_{L_t^\infty H_x^2} + |D(x)|_{H^4} \right\}$$

for some computable positive constant C .

(II) *Estimates of the regular terms* $\langle \nabla^k \mathcal{R}(d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle$:

Finally, we turn to estimating the regular terms $\langle \nabla^k \mathcal{R}(d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle$, which are divided into four parts: $\langle \nabla^k \mathcal{R}_i, \partial_t \nabla^k d_R^\epsilon \rangle (i = 1, 2, 3)$.

For the terms $\langle \nabla^k \mathcal{R}_1, \partial_t \nabla^k d_R^\epsilon \rangle$, by the Hölder inequality and the Sobolev embedding theorems we have:

$$\begin{aligned} & \sqrt{\epsilon} \langle \nabla^k (|\partial_t d_0 + D(x) \exp(-\frac{t}{\epsilon})|^2 D(x) \exp(-\frac{t}{\epsilon})), \partial_t \nabla^k d_R^\epsilon \rangle \\ & \lesssim \sqrt{\epsilon} \sum_{i+j+e=k} |\langle \nabla^i (\partial_t d_0 + D(x)) \nabla^j (\partial_t d_0 + D(x)) \nabla^e D(x), \partial_t \nabla^k d_R^\epsilon \rangle| \\ & \lesssim \sqrt{\epsilon} |D(x)|_{H^3} (|D(x)|_{H^2}^2 + |\partial_t d_0|_{L_t^\infty H_x^2}^2) |\partial_t d_R^\epsilon|_{H^2}, \end{aligned}$$

and

$$\begin{aligned} & \sqrt{\epsilon} \langle \nabla^k (|\nabla D(x)|^2 \exp(-\frac{2t}{\epsilon}) d_0), \partial_t \nabla^k d_R^\epsilon \rangle \\ & \lesssim \sqrt{\epsilon} \sum_{\substack{i+j+e=k \\ e \geq 1}} |\langle \nabla^{i+1} D(x) \nabla^{j+1} D(x) \nabla^e d_0, \partial_t \nabla^k d_0 \rangle| \\ & \quad + \sqrt{\epsilon} \sum_{i+j=k} |\langle \nabla^{i+1} D(x) \nabla^{j+1} D(x) d_0, \partial_t \nabla^k d_R^\epsilon \rangle| \\ & \lesssim \sqrt{\epsilon} (1 + |\nabla d_0|_{L_t^\infty H_x^2}) |D(x)|_{H^3}^2 |\partial_t d_R^\epsilon|_{H^2}. \end{aligned}$$

We can estimate the following two terms in a similar way, hence we get the following inequalities:

$$\begin{aligned} & \epsilon^{\frac{3}{2}} \langle \nabla^k (|\nabla D(x)|^2 \exp(-\frac{3t}{\epsilon}) D(x)), \partial_t \nabla^k d_R^\epsilon \rangle \\ & \lesssim \epsilon^{\frac{3}{2}} \sum_{i+j+e=k} |\langle \nabla^{i+1} D(x) \nabla^{j+1} D(x) \nabla^e D(x), \partial_t \nabla^k d_R^\epsilon \rangle| \\ & \lesssim \epsilon^{\frac{3}{2}} |D(x)|_{H^3}^3 |\partial_t d_R^\epsilon|_{H^2} \end{aligned}$$

and

$$\begin{aligned} & \sqrt{\epsilon} \langle \nabla^k (\nabla d_0 \cdot \nabla D(x)) \exp(-\frac{2t}{\epsilon}) D(x), \partial_t \nabla^k d_R^\epsilon \rangle \\ & \lesssim \sqrt{\epsilon} \sum_{i+j+e=k} |\langle \nabla^{i+1} d_0 \nabla^{j+1} D(x) \nabla^e D(x), \partial_t \nabla^k d_R^\epsilon \rangle| \\ & \lesssim \sqrt{\epsilon} |\nabla d_0|_{L_t^\infty H_x^4} |D(x)|_{H^3}^2 |\partial_t d_R^\epsilon|_{H^2}. \end{aligned}$$

So we have the estimate of $\langle \nabla^k \mathcal{R}_1, \partial_t \nabla^k d_R^\epsilon \rangle$ as follows:

$$\langle \nabla^k \mathcal{R}_1, \partial_t \nabla^k d_R^\epsilon \rangle \lesssim \sqrt{\epsilon} (1 + |\nabla d_0|_{L_t^\infty H_x^4} + |D(x)|_{H^3}) \times (|\partial_t d_0|_{L_t^\infty H_x^2}^2 + |D(x)|_{H^3}^2) |\partial_t d_R^\epsilon|_{H^2}. \quad (3.25)$$

For the terms $\langle \nabla^k \mathcal{R}_2, \partial_t \nabla^k d_R^\epsilon \rangle$, by using yet again the Hölder inequality and the Sobolev embedding theorems, we have

$$\begin{aligned} & \varepsilon \langle \nabla^k (\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})) \partial_t d_R^\epsilon D(x) \exp(-\frac{t}{\varepsilon}), \partial_t \nabla^k d_R^\epsilon \rangle \\ & \lesssim \varepsilon \sum_{\substack{i+j+e=k \\ e \geq 1}} |\langle (\nabla^i \partial_t d_0 + \nabla^i D(x) \exp(-\frac{t}{\varepsilon})) \nabla^j \partial_t d_R^\epsilon \nabla^e D(x) \exp(-\frac{t}{\varepsilon}), \partial_t \nabla^k d_R^\epsilon \rangle| \\ & \quad + \varepsilon \sum_{i+j=k} |\langle (\nabla^i \partial_t d_0 + \nabla^i D(x) \exp(-\frac{t}{\varepsilon})) \nabla^j \partial_t d_R^\epsilon D(x) \exp(-\frac{t}{\varepsilon}), \partial_t \nabla^k d_R^\epsilon \rangle| \\ & \lesssim \varepsilon (|\partial_t d_0|_{L_t^\infty H_x^4} + |D(x)|_{H^4}) |D(x)|_{H^4} |\partial_t d_R^\epsilon|_{H^2}^2, \end{aligned}$$

Similarly as for estimating the term $\varepsilon \langle \nabla^k (\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})) \partial_t d_R^\epsilon D(x) \exp(-\frac{t}{\varepsilon}), \partial_t \nabla^k d_R^\epsilon \rangle$, we can easily obtain the following estimates:

$$\begin{aligned} & \langle \nabla^k (|\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})|^2 d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle \lesssim (|\partial_t d_0|_{L_t^\infty H_x^4}^2 + |D(x)|_{H^4}^2) |d_R^\epsilon|_{H^2} |\partial_t d_R^\epsilon|_{H^2}, \\ & \langle \nabla^k (\nabla d_0 \cdot \nabla d_R^\epsilon D(x) \exp(-\frac{t}{\varepsilon})), \partial_t \nabla^k d_R^\epsilon \rangle \lesssim (|\nabla d_0|_{L_t^\infty H_x^4}^2 + |D(x)|_{H^4}^2) |\nabla d_R^\epsilon|_{H^2} |\partial_t d_R^\epsilon|_{H^2}, \\ & \langle \nabla^k (\nabla d_0 \cdot \nabla D(x) \exp(-\frac{t}{\varepsilon}) d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle \lesssim (|\nabla d_0|_{L_t^\infty H_x^4}^2 + |D(x)|_{H^4}^2) |d_R^\epsilon|_{H^2} |\partial_t d_R^\epsilon|_{H^2}. \end{aligned}$$

Observing the structure of the terms $\varepsilon \langle \nabla^k (|\nabla D(x)|^2 \exp(-\frac{2t}{\varepsilon}) d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle$ and $\varepsilon \langle \nabla^k (\nabla D(x) \cdot \nabla d_R^\epsilon \exp(-\frac{2t}{\varepsilon}) D(x)), \partial_t \nabla^k d_R^\epsilon \rangle$, one can similarly estimate the following terms:

$$\begin{aligned} & \varepsilon \langle \nabla^k (|\nabla D(x)|^2 \exp(-\frac{2t}{\varepsilon}) d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle \lesssim \varepsilon |D(x)|_{H^5}^2 |d_R^\epsilon|_{H^2} |\partial_t d_R^\epsilon|_{H^2}, \\ & \varepsilon \langle \nabla^k (\nabla D(x) \cdot \nabla d_R^\epsilon \exp(-\frac{2t}{\varepsilon}) D(x)), \partial_t \nabla^k d_R^\epsilon \rangle \lesssim \varepsilon |D(x)|_{H^5}^2 |\nabla d_R^\epsilon|_{H^2} |\partial_t d_R^\epsilon|_{H^2}. \end{aligned}$$

Furthermore we get:

$$\begin{aligned} & \langle \nabla^k [(\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})) \cdot \partial_t d_R^\epsilon d_0], \partial_t \nabla^k d_R^\epsilon \rangle \\ & \lesssim \sum_{\substack{i+j+e=k \\ e \geq 1}} |\langle \nabla^i (\partial_t d_0 + D(x)) \nabla^j \partial_t d_R^\epsilon \nabla^e d_0, \partial_t \nabla^k d_R^\epsilon \rangle| \\ & \quad + \sum_{i+j=k} |\langle \nabla^i (\partial_t d_0 + D(x)) \nabla^j \partial_t d_R^\epsilon d_0, \partial_t \nabla^k d_R^\epsilon \rangle| \\ & \lesssim (1 + |\nabla d_0|_{L_t^\infty H_x^3}) (|\partial_t d_0|_{L_t^\infty H_x^4} + |D(x)|_{H^4}) |\partial_t d_R^\epsilon|_{H^2}^2. \end{aligned}$$

Similarly as for estimating $\langle \nabla^k [(\partial_t d_0 + D(x) \exp(-\frac{t}{\varepsilon})) \cdot \partial_t d_R^\epsilon d_0], \partial_t \nabla^k d_R^\epsilon \rangle$, one can also get:

$$\langle \nabla^k (\nabla D(x) \cdot \nabla d_R^\epsilon \exp(-\frac{t}{\varepsilon}) d_0), \partial_t \nabla^k d_R^\epsilon \rangle \lesssim (1 + |\nabla d_0|_{L_t^\infty H_x^3}) |D(x)|_{H^5} |\nabla d_R^\epsilon|_{H^2} |\partial_t d_R^\epsilon|_{H^2}.$$

Thus we have the following estimate of $\langle \nabla^k \mathcal{R}_2, \partial_t \nabla^k d_R^\epsilon \rangle$:

$$\begin{aligned} \langle \nabla^k \mathcal{R}_2, \partial_t \nabla^k d_R^\epsilon \rangle & \lesssim (|\partial_t d_0|_{L_t^\infty H_x^4} + |D(x)|_{H^5} + |\partial_t d_0|_{L_t^\infty H_x^4}^2 + |D(x)|_{H^5}^2 + |\nabla d_0|_{L_t^\infty H_x^3}^2) \\ & \quad \times (|d_R^\epsilon|_{H^2} + |\nabla d_R^\epsilon|_{H^2} + |\partial_t d_R^\epsilon|_{H^2}) |\partial_t d_R^\epsilon|_{H^2}. \end{aligned} \tag{3.26}$$

For the terms $\langle \nabla^k \mathcal{R}_3, \partial_t \nabla^k d_R^\epsilon \rangle$, one can use an estimate similar to the one for the term $\varepsilon \langle \nabla^k (\partial_t d_0 \cdot D(x) \exp(-\frac{t}{\varepsilon}) \partial_t d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle$ to get:

$$\begin{aligned} & \varepsilon^{\frac{3}{2}} \langle \nabla^k (D(x) \exp(-\frac{t}{\varepsilon}) |\partial_t d_R^\epsilon|^2), \partial_t \nabla^k d_R^\epsilon \rangle \lesssim \varepsilon^{\frac{3}{2}} |D(x)|_{H^4} |\partial_t d_R^\epsilon|_{H^2}^3, \\ & \sqrt{\varepsilon} \langle \nabla^k (\nabla D(x) \cdot \nabla d_R^\epsilon \exp(-\frac{t}{\varepsilon}) d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle \lesssim \sqrt{\varepsilon} |D(x)|_{H^5} (|d_R^\epsilon|_{H^2}^2 + |\nabla d_R^\epsilon|_{H^2}^2) |\partial_t d_R^\epsilon|_{H^2}, \\ & \sqrt{\varepsilon} \langle \nabla^k (|\nabla d_R^\epsilon|^2 D(x) \exp(-\frac{t}{\varepsilon})), \partial_t \nabla^k d_R^\epsilon \rangle \lesssim \sqrt{\varepsilon} |D(x)|_{H^4} |\nabla d_R^\epsilon|_{H^2}^2 |\partial_t d_R^\epsilon|_{H^2}. \end{aligned}$$

Reasoning analogously as in estimating $\sqrt{\varepsilon}\langle \nabla^k [(\partial_t d_0 + D(x)\exp(-\frac{t}{\varepsilon})) \cdot \partial_t d_R^\varepsilon d_R^\varepsilon], \partial_t \nabla^k d_R^\varepsilon \rangle$ we have

$$\begin{aligned} & \sqrt{\varepsilon}\langle \nabla^k [(\partial_t d_0 + D(x)\exp(-\frac{t}{\varepsilon})) \cdot \partial_t d_R^\varepsilon d_R^\varepsilon], \partial_t \nabla^k d_R^\varepsilon \rangle \\ & \lesssim \sqrt{\varepsilon} \sum_{\substack{i+j+e=k \\ e \geq 1}} |\langle \nabla^i (\partial_t d_0 + D(x)) \cdot \nabla^j \partial_t d_R^\varepsilon \nabla^e d_R^\varepsilon, \partial_t \nabla^k d_R^\varepsilon \rangle| \\ & \quad + \sqrt{\varepsilon} \sum_{i+j=k} |\langle \nabla^i (\partial_t d_0 + D(x)) \cdot \nabla^j \partial_t d_R^\varepsilon d_R^\varepsilon, \partial_t \nabla^k d_R^\varepsilon \rangle| \\ & \lesssim \sqrt{\varepsilon} (|\partial_t d_0|_{L_t^\infty H_x^4} + |D(x)|_{H^4}) (|d_R^\varepsilon|_{H^2} + |\nabla d_R^\varepsilon|_{H^2}) |\partial_t d_R^\varepsilon|_{H^2}^2 \end{aligned}$$

and furthermore, using that $|d_0| = 1$, it is easy to obtain

$$\sqrt{\varepsilon}\langle \nabla^k (|\partial_t d_R^\varepsilon|^2 d_0), \partial_t \nabla^k d_R^\varepsilon \rangle \lesssim \sqrt{\varepsilon} (1 + |\nabla d_0|_{L_t^\infty H_x^3}) |\partial_t d_R^\varepsilon|_{H^2}^3.$$

Summarizing, we get the estimate of $\langle \nabla^k \mathcal{R}_3, \partial_t \nabla^k d_R^\varepsilon \rangle$ as follows:

$$\begin{aligned} \langle \nabla^k \mathcal{R}_3, \partial_t \nabla^k d_R^\varepsilon \rangle & \lesssim \sqrt{\varepsilon} (1 + |\partial_t d_0|_{L_t^\infty H_x^4} + |\nabla d_0|_{L_t^\infty H_x^3} + |D(x)|_{H^5}) \\ & \quad \times (1 + |\partial_t d_R^\varepsilon|_{H^2}^2 + |d_R^\varepsilon|_{H^2}^2 + |\nabla d_R^\varepsilon|_{H^2}^2) |\partial_t d_R^\varepsilon|_{H^2}. \end{aligned} \tag{3.27}$$

For the terms $\langle \nabla^k \mathcal{R}_4, \partial_t \nabla^k d_R^\varepsilon \rangle$, we get, by the Hölder inequality and Sobolev embedding theorems:

$$\begin{aligned} \varepsilon \langle \nabla^k (|\partial_t d_R^\varepsilon|^2 d_R^\varepsilon), \partial_t \nabla^k d_R^\varepsilon \rangle & = \varepsilon \sum_{\substack{i+j+e=k \\ e \geq 1}} \langle \nabla^i \partial_t d_R^\varepsilon \nabla^j \partial_t d_R^\varepsilon \nabla^e d_R^\varepsilon, \partial_t \nabla^k d_R^\varepsilon \rangle \\ & \quad + \varepsilon \sum_{i+j=k} \langle \nabla^i \partial_t d_R^\varepsilon \nabla^j \partial_t d_R^\varepsilon d_R^\varepsilon, \partial_t \nabla^k d_R^\varepsilon \rangle \\ & \lesssim \varepsilon (|d_R^\varepsilon|_{H^2} + |\nabla d_R^\varepsilon|_{H^2}) |\partial_t d_R^\varepsilon|_{H^2}^3 \end{aligned}$$

and then, similarly:

$$\langle \nabla^k (|\nabla d_R^\varepsilon|^2 d_R^\varepsilon), \partial_t \nabla^k d_R^\varepsilon \rangle \lesssim (|d_R^\varepsilon|_{H^2} + |\nabla d_R^\varepsilon|_{H^2}) |\nabla d_R^\varepsilon|_{H^2}^2 |\partial_t d_R^\varepsilon|_{H^2}.$$

So we obtain the estimate of $\langle \nabla^k \mathcal{R}_4, \partial_t \nabla^k d_R^\varepsilon \rangle$ as follows:

$$\langle \nabla^k \mathcal{R}_4, \partial_t \nabla^k d_R^\varepsilon \rangle \lesssim (|d_R^\varepsilon|_{H^2} + |\nabla d_R^\varepsilon|_{H^2}) (|\nabla d_R^\varepsilon|_{H^2}^2 + \varepsilon |\partial_t d_R^\varepsilon|_{H^2}^2) |\partial_t d_R^\varepsilon|_{H^2}. \tag{3.28}$$

Then the inequalities (3.25), (3.26), (3.27) and (3.28) give the estimate of the regular terms $\langle \nabla^k \mathcal{R}(d_R^\varepsilon), \partial_t \nabla^k d_R^\varepsilon \rangle$ as follows:

$$\begin{aligned} \langle \nabla^k \mathcal{R}(d_R^\varepsilon), \partial_t \nabla^k d_R^\varepsilon \rangle & \leq C_{k2} [(|d_R^\varepsilon|_{H^2} + |\nabla d_R^\varepsilon|_{H^2}) (1 + |\nabla d_R^\varepsilon|_{H^2}^2 + \varepsilon |\partial_t d_R^\varepsilon|_{H^2}^2) \\ & \quad + \sqrt{\varepsilon} (1 + |\partial_t d_R^\varepsilon|_{H^2}^2 + |d_R^\varepsilon|_{H^2}^2 + |\nabla d_R^\varepsilon|_{H^2}^2) + |\partial_t d_R^\varepsilon|_{H^2}] |\partial_t d_R^\varepsilon|_{H^2}, \end{aligned} \tag{3.29}$$

where the constant C_{k2} is

$$\begin{aligned} C_{k2} & = C (1 + |\partial_t d_0|_{L_t^\infty H_x^4} + |\nabla d_0|_{L_t^\infty H_x^4} + |D(x)|_{H^5}) \\ & \quad \times (1 + |\partial_t d_0|_{L_t^\infty H_x^4}^2 + |D(x)|_{H^5}^2 + |\nabla d_0|_{L_t^\infty H_x^4}^2) > 0, \end{aligned}$$

and C is a computable positive constant.

Therefore, by substituting the inequalities (3.24) and (3.29) into (3.20) one has

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} (|\partial_t \nabla^k d_R^\epsilon|_{L^2}^2 + \frac{1}{\epsilon} |\nabla^{k+1} d_R^\epsilon|_{L^2}^2) + \frac{1}{\epsilon} |\partial_t \nabla^k d_R^\epsilon|_{L^2}^2 \\ & \leq C_k \left\{ (1 + |d_R^\epsilon|_{H^2}^2 + |\nabla d_R^\epsilon|_{H^2}^2 + |\frac{d_R^\epsilon}{\sqrt{\epsilon}}|_{H^2} + |\frac{\nabla d_R^\epsilon}{\sqrt{\epsilon}}|_{H^2}) |\frac{\partial_t d_R^\epsilon}{\sqrt{\epsilon}}|_{H^2} \right. \\ & \quad + [(|d_R^\epsilon|_{H^2} + |\nabla d_R^\epsilon|_{H^2}) (1 + |\nabla d_R^\epsilon|_{H^2}^2 + \epsilon |\partial_t d_R^\epsilon|_{H^2}^2) \\ & \quad \left. + \sqrt{\epsilon} (1 + |\partial_t d_R^\epsilon|_{H^2}^2 + |d_R^\epsilon|_{H^2}^2 + |\nabla d_R^\epsilon|_{H^2}^2) + |\partial_t d_R^\epsilon|_{H^2} \right] |\partial_t d_R^\epsilon|_{H^2} \Big\}, \end{aligned} \tag{3.30}$$

where the positive constant C_k is

$$C_k = C(1 + |\partial_{tt} d_0|_{L^\infty(0,T;H^2)} + |\partial_t d_0|_{L^\infty(0,T;H^4)}^3 + |\nabla d_0|_{L^\infty(0,T;H^4)}^3 + |D(x)|_{H^5}^3) > 0$$

and $C > 0$ is a computable constant.

(III) For the estimate of $|\nabla^k d_R^\epsilon|_{L^2}$ ($k = 1, 2$):

Applying ∇^k ($k = 1, 2$) to the remainder Equation (1.15), multiplying by $\nabla^k d_R^\epsilon$, integrating over \mathbb{R}^3 and by parts, we have

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} [|\nabla^k \partial_t d_R^\epsilon + \nabla^k d_R^\epsilon|_{L^2}^2 + (\frac{1}{\epsilon} - 1) |\nabla^k d_R^\epsilon|_{L^2}^2 - |\partial_t \nabla^k d_R^\epsilon|_{L^2}^2] - |\partial_t \nabla^k d_R^\epsilon|_{L^2}^2 + \frac{1}{\epsilon} |\nabla^{k+1} d_R^\epsilon|_{L^2}^2 \\ & = \langle \nabla^k \mathcal{S}(d_R^\epsilon), \nabla^k d_R^\epsilon \rangle + \langle \nabla^k \mathcal{R}(d_R^\epsilon), \nabla^k d_R^\epsilon \rangle. \end{aligned} \tag{3.31}$$

Similarly as in the estimates of the terms $\langle \nabla^k \mathcal{S}(d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle$ and $\langle \nabla^k \mathcal{R}(d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle$ in the inequalities (3.24) and (3.29), respectively, we can analogously estimate $\langle \nabla^k \mathcal{S}(d_R^\epsilon), \nabla^k d_R^\epsilon \rangle$ and $\langle \nabla^k \mathcal{R}(d_R^\epsilon), \nabla^k d_R^\epsilon \rangle$ as follows:

$$\langle \nabla^k \mathcal{S}(d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle \leq C_{k1} (1 + |d_R^\epsilon|_{H^2}^2 + |\nabla d_R^\epsilon|_{H^2}^2 + |\frac{d_R^\epsilon}{\sqrt{\epsilon}}|_{H^2} + |\frac{\nabla d_R^\epsilon}{\sqrt{\epsilon}}|_{H^2}) |\frac{d_R^\epsilon}{\sqrt{\epsilon}}|_{H^2}, \tag{3.32}$$

and

$$\begin{aligned} \langle \nabla^k \mathcal{R}(d_R^\epsilon), \partial_t \nabla^k d_R^\epsilon \rangle & \leq C_{k2} [(|d_R^\epsilon|_{H^2} + |\nabla d_R^\epsilon|_{H^2}) (1 + |\nabla d_R^\epsilon|_{H^2}^2 + \epsilon |\partial_t d_R^\epsilon|_{H^2}^2) \\ & \quad + \sqrt{\epsilon} (1 + |\partial_t d_R^\epsilon|_{H^2}^2 + |d_R^\epsilon|_{H^2}^2 + |\nabla d_R^\epsilon|_{H^2}^2) + |\partial_t d_R^\epsilon|_{H^2}] |d_R^\epsilon|_{H^2}. \end{aligned} \tag{3.33}$$

By plugging the inequalities (3.32) and (3.33) into the equality (3.31), we get the following estimate:

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} [|\nabla^k \partial_t d_R^\epsilon + \nabla^k d_R^\epsilon|_{L^2}^2 + (\frac{1}{\epsilon} - 1) |\nabla^k d_R^\epsilon|_{L^2}^2 \\ & \quad - |\partial_t \nabla^k d_R^\epsilon|_{L^2}^2] - |\partial_t \nabla^k d_R^\epsilon|_{L^2}^2 + \frac{1}{\epsilon} |\nabla^{k+1} d_R^\epsilon|_{L^2}^2 \\ & \leq C_k \left\{ (1 + |d_R^\epsilon|_{H^2}^2 + |\nabla d_R^\epsilon|_{H^2}^2 + |\frac{d_R^\epsilon}{\sqrt{\epsilon}}|_{H^2} + |\frac{\nabla d_R^\epsilon}{\sqrt{\epsilon}}|_{H^2}) |\frac{d_R^\epsilon}{\sqrt{\epsilon}}|_{H^2} \right. \\ & \quad + [(|d_R^\epsilon|_{H^2} + |\nabla d_R^\epsilon|_{H^2}) (1 + |\nabla d_R^\epsilon|_{H^2}^2 + \epsilon |\partial_t d_R^\epsilon|_{H^2}^2) \\ & \quad \left. + \sqrt{\epsilon} (1 + |\partial_t d_R^\epsilon|_{H^2}^2 + |d_R^\epsilon|_{H^2}^2 + |\nabla d_R^\epsilon|_{H^2}^2) + |\partial_t d_R^\epsilon|_{H^2} \right] |d_R^\epsilon|_{H^2} \Big\}. \end{aligned} \tag{3.34}$$

Multiplying the inequality (3.34) by $\frac{1}{2}$ and adding it to the inequality (3.30), we obtain the higher order estimate:

$$\begin{aligned} & \frac{1}{4} \frac{d}{dt} [|\partial_t \nabla^k d_R^\epsilon|_{L^2}^2 + (\frac{1}{\epsilon} - 1) |\nabla^k d_R^\epsilon|_{L^2}^2 + \frac{2}{\epsilon} |\nabla^{k+1} d_R^\epsilon|_{L^2}^2 + |\nabla^k \partial_t d_R^\epsilon + \nabla^k d_R^\epsilon|_{L^2}^2] \\ & \quad + (\frac{1}{\epsilon} - \frac{1}{2}) |\partial_t \nabla^k d_R^\epsilon|_{L^2}^2 + \frac{1}{2\epsilon} |\nabla^{k+1} d_R^\epsilon|_{L^2}^2 \end{aligned}$$

$$\begin{aligned}
&\leq \frac{3}{2}C_k \left\{ (1 + |\mathrm{d}_R^\epsilon|^2_{H^2} + |\nabla \mathrm{d}_R^\epsilon|^2_{H^2} + |\frac{\mathrm{d}_R^\epsilon}{\sqrt{\epsilon}}|_{H^2} + |\frac{\nabla \mathrm{d}_R^\epsilon}{\sqrt{\epsilon}}|_{H^2}) (|\frac{\mathrm{d}_R^\epsilon}{\sqrt{\epsilon}}|_{H^2} + |\frac{\partial_t \mathrm{d}_R^\epsilon}{\sqrt{\epsilon}}|_{H^2}) \right. \\
&\quad + [(|\mathrm{d}_R^\epsilon|_{H^2} + |\nabla \mathrm{d}_R^\epsilon|_{H^2}) (1 + |\nabla \mathrm{d}_R^\epsilon|^2_{H^2} + \varepsilon |\partial_t \mathrm{d}_R^\epsilon|^2_{H^2}) + |\partial_t \mathrm{d}_R^\epsilon|_{H^2} \\
&\quad \left. + \sqrt{\varepsilon} (1 + |\partial_t \mathrm{d}_R^\epsilon|^2_{H^2} + |\mathrm{d}_R^\epsilon|^2_{H^2} + |\nabla \mathrm{d}_R^\epsilon|^2_{H^2})] (|\mathrm{d}_R^\epsilon|_{H^2} + |\partial_t \mathrm{d}_R^\epsilon|_{H^2}) \right\}. \quad (3.35)
\end{aligned}$$

Therefore, combining the L^2 -estimate (3.19) and the k th-order estimate (3.35) for $k = 1, 2$, we obtain:

$$\begin{aligned}
&\frac{1}{4} \frac{\mathrm{d}}{\mathrm{d}t} \left[|\partial_t \mathrm{d}_R^\epsilon|^2_{H^2} + (\frac{1}{\varepsilon} - 1) |\mathrm{d}_R^\epsilon|^2_{H^2} + \frac{2}{\varepsilon} |\nabla \mathrm{d}_R^\epsilon|^2_{H^2} + |\partial_t \mathrm{d}_R^\epsilon + \mathrm{d}_R^\epsilon|^2_{H^2} \right] \\
&\quad + (\frac{1}{\varepsilon} - \frac{1}{2}) |\partial_t \mathrm{d}_R^\epsilon|^2_{H^2} + \frac{1}{2\varepsilon} |\nabla \mathrm{d}_R^\epsilon|^2_{H^2} \\
&\leq \tilde{C}_k \left\{ (1 + |\mathrm{d}_R^\epsilon|^2_{H^2} + |\nabla \mathrm{d}_R^\epsilon|^2_{H^2} + |\frac{\mathrm{d}_R^\epsilon}{\sqrt{\varepsilon}}|_{H^2} + |\frac{\nabla \mathrm{d}_R^\epsilon}{\sqrt{\varepsilon}}|_{H^2}) (|\frac{\mathrm{d}_R^\epsilon}{\sqrt{\varepsilon}}|_{H^2} + |\frac{\partial_t \mathrm{d}_R^\epsilon}{\sqrt{\varepsilon}}|_{H^2}) \right. \\
&\quad + [(|\mathrm{d}_R^\epsilon|_{H^2} + |\nabla \mathrm{d}_R^\epsilon|_{H^2}) (1 + |\nabla \mathrm{d}_R^\epsilon|^2_{H^2} + \varepsilon |\partial_t \mathrm{d}_R^\epsilon|^2_{H^2}) + |\partial_t \mathrm{d}_R^\epsilon|_{H^2} \\
&\quad \left. + \sqrt{\varepsilon} (1 + |\partial_t \mathrm{d}_R^\epsilon|^2_{H^2} + |\mathrm{d}_R^\epsilon|^2_{H^2} + |\nabla \mathrm{d}_R^\epsilon|^2_{H^2})] (|\mathrm{d}_R^\epsilon|_{H^2} + |\partial_t \mathrm{d}_R^\epsilon|_{H^2}) \right\}, \quad (3.36)
\end{aligned}$$

where the positive constant \tilde{C}_k is

$$\tilde{C}_k = C(1 + |\partial_{tt} \mathrm{d}_0|_{L^\infty(0, T; H^2)} + |\partial_t \mathrm{d}_0|_{L^\infty(0, T; H^4)}^3 + |\nabla \mathrm{d}_0|_{L^\infty(0, T; H^4)}^3 + |\mathrm{D}(x)|_{H^5}^3) > 0$$

and $C > 0$ is a computable constant. Then, by the definition of the energy functionals $E_\varepsilon(t)$ and $F_\varepsilon(t)$, and the condition $0 < \varepsilon < \frac{1}{2}$, the H^2 -estimate (3.36) implies that

$$\frac{\mathrm{d}}{\mathrm{d}t} E_\varepsilon(t) + 4F_\varepsilon(t) \leq C' \left\{ E_\varepsilon^{\frac{1}{2}}(t) + E_\varepsilon(t) + \varepsilon^{\frac{1}{2}} E_\varepsilon^{\frac{3}{2}}(t) + \varepsilon^{\frac{3}{2}} E_\varepsilon^2(t) + [1 + E_\varepsilon^{\frac{1}{2}}(t) + \varepsilon E_\varepsilon(t)] F_\varepsilon^{\frac{1}{2}}(t) \right\},$$

where $C' > 0$, which immediately implies the claimed inequality (3.1) by using Young's inequality. Consequently, the proof of Lemma 3.1 is completed. \square

4. The Proof of Theorem 1.1

In this section we will provide the proof of Theorem 1.1, by using the uniform energy bounds (3.1) in Section 3. Before doing this, we note that for any fixed inertia constant $\varepsilon > 0$ the well-posedness of the remainder system (1.15)-(1.16) can be stated as follows:

PROPOSITION 4.1. *Given $\mathrm{d}^{in} : \mathbb{R}^3 \rightarrow \mathbb{S}^2$ and $\tilde{\mathrm{d}}^{in} : \mathbb{R}^3 \rightarrow \mathbb{R}^3$ satisfying $\nabla \mathrm{d}^{in} \in H^6$, $\tilde{\mathrm{d}}^{in} \in H^4$ with $\mathrm{d}^{in} \cdot \tilde{\mathrm{d}}^{in} \equiv 0$, we define $\mathrm{D}(x) \equiv \tilde{\mathrm{d}}^{in}(x) - \Delta \mathrm{d}^{in}(x) - |\nabla \mathrm{d}^{in}(x)|^2 \mathrm{d}^{in}(x)$ and denote $M = |\mathrm{D}|_{H^2}^2 + 2|\nabla \mathrm{D}|_{H^2}^2 < \infty$.*

Then, for any fixed $\varepsilon \in (0, \frac{1}{2})$, there exists a time $T^\varepsilon = \min\{T, \frac{1}{C} \ln(\frac{1+\varepsilon M}{\varepsilon(1+M)})\} > 0$, where $T, C > 0$ are provided in Proposition 1.1, Lemma 3.1, respectively, such that the remainder Equation (1.15) with the initial conditions (1.16) admits a unique solution $\mathrm{d}_R^\varepsilon \in C([0, T^\varepsilon]; H^3)$ and $\partial_t \mathrm{d}_R^\varepsilon \in C([0, T^\varepsilon]; H^2)$. Moreover, the solution d_R^ε satisfies the inequality

$$|\partial_t \mathrm{d}_R^\varepsilon(t)|_{H^2}^2 + \frac{1}{\varepsilon} |\mathrm{d}_R^\varepsilon(t)|_{H^3}^2 \leq \frac{2M e^{Ct}}{1 + \varepsilon M - \varepsilon(1+M)e^{Ct}} \quad (4.1)$$

for all $t \in [0, T^\varepsilon]$.

Proof. We employ a mollifier argument to prove this proposition. For any fixed $\varepsilon > 0$ we can directly construct a system approximating (1.15)-(1.16) as follows:

$$\begin{cases} \varepsilon \partial_t \mathrm{w}_\eta^\varepsilon = -\mathcal{J}_\eta \mathrm{w}_\eta^\varepsilon + \mathcal{J}_\eta \Delta \mathrm{d}_{R, \eta}^\varepsilon + \varepsilon \mathcal{J}_\eta \mathcal{S}(\mathcal{J}_\eta \mathrm{d}_{R, \eta}^\varepsilon) + \varepsilon \mathcal{J}_\eta \mathcal{R}(\mathcal{J}_\eta \mathrm{d}_{R, \eta}^\varepsilon), \\ \partial_t \mathrm{d}_{R, \eta}^\varepsilon = \mathrm{w}_\eta^\varepsilon, \\ \mathrm{d}_{R, \eta}^\varepsilon|_{t=0} = \sqrt{\varepsilon} \mathcal{J}_\eta \mathrm{D}(x), \quad \mathrm{w}_\eta^\varepsilon|_{t=0} = 0, \end{cases} \quad (4.2)$$

where the mollifier operator \mathcal{J}_η is defined as

$$\mathcal{J}_\eta f = \mathcal{F}^{-1}(\mathbf{1}_{|\xi| \leq \frac{1}{\eta}} \mathcal{F}(f)(\xi)),$$

where the symbol \mathcal{F} denotes the standard Fourier transform operator and \mathcal{F}^{-1} is the inverse Fourier transform operator. By ODE theory in Hilbert spaces one can prove the existence and uniqueness of the approximate system (4.2) on the maximal time interval $[0, T_\eta^\varepsilon)$. Then by the fact $\mathcal{J}_\eta^2 = \mathcal{J}_\eta$ and the uniqueness of (4.2) we know that $\mathcal{J}_\eta d_{R,\eta}^\varepsilon = d_{R,\eta}^\varepsilon$ and $\mathcal{J}_\eta w_\eta^\varepsilon = w_\eta^\varepsilon$. Thus by the analogous energy estimate shown in Lemma 3.1 applied to the approximate system (4.2), one can obtain the following energy inequality for $d_{R,\eta}^\varepsilon$ and w_η^ε

$$\frac{d}{dt} E_{\varepsilon,\eta}(t) + 3F_{\varepsilon,\eta}(t) \leq C[1 + E_{\varepsilon,\eta}(t)] [1 + \varepsilon E_{\varepsilon,\eta}(t)] \tag{4.3}$$

for all $t \in [0, T_\eta^\varepsilon)$, where the positive constant $C > 0$ is independent of ε and η , and the energy functionals $E_{\varepsilon,\eta}(t)$, $F_{\varepsilon,\eta}(t)$ are of the same forms as $E_\varepsilon(t)$, $F_\varepsilon(t)$ defined in Section 3 (replacing d_R^ε by $d_{R,\eta}^\varepsilon$), respectively.

Since $d_{R,\eta}^\varepsilon$ satisfies the initial conditions $\partial_t d_{R,\eta}^\varepsilon(0, x) = 0$ and $d_{R,\eta}^\varepsilon(0, x) = \sqrt{\varepsilon} \mathcal{J}_\eta D(x)$, we know that for $\varepsilon \in (0, \frac{1}{2})$

$$\begin{aligned} E_{\varepsilon,\eta}(0) &= |\partial_t d_{R,\eta}^\varepsilon(0, \cdot)|_{H^2}^2 + (\frac{1}{\varepsilon} - 1) |d_{R,\eta}^\varepsilon(0, \cdot)|_{H^2}^2 \\ &\quad + \frac{2}{\varepsilon} |\nabla d_{R,\eta}^\varepsilon(0, \cdot)|_{H^2}^2 + |\partial_t d_{R,\eta}^\varepsilon(0, \cdot) + d_{R,\eta}^\varepsilon(0, \cdot)|_{H^2}^2 \\ &= (\frac{1}{\varepsilon} - 1) |\sqrt{\varepsilon} \mathcal{J}_\eta D|_{H^2}^2 + \frac{2}{\varepsilon} |\sqrt{\varepsilon} \nabla \mathcal{J}_\eta D|_{H^2}^2 + |\sqrt{\varepsilon} \mathcal{J}_\eta D|_{H^2}^2 \\ &\leq (1 - \varepsilon) |D|_{H^2}^2 + 2 |\nabla D|_{H^2}^2 + \varepsilon |D|_{H^2}^2 \\ &= |D|_{H^2}^2 + 2 |\nabla D|_{H^2}^2 = M < \infty. \end{aligned} \tag{4.4}$$

Then, one can solve the ODE inequality (4.3) with the initial condition (4.4), obtaining that

$$\frac{1 + E_{\varepsilon,\eta}(t)}{1 + \varepsilon E_{\varepsilon,\eta}(t)} \leq \frac{1 + E_{\varepsilon,\eta}(0)}{1 + \varepsilon E_{\varepsilon,\eta}(0)} e^{C(1-\varepsilon)t} \leq \frac{1 + M}{1 + \varepsilon M} e^{Ct}$$

holds for all $t \in [0, T_\eta^\varepsilon)$. Consequently, for all $t \in [0, \min\{T_\eta^\varepsilon, T, \frac{1}{C} \ln(\frac{1+\varepsilon M}{\varepsilon(1+M)})\})$ we know that

$$E_{\varepsilon,\eta}(t) \leq \frac{(1 + M)e^{Ct}}{1 + \varepsilon M - \varepsilon(1 + M)e^{Ct}}. \tag{4.5}$$

Notice that the continuity of $E_{\varepsilon,\eta}(t)$ and the maximality of $T_\eta^\varepsilon > 0$ imply that

$$T_\eta^\varepsilon \geq \frac{1}{C} \ln\left(\frac{1+\varepsilon M}{\varepsilon(1+M)}\right) > 0.$$

Hence the inequality (4.5) holds for all $t \in [0, T^\varepsilon)$ uniformly in $\eta > 0$, where

$$T^\varepsilon = \min\{T, \frac{1}{C} \ln\left(\frac{1+\varepsilon M}{\varepsilon(1+M)}\right)\} > 0.$$

Finally, we can finish the proof of this proposition by standard compactness methods and taking the limit as $\eta \rightarrow 0$. The uniqueness issue can be reduced to the uniqueness of the damped wave map system (1.8) which can be obtained by methods analogous to

those in the book of Shatah and Struwe [16]. For convenience, we omit the details of the proof. \square

The Proof of Theorem 1.1. Now, based on the energy estimate (4.1) in Proposition 4.1, we verify Theorem 1.1. We observe that the function

$$f(\varepsilon) := \frac{1}{C} \ln \left(\frac{1 + \varepsilon M}{\varepsilon(1 + M)} \right)$$

is strictly decreasing in $\varepsilon \in (0, \frac{1}{2})$ and $\lim_{\varepsilon \searrow 0} f(\varepsilon) = +\infty$. Consequently, we can choose

$$\varepsilon_0 = \min \left\{ \frac{1}{2}, \frac{1}{(1+M)e^{CT-M}} \right\} \in (0, \frac{1}{2})$$

such that for any $\varepsilon \in (0, \varepsilon_0)$

$$f(\varepsilon) = \frac{1}{C} \ln \left(\frac{1 + \varepsilon M}{\varepsilon(1 + M)} \right) > T.$$

As a result, for the number T^ε determined in Proposition 4.1, we have that $T^\varepsilon \equiv T$ for all $\varepsilon \in (0, \varepsilon_0)$. Then, the inequality (4.1) in Proposition 4.1 implies that

$$|\partial_t d_R^\varepsilon|_{L^\infty(0,T;H^2)}^2 + \frac{1}{\varepsilon} |d_R^\varepsilon|_{L^\infty(0,T;H^3)}^2 \leq \frac{2Me^{CT}}{1+\varepsilon_0 M - \varepsilon_0(1+M)e^{CT}} := C_0 < \infty,$$

and the proof of Theorem 1.1 is completed.

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