GAUSSIAN PROCESSES ASSOCIATED TO INFINITE BEAD-SPRING NETWORKS II: BEADS WITH MASS AND THE VANISHING MASS LIMIT[∗]

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Abstract. We construct families of Gaussian processes $x_{\varepsilon}(t,n)$, $t \in [0,\infty)$, $n \in \mathbb{Z}$, modeling a class of infinite networks of stochastically fluctuating, interacting beads, of small mass, proportional to ε . We examine covariances $\mathbb{E}(x_{\varepsilon}(t_1,n_1)x_{\varepsilon}(t_2,n_2))$ and draw conclusions about the subdiffusive nature of these processes, with particular attention to the behavior as $\varepsilon \to 0$. This complements previous work of the author, which in turn was influenced by work of McKinley, Yao, and Forest.

Key words. Gaussian processes, stochastic differential equations, singular perturbation.

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1. Introduction

In [7], following earlier work of [5] and [4], we studied the behavior of Gaussian processes that can be described as follows. Let $\ell^2(\mathbb{Z})$ denote the space of functions $a:\mathbb{Z}\to\mathbb{C}$ such that $\sum |a(n)|^2<\infty$ (here $\mathbb Z$ denotes the set of integers and $\mathbb C$ the set of complex numbers), and let L be a negative semidefinite, self adjoint operator on $\ell^2(\mathbb{Z})$. We assume finitely supported elements of $\ell^2(\mathbb{Z})$ belong to the domain of L, so we can write

$$
Ly(n) = \sum_{m \in \mathbb{Z}} \lambda(n, m) y(m).
$$
 (1.1)

Self adjointness implies $\lambda(n,m) = \overline{\lambda(m,n)}$. We assume

$$
\lambda(n,m) \in \mathbb{R}, \quad \text{hence} \quad \lambda(m,n) = \lambda(n,m). \tag{1.2}
$$

The process $x(t) = (x(t, n))$ studied in [7] solves the infinite system of stochastic differential equations

$$
dx(t,n) = Lx(t,n)dt + \sigma dW_n(t), \quad x(0,n) = 0,
$$
\n(1.3)

for $n \in \mathbb{Z}$, $t \geq 0$. Here W_n are independent, identically distributed Wiener processes. The system (1.3) provides a model for the motion of a polymer, pictured as a network of beads that interact and are also independently randomly jittered, as in Brownian motion. The particular case

$$
Ly(n) = y(n-1) - 2y(n) + y(n+1)
$$
\n(1.4)

gives rise to what is called the Rouse chain model; see [5] and [4] for further details and references to the literature.

In [7], the solution to (1.3) was constructed in the form

$$
x(t,n) = \sigma \int_0^t \sum_{m \in \mathbb{Z}} h(t-s,n,m) dW_m(s), \qquad (1.5)
$$

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where

$$
e^{tL}y(n) = \sum_{m} h(t, n, m)y(m).
$$
\n(1.6)

It was shown that the series in (1.5) converges and defines a Gaussian process, with mean 0. Formulas were derived for $\mathbb{E}(x(t_1,n_1)x(t_2,n_2))$, with special consideration of

$$
\mathbb{E}(x(t,n)^2), \text{ and } \mathbb{E}(|x(t,n_1)-x(t,n_2)|^2). \tag{1.7}
$$

The analysis of the first expectation in (1.7) recovered results of [5] and [4] on subdiffusivity of $x(t,n)$, and the analysis of the second expectation in (1.7), and also of $\mathbb{E}(x(t,n_1)x(t,n_2))$, provided information on the joint distribution of $x(t,n_1)$ and $x(t, n₂)$.

As pointed out in [5] and [4], the system (1.3) is the $\varepsilon = 0$ case of the second order system

$$
\varepsilon x_{\varepsilon}''(t,n) + x_{\varepsilon}'(t,n) = Lx_{\varepsilon}(t,n) + \sigma W_n'(t), \tag{1.8}
$$

with prime denoting the t-derivative. Here ε is proportional to the mass of each bead. It is reasonable to consider ε to be positive but quite small. Thus it is of interest to study the solution $x_{\varepsilon}(t,n)$ to (1.8), with particular interest in the behavior as $\varepsilon \searrow 0$. This paper addresses that task. We take initial data

$$
x_{\varepsilon}(0,n) = 0, \quad x'_{\varepsilon}(0,n) = 0, \quad \forall n \in \mathbb{Z}.
$$
 (1.9)

Since (1.8) changes type when ε reaches 0, this is a singular perturbation problem. We first tackle it under an additional condition on L, namely that it be a bounded operator on $\ell^2(\mathbb{Z})$, with operator norm $||L|| < \infty$. This condition holds for (1.4) and for many (arguably, for most) other examples arising in the bead-spring setting. Other examples include graph Laplacians, shown to be bounded in [7], in the case of infinite graphs, following results exposed for finite graphs in [1]. We produce a formula for the solution to (1.8) – (1.9) valid for

$$
0 < \varepsilon < \frac{1}{4||L||},\tag{1.10}
$$

and study its behavior as $\varepsilon \searrow 0$. (In §6 we drop the hypothesis that L be bounded and allow arbitrary $\varepsilon > 0$.)

To see how such a formula arises, let us rewrite (1.5) as

$$
x(t) = \sigma \int_0^t e^{(t-s)L} \, dW(s),\tag{1.11}
$$

to celebrate how it comes from Duhamel's formula. To obtain an analogue for (1.8), we set $v_{\varepsilon}(t) = x'_{\varepsilon}(t)$, i.e., $v_{\varepsilon}(t,n) = x'_{\varepsilon}(t,n)$, and rewrite (1.8) as a first order system

$$
\frac{d}{dt}\begin{pmatrix} x_{\varepsilon} \\ v_{\varepsilon} \end{pmatrix} = X_{\varepsilon} \begin{pmatrix} x_{\varepsilon} \\ v_{\varepsilon} \end{pmatrix} + \begin{pmatrix} 0 \\ \beta \sigma W'(t) \end{pmatrix},
$$
\n(1.12)

where

$$
X_{\varepsilon} = \begin{pmatrix} 0 & I \\ \beta L & -\beta I \end{pmatrix}.
$$
 (1.13)

Here and below, we set

$$
\beta = \frac{1}{\varepsilon}.\tag{1.14}
$$

In (1.8) and (1.12) , we use the "white noise" formalism $W'(t)$. The system (1.12) is of course a Wiener-Itô stochastic differential equation, which can be written

$$
d\binom{x_{\varepsilon}}{v_{\varepsilon}} = X_{\varepsilon}\binom{x_{\varepsilon}}{v_{\varepsilon}} dt + \binom{0}{\beta \sigma dW(t)}.
$$

Taking into account the initial data (1.9), the Duhamel formula gives

$$
\begin{aligned}\n\begin{pmatrix}\nx_{\varepsilon}(t) \\
v_{\varepsilon}(t)\n\end{pmatrix} &= \sigma \int_0^t e^{(t-s)X_{\varepsilon}} \begin{pmatrix} 0 \\
\beta W'(s) \end{pmatrix} ds \\
&= \sigma \int_0^t e^{(t-s)X_{\varepsilon}} \begin{pmatrix} 0 \\
\beta dW(s) \end{pmatrix}.\n\end{aligned} \tag{1.15}
$$

To compute $e^{sX_{\varepsilon}}$, we note that by the spectral theorem (cf. [6], Chapter 7) we can treat L as a real number and X_{ε} as a real 2×2 matrix, with "eigenvalues"

$$
\lambda_{\pm}(\beta, L) = -\frac{\beta}{2}I \pm \frac{\beta}{2}(I + 4\varepsilon L)^{1/2},\tag{1.16}
$$

and "eigenvectors"

$$
\begin{pmatrix} 1 \\ \lambda_{\pm}(\beta, L) \end{pmatrix} . \tag{1.17}
$$

One then calculates

$$
e^{tX_{\varepsilon}}\binom{0}{1} = \binom{(e^{t\lambda_+} - e^{t\lambda_-})/(\lambda_+ - \lambda_-)}{(\lambda_+ e^{t\lambda_+} - \lambda_- e^{t\lambda_-})/(\lambda_+ - \lambda_-)}.
$$
\n(1.18)

Thus (1.15) yields

$$
x_{\varepsilon}(t) = \sigma \int_0^t \left[A_\beta^+(t-s) - A_\beta^-(t-s)\right] dW(s),\tag{1.19}
$$

where

$$
A_{\beta}^{\pm}(s) = (I + 4\varepsilon L)^{-1/2} e^{s\lambda_{\pm}(\beta, L)}, \qquad (1.20)
$$

and $\lambda_{+}(\beta,L)$, given by (1.16), are bounded, negative semidefinite, self adjoint operators on $\ell^2(\mathbb{Z})$, as long as (1.10) holds. We have the task to show that the right side of (1.19) is a well defined Gaussian process and to investigate its properties, with particular attention to the behavior as $\varepsilon \searrow 0$, i.e., as $\beta \nearrow \infty$.

For use in subsequent sections, in §2 we collect some results on a class of vector stochastic integrals of the form

$$
x(t) = \int_0^t A(t - s) dW(s),
$$
\n(1.21)

where $\{A(s), A(s)^*: s \geq 0\}$ are strongly continuous families of bounded linear operators on $\ell^2(\mathbb{Z})$. Here, $x(t) = (x(t,n), n \in \mathbb{Z})$. We show that for each n, $x(t,n)$ is well defined

and is a continuous function of $t \in [0, \infty)$ with values in $L^2(X, \nu)$, where (X, ν) is a naturally constructed probability space (see §2 for details). Also, for each $t \geq 0$, $n \in \mathbb{Z}$, $x(t,n)$ is a Gaussian random variable with mean zero. These results can be established via material in Chapter 4 of [2], but the setting here is more elementary. For the convenience of readers not familiar with infinite dimensional stochastic analysis, we give short, direct demonstrations of the needed formulas, as a consequence of classical work of Paley, Wiener, and Zygmund. Formulas established in §2 include

$$
\mathbb{E}(|x(t,n)|^2) = \int_0^t \|A(s)^*\delta_n\|_{\ell^2}^2 ds,
$$
\n(1.22)

and more generally

$$
\mathbb{E}(x(t,n_1)\overline{x(t,n_2)}) = \int_0^t (A(s)^* \delta_{n_2}, A(s)^* \delta_{n_1}) ds,
$$
\n(1.23)

where $\{\delta_n : n \in \mathbb{Z}\}\$ is the orthonormal basis of $\ell^2(\mathbb{Z})$ given by $\delta_n(m) = 1$ if $m = n, 0$ otherwise. If $A(s)$ and $A(s)^*$ commute for all s, one can erase the asterisks in (1.22) and (1.23).

In §3 we apply the results of §2 to $A(s) = A_{\beta}^{\pm}(s)$, given by (1.20), and construct

$$
x_{\varepsilon}(t,n) = x_{\varepsilon}^{+}(t,n) - x_{\varepsilon}^{-}(t,n),
$$
\n(1.24)

when (1.10) holds. Here $x_{\varepsilon}^{\pm}(t,n)$ is the *n*th component of

$$
x_{\varepsilon}^{\pm}(t) = \sigma \int_0^t A_{\beta}^{\pm}(t-s) dW(s),
$$

with A_{β}^{\pm} as in (1.20). We compare $x_{\varepsilon}(t,n)$ to the solution to (1.3), given by (1.5), which we now denote $x_0(t,n)$. We show that

$$
\mathbb{E}(|x_{\varepsilon}^+(t,n)-x_0(t,n)|^2) \le C\varepsilon \mathbb{E}(x_0(t,n)^2),\tag{1.25}
$$

and

$$
\mathbb{E}(x_{\varepsilon}^-(t,n)^2) \le C\sigma^2(1 - e^{-t/\varepsilon})\varepsilon,\tag{1.26}
$$

provided $0 < \varepsilon \le a / ||L||$, with $a < 1/4$; see Theorem 3.1. These estimates imply that whenever $x_0(t,n)$ is subdiffusive, i.e.,

$$
\frac{1}{t}\mathbb{E}(x_0(t,n)^2)\longrightarrow 0 \text{ as } t \nearrow \infty,
$$
\n(1.27)

the processes $x_{\varepsilon}(t,n)$ are uniformly subdiffusive, for ε in such an interval.

In §4 we note that the processes $x_{\varepsilon}^{\pm}(t,n)$ are differentiable (as functions of t with values in $L^2(X,\nu)$, for ε satisfying (1.10), and study

$$
v_{\varepsilon}^{\pm}(t,n) = \frac{d}{dt}x_{\varepsilon}^{\pm}(t,n). \tag{1.28}
$$

At least one of these must blow up as $\varepsilon \searrow 0$, since $x_0(t,n)$ is not differentiable; as it turns out, $v_{\varepsilon}^{-}(t, n)$ blows up. We show that

$$
\mathbb{E}(v_{\varepsilon}^+(t,n)^2) \le C \mathbb{E}(x_0(t,n)^2),\tag{1.29}
$$

but

$$
\mathbb{E}(v_{\varepsilon}^-(t,n)^2) \ge \frac{\sigma^2}{4\varepsilon} (1 - e^{-t/\varepsilon}).
$$
\n(1.30)

In §5 we convert formulas for $\mathbb{E}(x_\varepsilon(t,n)^2)$ into integral formulas, arising from a spectral representation of L, and examine the asymptotic behavior as $t\nearrow\infty$, including more precise versions of the subdiffusivity result (1.27) and their counterparts for $\mathbb{E}(x_{\varepsilon}(t,n)^2)$; see Theorems 5.1–5.2.

Results of §§3–5 use the hypothesis (1.10). We obtain estimates valid uniformly for $0 < \varepsilon < a/||L||$, given $a < 1/4$. In §6 we extend the scope of our investigation, in two ways. First, we replace (1.10) by

$$
0 < \varepsilon < \infty. \tag{1.31}
$$

Second, we remove the hypothesis that L be bounded. In this more general setting, frequently $-1/4\varepsilon$ belongs to the spectrum of L and represents a transition from overdamping to underdamping in the system (1.8). The operators $A_{\beta}^{\pm}(s)$ in (1.20) are then not bounded, and the processes $x_{\varepsilon}^{\pm}(t,n)$ do not exist. However,

$$
A_{\beta}(s) = A_{\beta}^{+}(s) - A_{\beta}^{-}(s)
$$
\n(1.32)

is bounded. In fact, from (1.20) we obtain

$$
A_{\beta}(s) = s\beta e^{-s\beta/2} H\left(\frac{s\beta}{2}(I + 4\varepsilon L)^{1/2}\right),\tag{1.33}
$$

where H is the entire holomorphic, even function on $\mathbb C$ given by

$$
H(z) = \frac{\sinh z}{z}, \quad H(0) = 1.
$$
 (1.34)

Using this, we show that the processes $x_\varepsilon(t,n)$ exist. We obtain formulas for $\mathbb{E}(x_{\varepsilon}(t,n)^2)$, etc., extending those obtained earlier for ε satisfying (1.10). Making use of these results, we extend the scope of results of §5. Our main results in this section are given in Theorems 6.1–6.2.

2. A class of vector stochastic integrals

In this section we provide some useful formulas for vector stochastic integrals of the form

$$
x(t) = \int_0^t A(t - s) dW(s),
$$
\n(2.1)

where, for each $s \geq 0$,

$$
A(s) : \ell^2(\mathbb{Z}) \longrightarrow \ell^2(\mathbb{Z})
$$
\n^(2.2)

is a bounded linear operator. For simplicity we assume

 $A(s)$ and $A(s)^*$ are strongly continuous in $s \in [0,\infty)$, (2.3)

though the calculations below will make it clear that we can relax this hypothesis. Written out more fully, (2.1) takes the form

$$
x(t,n) = \int_0^t \sum_{m \in \mathbb{Z}} a(t-s,n,m) dW_m(s), \qquad (2.4)
$$

where, for $y \in \ell^2(\mathbb{Z}),$

$$
A(s)y(n) = \sum_{m} a(s,n,m)y(m).
$$
 (2.5)

The operators arising in (1.19) are self adjoint (for $0 < \varepsilon < 1/4||L||$) and reality preserving, but we do not need these properties for the development here. Consequently, the processes (2.4) might be complex valued. Note the adjoint $A(s)^*$ of $A(s)$ satisfies

$$
A(s)^* y(m) = \sum_{m} a^*(s, n, m) y(m), \quad a^*(s, n, m) = \overline{a(s, m, n)},
$$
\n(2.6)

and that

$$
a(s,n,m) = A(s)\delta_m(n), \quad a^*(s,n,m) = A(s)^*\delta_m(n), \tag{2.7}
$$

where $\delta_n \in \ell^2(\mathbb{Z})$ is given by

$$
\delta_n(m) = \delta_{n,m} = \begin{cases} 1 & \text{if } m = n, \\ 0 & \text{otherwise.} \end{cases}
$$
 (2.8)

As stated in the introduction, $\{W_n : n \in \mathbb{Z}\}\$ is a collection of independent, identically distributed Wiener processes. In more detail, let $B(t)$ be the Wiener process (Brownian motion), which is a continuous family $B(t) \in L^2(\Omega, \mu)$, where Ω is path space and μ is Wiener measure. Then set $\Omega_n = \Omega$, $\mu_n = \mu$, for $n \in \mathbb{Z}$, and take the product space (with product measure)

$$
(X,\nu) = \prod_{n \in \mathbb{Z}} (\Omega_n, \mu_n).
$$
\n(2.9)

We obtain (2.4) as

$$
x(t,n) = \lim_{K \to \infty} \sum_{m=-K}^{K} \xi_m(t,n),
$$
\n(2.10)

where

$$
\xi_m(t,n) = \int_0^t a(t-s,n,m) \, dW_m(s). \tag{2.11}
$$

Our first task is to establish convergence in $L^2(X,\nu)$ of the right side of (2.10). Note that

$$
m \neq m' \Longrightarrow \xi_m(t, n) \perp \xi_{m'}(t, n) \quad \text{in} \quad L^2(X, \nu), \tag{2.12}
$$

so it suffices to bound $\sum_m \mathbb{E}(|\xi_m(t,n)|^2)$. To get this, note that

$$
\mathbb{E}(|\xi_m(t,n)|^2) = \int_0^t a(s,n,m)a^*(s,m,n)ds,
$$
\n(2.13)

which is the classical Paley-Wiener-Zygmund identity (cf. [3], §2.1). Hence

$$
\sum_{m} \mathbb{E}(|\xi_m(t,n)|^2) = \sum_{m} \int_0^t a(s,n,m) a^*(s,m,n) ds
$$

$$
= \sum_{m} \int_{0}^{t} a(s,n,m)A(s)^{*}\delta_{n}(m) ds
$$

\n
$$
= \int_{0}^{t} A(s)A(s)^{*}\delta_{n}(n) ds
$$

\n
$$
= \int_{0}^{t} (A(s)A(s)^{*}\delta_{n}, \delta_{n}) ds
$$

\n
$$
= \int_{0}^{t} ||A(s)^{*}\delta_{n}||_{\ell^{2}}^{2} ds.
$$
 (2.14)

Here $\delta_n \in \ell^2(\mathbb{Z})$ is given by (2.8). Thus we have convergence in (2.10), and

$$
\mathbb{E}(|x(t,n)|^2) = \int_0^t \|A(s)^*\delta_n\|_{\ell^2}^2 ds.
$$
\n(2.15)

The nature of the convergence implies that for each $n \in \mathbb{Z}$, $t \geq 0$, $x(t,n)$ is a Gaussian random variable on (X,ν) with mean 0.

We next aim to show that, under the hypotheses in (2.3) , $x(t,n)$ is a continuous function of $t \in [0, \infty)$ with values in $L^2(X, \nu)$, for each n. In preparation for this, we note that

$$
\mathbb{E}(x(t_1,n)\overline{x(t_2,n)}) = \sum_{k} \int_0^{t_1 \wedge t_2} a(t_1-s,n,k)a^*(t_2-s,k,n) ds.
$$
 (2.16)

Here $t_1 \wedge t_2 = \min(t_1, t_2)$. We have

$$
\sum_{k} a(t_1 - s, n, k)a^*(t_2 - s, k, n) = \sum_{k} a(t_1 - s, n, k)A(t_2 - s)^*\delta_n(k)
$$

= $A(t_1 - a)A(t_2 - s)^*\delta_n(n)$
= $(A(t_2 - s)^*\delta_n, A(t_1 - s)^*\delta_n),$ (2.17)

so

$$
\mathbb{E}(x(t_1,n),\overline{x(t_2,n)}) = \int_0^{t_1 \wedge t_2} (A(t_2-s)^* \delta_n, A(t_1-s)^* \delta_n) ds. \tag{2.18}
$$

Now

$$
\mathbb{E}(|x(t_1,n)-x(t_2,n)|^2) = \mathbb{E}(x(t_1,n)^2) + \mathbb{E}(x(t_2,n)^2) - 2\operatorname{Re}\mathbb{E}(x(t_1,n)\overline{x(t_2,n)}),
$$
\n(2.19)

so (2.18) gives (say if $0 \le t_1 \le t_2$)

$$
\mathbb{E}(|x(t_1,n) - x(t_2,n)|^2)
$$
\n
$$
= \int_{t_1}^{t_2} ||A(t_2 - s)^* \delta_n||_{\ell^2}^2 ds + \int_0^{t_1 \wedge t_2} \left\{ (A(t_1 - s)^* \delta_n, A(t_1 - s)^* \delta_n) + (A(t_2 - s)^* \delta_n, A(t_2 - s)^* \delta_n) - 2\operatorname{Re}(A(t_2 - s)^* \delta_n, A(t_1 - s)^* \delta_n) \right\} ds. \quad (2.20)
$$

The first integral on the right side of (2.20) is $\leq C|t_1-t_2|$. We can write the second integral as

$$
\operatorname{Re} \int_0^{t_1 \wedge t_2} \left\{ \left([A(t_1 - s)^* - A(t_2 - s)^*] \delta_n, A(t_1 - s)^* \delta_n \right) + (A(t_2 - s)^* \delta_n, [A(t_2 - s)^* - A(t_1 - s)^*] \delta_n \right\} ds.
$$
\n(2.21)

Given (2.3), the fact that (2.21) tends to 0 as $t_1 \rightarrow t_2$ follows readily from the Lebesgue dominated convergence theorem. Let us summarize what we have established.

PROPOSITION 2.1. Under the hypotheses (2.2) – (2.3) , the formula (2.4) gives for each $n \in \mathbb{Z}$ a well defined $x(t,n)$, a continuous function of $t \in [0,\infty)$ with values in $L^2(X,\nu)$, for each t,n a Gaussian randon variable with mean 0, satisfying the identities (2.15) and (2.18).

We complement (2.15) with the following computation, derived similarly.

$$
\mathbb{E}(x(t,n_1)\overline{x(t,n_2)}) = \sum_{k} \int_0^t a(t-s,n_1,k)\overline{a(t-s,n_2,k)}ds
$$

=
$$
\sum_{k} \int_0^t a(t-s,n_1,k)a^*(t-s,k,n_2)ds.
$$
 (2.22)

Parallel to (2.17), we then get

$$
\mathbb{E}(x(t,n_1)\overline{x(t,n_2)}) = \int_0^t (A(t-s)^*\delta_{n_2}, A(t-s)^*\delta_{n_1}) ds
$$

=
$$
\int_0^t (A(s)^*\delta_{n_2}, A(s)^*\delta_{n_1}) ds.
$$
 (2.23)

Combining (2.15) and (2.23) , we have

$$
\mathbb{E}(|x(t,n_1) - x(t,n_2)|^2)
$$
\n
$$
= \mathbb{E}(|x(t,n_1)|^2) + \mathbb{E}(|x(t,n_2)|^2) - 2\operatorname{Re}\mathbb{E}(x(t,n_1)\overline{x(t,n_2)})
$$
\n
$$
= \int_0^t \left\{ \|A(s)^*\delta_{n_1}\|_{\ell^2}^2 + \|A(s)^*\delta_{n_2}\|_{\ell^2}^2 - 2\operatorname{Re}(A(s)^*\delta_{n_2}, A(s)^*\delta_{n_1}) \right\} ds
$$
\n
$$
= \int_0^t \|A(s)^*(\delta_{n_1} - \delta_{n_2})\|_{\ell^2}^2 ds. \tag{2.24}
$$

We now give a condition under which the components $x(t,n)$ of the process (2.1) are differentiable, as functions of t with values in $L^2(X, \nu)$. Let us add to (2.3) the hypothesis

$$
A'(s) \text{ and } A'(s)^* \text{ are strongly continuous in } s \in [0, \infty). \tag{2.25}
$$

Then, as in the scalar case, Wiener's integration by parts formula holds for (2.1):

$$
x(t) = \int_0^t A'(t-s)W(s)ds + A(0)W(t).
$$
 (2.26)

We have the following.

PROPOSITION 2.2. In the setting of Proposition 2.1, if also (2.25) holds and $A(0) = 0$, then $x(t,n)$ is differentiable for each $n \in \mathbb{Z}$, and $x'(t,n)$ is a continuous function of $t \in [0, \infty)$ with values in $L^2(X, \nu)$.

Proof. Let us temporarily assume that (2.25) also holds for $A''(s)$. Then we differentiate (2.26) and get (provided $A(0)=0$)

$$
x'(t) = \int_0^t A''(t-s)W(s)ds + A'(0)W(t).
$$
 (2.27)

Applying (2.26) with A replaced by A' then gives

$$
x'(t) = \int_0^t A'(t-s)dW(s).
$$
 (2.28)

A mollification and approximation argument gives (2.28) without the additional assumption on A'' . \Box

Returning to the computations (2.13) – (2.23) , note that if $A(s)$ is self adjoint for all s, all the asterisks can be removed, and if these operators are reality preserving, all the overlines can be removed. Furthermore,

$$
A(s)^* A(s) = A(s)A(s)^* \n\implies ||A(s)^*\delta_n||_{\ell^2} = ||A(s)\delta_n||_{\ell^2}, \text{ and} \n(A(s)^*\delta_{n_2}, A(s)^*\delta_{n_1}) = (A(s)\delta_{n_2}, A(s)\delta_{n_1}),
$$
\n(2.29)

so we have the following.

PROPOSITION 2.3. In the setting of Proposition 2.1, if also $A(s)$ is normal for all $s \geq 0$, then

$$
\mathbb{E}(|x(t,n)|^2) = \int_0^t \|A(s)\delta_n\|_{\ell^2}^2 ds,
$$
\n(2.30)

and more generally

$$
\mathbb{E}(x(t,n_1)\overline{x(t,n_2)}) = \int_0^t (A(s)\delta_{n_2}, A(s)\delta_{n_1}) ds.
$$
 (2.31)

3. The processes $x_{\varepsilon}(t,n)$

Results of §2 on the vector stochastic integral $\int_0^t A(t-s)dW(s)$ apply to (1.19) with

$$
A(s) = \sigma[A_{\beta}^{+}(s) - A_{\beta}^{-}(s)],
$$

\n
$$
A_{\beta}^{\pm}(s) = (I + 4\varepsilon L)^{-1/2} e^{s\lambda_{\pm}(\beta, L)},
$$

\n
$$
\lambda_{\pm}(\beta, L) = -\frac{\beta}{2} I \pm \frac{\beta}{2} (I + 4\varepsilon L)^{1/2}.
$$
\n(3.1)

In the current setting, L is a bounded, reality preserving, negative semidefinite, self adjoint operator on $\ell^2(\mathbb{Z})$, $0 < \varepsilon < 1/4||L||$, and $\beta = 1/\varepsilon$. Hence $\lambda_{\pm}(\beta, L)$ are negative semidefinite, self adjoint operators on $\ell^2(\mathbb{Z})$. Thus, for each such ε , $x_{\varepsilon}(t)$ = $(x_\varepsilon(t,n), n\in\mathbb{Z})$ has the property that, for each $n\in\mathbb{Z}$, $x_\varepsilon(t,n)$ is a continuous function of $t \in [0, \infty)$ with values in $L^2(X, \nu)$, and for each $t \geq 0$ is a real valued Gaussian random variable with mean 0. For further analysis, it is convenient (using (1.19)) to write

$$
x_{\varepsilon}(t,n) = x_{\varepsilon}^{+}(t,n) - x_{\varepsilon}^{-}(t,n),
$$
\n(3.2)

where

$$
x_{\varepsilon}^{\pm}(t) = \sigma \int_0^t A_{\beta}^{\pm}(t-s) dW(s).
$$
 (3.3)

The formula (2.15) gives

$$
\mathbb{E}(x_{\varepsilon}^{\pm}(t,n)^{2}) = \sigma^{2} \int_{0}^{t} \|A_{\beta}^{\pm}(s)\delta_{n}\|_{\ell^{2}}^{2} ds.
$$
 (3.4)

Note that Spec $\lambda_-(\beta,L) \subset (-\infty,-\beta/2]$, so we have the operator norm estimate

$$
||A_{\beta}^{-}(s)|| \le ||(I + 4\varepsilon L)^{-1/2}||e^{-s\beta/2}, \tag{3.5}
$$

and we get

$$
\mathbb{E}(x_{\varepsilon}^{-}(t,n)^{2}) \le C\sigma^{2}\varepsilon(1-e^{-\beta t}), \quad 0 < \varepsilon \le \frac{a}{\|L\|}, \quad a < \frac{1}{4}, \tag{3.6}
$$

with C independent of $t \in [0,\infty)$.

In order to analyze $x_{\varepsilon}^+(t,n)$, note that, as long as (1.10) holds,

$$
(I + 4\varepsilon L)^{1/2} = I + 2\varepsilon L \Phi(4\varepsilon L),\tag{3.7}
$$

with $\Phi(\lambda)$ given by

$$
(1+\lambda)^{1/2} = 1 + \frac{1}{2}\lambda - \frac{1}{8}\lambda^2 + \cdots
$$

= $1 + \frac{1}{2}\lambda \left(1 - \frac{1}{4}\lambda + \cdots\right)$
= $1 + \frac{1}{2}\lambda \Phi(\lambda).$ (3.8)

Note that $\Phi(\lambda)$ is holomorphic on $\{\lambda \in \mathbb{C} : |\lambda| < 1\}$ and

 $\Phi(0) = 1, \quad \Phi(\lambda) > 0 \quad \text{for} \quad \lambda \in (-1, 1).$ (3.9)

Hence

$$
\lambda_{+}(\beta, L) = L\Phi(4\varepsilon L),\tag{3.10}
$$

so

$$
A_{\beta}^{+}(s) = (I + 4\varepsilon L)^{-1/2} e^{sL\Phi(4\varepsilon L)}, \qquad (3.11)
$$

and

$$
x_{\varepsilon}^{+}(t) = \sigma \int_{0}^{t} (I + 4\varepsilon L)^{-1/2} e^{(t-s)L\Phi(4\varepsilon L)} dW(s).
$$
 (3.12)

Hence

$$
\mathbb{E}(x_{\varepsilon}^{+}(t,n)^{2}) = \sigma^{2} \int_{0}^{t} \|(I + 4\varepsilon L)^{-1/2} e^{sL\Phi(4\varepsilon L)} \delta_{n}\|_{\ell^{2}}^{2} ds
$$
\n
$$
= \sigma^{2} \int_{0}^{t} (e^{2sL\Phi(4\varepsilon L)} (I + 4\varepsilon L)^{-1} \delta_{n}, \delta_{n}) ds.
$$
\n(3.13)

If we set

$$
G(\lambda) = \int_0^1 e^{-s\lambda} ds = \begin{cases} \frac{1 - e^{-\lambda}}{\lambda} & , \lambda > 0, \\ 1 & , \lambda = 0, \end{cases}
$$
 (3.14)

we can write (3.13) as

$$
\mathbb{E}(x_{\varepsilon}^+(t,n)^2) = \sigma^2 t \big(G(-2tL\Phi(4\varepsilon L))(I + 4\varepsilon L)^{-1} \delta_n, \delta_n \big). \tag{3.15}
$$

In $\S 5$ we will investigate large t behavior of this.

At this point, it is natural to compare $x_{\varepsilon}^{+}(t)$ with the solution $x_{0}(t)$ to (1.3), given by (1.11), i.e.,

$$
x_0(t) = \sigma \int_0^t e^{(t-s)L} dW(s).
$$
 (3.16)

Note that, parallel to (3.13)–(3.15),

$$
\mathbb{E}(x_0(t,n)^2) = \sigma^2 \int_0^t \|e^{sL} \delta_n\|_{\ell^2}^2 ds
$$

= $\sigma^2 t (G(-2tL)\delta_n, \delta_n).$ (3.17)

Applying (2.15) to the difference of (3.12) and (3.16) gives

$$
\mathbb{E}(|x_{\varepsilon}^{+}(t,n)-x_{0}(t,n)|^{2})
$$
\n
$$
=\sigma^{2} \int_{0}^{t} \|[A_{\beta}^{+}(t-s)-e^{(t-s)L}]\delta_{n}\|_{\ell^{2}}^{2} ds
$$
\n
$$
=\sigma^{2} \int_{0}^{t} \|[(I+4\varepsilon L)^{-1/2} e^{sL\Phi(4\varepsilon L)} - e^{sL}]\delta_{n}\|_{\ell^{2}}^{2} ds.
$$
\n(3.18)

Using $(a+b)^2 \leq 2a^2 + 2b^2$, we can write

$$
\mathbb{E}(|x_{\varepsilon}^+(t,n)-x_0(t,n)|^2) \le 2\sigma^2(A_{\varepsilon}(t,n)+B_{\varepsilon}(t,n)),\tag{3.19}
$$

where

$$
A_{\varepsilon}(t,n) = \int_0^t \|(I + 4\varepsilon L)^{-1/2} [e^{sL\Phi(4\varepsilon L)} - e^{sL}] \delta_n\|_{\ell^2}^2 ds,
$$

\n
$$
B_{\varepsilon}(t,n) = \int_0^t \|[(I + 4\varepsilon L)^{-1/2} - I] e^{sL} \delta_n\|_{\ell^2}^2 ds.
$$
\n(3.20)

Noting that

$$
||(I + 4\varepsilon L)^{-1/2} - I|| \le C\varepsilon, \quad \text{for } 0 < \varepsilon \le \frac{a}{||L||}, \quad a < \frac{1}{4}, \tag{3.21}
$$

and comparing (3.17), we have

$$
\sigma^2 B_{\varepsilon}(t,n) \le C \varepsilon \mathbb{E}(x_0(t,n)^2). \tag{3.22}
$$

We also have

$$
A_{\varepsilon}(t,n) \le C\widetilde{A}_{\varepsilon}(t,n),\tag{3.23}
$$

where

$$
\widetilde{A}_{\varepsilon}(t,n) = \int_0^t \|\big[e^{sL\Phi(4\varepsilon L)} - e^{sL}\big]\delta_n\|_{\ell^2}^2 ds.
$$
\n(3.24)

To proceed, recall from (3.8) that

$$
(1+\lambda)^{1/2} = 1 + \frac{1}{2}\lambda \Phi(\lambda),
$$
\n(3.25)

where $\Phi(\lambda)$ is given by

$$
\Phi(\lambda) = 1 - \frac{1}{4}\lambda + \dots = 1 - \frac{\lambda}{4}\psi(\lambda), \quad \psi(0) = 1,
$$
\n(3.26)

with $\psi(\lambda)$ holomorphic in $\{\lambda \in \mathbb{C} : |\lambda| < 1\}$, real and positive for $\lambda \in (-1,1)$. The positivity can be seen from the concavity of $(1+\lambda)^{1/2}$, which implies $(1+\lambda)^{1/2} \leq 1+\lambda/2$ on $(-1,1)$, hence $\Phi(\lambda) \le 1$ on $[0,1)$ and ≥ 1 on $(-1,0]$. Hence

$$
e^{sL\Phi(4\varepsilon L)} - e^{sL} = \left(e^{-s\varepsilon L^2\psi(4\varepsilon L)} - I\right)e^{sL},\tag{3.27}
$$

and we have

$$
\widetilde{A}_{\varepsilon}(t,n) = \int_0^t \|(I - e^{-s\varepsilon L^2 \psi(4\varepsilon L)})e^{sL}\delta_n\|_{\ell^2}^2 ds,
$$
\n(3.28)

which gives

$$
\sigma^2 \widetilde{A}_{\varepsilon}(t,n) \le \sup_{0 < s < t} \|I - e^{-s\varepsilon L^2 \psi(4\varepsilon L)}\|^2 \mathbb{E}(x_0(t,n)^2). \tag{3.29}
$$

If we take $a \in (0,1/4)$ and set

$$
\alpha = \sup_{0 < \varepsilon \le a / ||L||} ||L^2 \psi(4\varepsilon L)||,\tag{3.30}
$$

then, since $L^2\psi(4\varepsilon L)$ is positive semidefinite, we have

$$
\sigma^2 \widetilde{A}_{\varepsilon}(t,n) \le (1 - e^{-\alpha \varepsilon t})^2 \mathbb{E}(x_0(t,n)^2),\tag{3.31}
$$

provided

$$
0 < \varepsilon \le \frac{a}{\|L\|}, \quad a < \frac{1}{4}.\tag{3.32}
$$

The factor in front of $\mathbb{E}(x_0(t,n)^2)$ in (3.31) is $O(\varepsilon)$ for t in each bounded interval in [0,∞), but one loses uniformity as $t\nearrow\infty$. In fact, (3.31) is not optimal. We proceed to derive a stronger estimate. Writing

$$
(I - e^{-s\varepsilon L^2 \psi(4\varepsilon L)})e^{sL} = (e^{sL/2} - e^{sL/2 - s\varepsilon L^2 \psi(4\varepsilon L)})e^{sL/2},
$$
\n(3.33)

we have

$$
\widetilde{A}_{\varepsilon}(t,n) \le A_{\varepsilon}^{\#}(t)^{2} \int_{0}^{t} \|e^{sL/2} \delta_{n}\|_{\ell^{2}}^{2} ds,
$$
\n(3.34)

with

$$
A_{\varepsilon}^{\#}(t) = \sup_{0 < s < t} \|e^{sL/2} - e^{sL/2 - s\varepsilon L^2 \psi(4\varepsilon L)}\|
$$
\n
$$
\leq \sup_{0 < s < t, 0 \leq \Lambda \leq \|L\|} \left| e^{-s\Lambda/2} - e^{-s\Lambda/2 - s\varepsilon \Lambda^2 \psi(-4\varepsilon \Lambda)} \right|,\tag{3.35}
$$

the latter inequality by the spectral theorem. Now, over the range $0 \leq \Lambda \leq ||L||$,

$$
\varphi = \Lambda \psi(-4\varepsilon \Lambda) \Longrightarrow 0 \le \varphi \le \Lambda_0,\tag{3.36}
$$

as long as (3.32) holds, where $\Lambda_0 = ||L|| \sup_{(-1,1)} \psi(\lambda)$, and we obtain

$$
A_{\varepsilon}^{\#}(t) \le \sup_{0 < s < t, \Lambda \ge 0, 0 \le \varphi \le \Lambda_0} \left| e^{-s\Lambda/2} - e^{-s\Lambda/2 - s\varepsilon\Lambda\varphi} \right|.
$$
\n(3.37)

Taking $s\Lambda \mapsto \Lambda$, we get

$$
A_{\varepsilon}^{\#}(t) \le \sup_{\Lambda \ge 0, 0 \le \varphi \le \Lambda_0} |e^{-\Lambda/2} - e^{-\Lambda/2 - \varepsilon \varphi \Lambda}|
$$

\n
$$
\le \sup_{\Lambda \ge 0, 0 \le \varphi \le \Lambda_0} \varepsilon \varphi \Lambda e^{-\Lambda/2}
$$

\n
$$
\le \varepsilon \Lambda_0,
$$
 (3.38)

since sup $\Lambda e^{-\Lambda/2} = 2/e < 1$. Note that this estimate is independent of t. Meanwhile,

$$
\sigma^2 \int_0^t \|e^{sL/2} \delta_n\|_{\ell^2}^2 ds = 2\sigma^2 \int_0^{t/2} \|e^{sL} \delta_n\|_{\ell^2}^2 ds
$$

= $2\mathbb{E}\Big(x_0 \Big(\frac{t}{2}, n\Big)^2\Big)$
 $\leq 2\mathbb{E}(x_0(t, n)^2),$ (3.39)

so (3.34) and (3.38) yield

$$
\sigma^2 \widetilde{A}_{\varepsilon}(t,n) \le 2\Lambda_0 \varepsilon \mathbb{E}(x_0(t,n)^2). \tag{3.40}
$$

Let us collect the main results established above.

THEOREM 3.1. As long as (3.32) holds, the formulas (3.2)–(3.3) give, for each $n \in \mathbb{Z}$, a mean zero Gaussian process $t \mapsto x_{\varepsilon}(t,n) = x_{\varepsilon}^+(t,n) - x_{\varepsilon}^-(t,n)$, a continuous function of $t \in [0,\infty)$ with values in $L^2(X,\nu)$. Furthermore, there exist $C, \alpha \in (0,\infty)$ such that when (3.32) holds and $x_0(t,n)$ is given by (3.16), then for all $n \in \mathbb{Z}$, $t \ge 0$,

$$
\mathbb{E}(|x_{\varepsilon}^+(t,n)-x_0(t,n)|^2) \le C\varepsilon \mathbb{E}(x_0(t,n)^2),\tag{3.41}
$$

and

$$
\mathbb{E}(x_{\varepsilon}^{-}(t,n)^{2}) \leq C\sigma^{2}(1 - e^{-t/\varepsilon})\varepsilon.
$$
\n(3.42)

We record formulas for the covariance of $x_{\varepsilon}^{\pm}(t,n_1)$ and $x_{\varepsilon}^{\pm}(t,n_2)$. By (2.23), we have (with coherent choice of signs)

$$
\mathbb{E}(x_{\varepsilon}^{\pm}(t,n_1)x_{\varepsilon}^{\pm}(t,n_2))
$$
\n
$$
= \sigma^2 \int_0^t (A_{\beta}^{\pm}(s)\delta_{n_1}, A_{\beta}^{\pm}(s)\delta_{n_2})
$$
\n
$$
= \sigma^2 \int_0^t ((I + 4\varepsilon L)^{-1} e^{2s\lambda_{\pm}(\beta, L)} \delta_{n_1}, \delta_{n_2}) ds. \tag{3.43}
$$

In particular, using (3.10),

$$
\mathbb{E}(x_{\varepsilon}^{+}(t,n_{1})x_{\varepsilon}^{+}(t,n_{2}))
$$
\n
$$
= \sigma^{2} \int_{0}^{t} ((I + 4\varepsilon L)^{-1} e^{2sL\Phi(4\varepsilon L)} \delta_{n_{1}}, \delta_{n_{2}}) ds
$$
\n
$$
= \sigma^{2} t \big((I + 4\varepsilon L)^{-1} G(-2tL\Phi(4\varepsilon L)) \delta_{n_{1}}, \delta_{n_{2}} \big).
$$
\n(3.44)

4. The processes $v_{\varepsilon}(t,n)$

From (3.1) we see that $A_{\beta}^{\pm}(0) = (I + 4\varepsilon L)^{-1/2}$ for $0 < \varepsilon < 1/4||L||$, so $A(0) = 0$ and, by Proposition 2.2, $x_{\varepsilon}(t,n)$ is differentiable, as a function of t, with values in $L^2(X,\nu)$, for each $n \in \mathbb{Z}$. By (2.28) ,

$$
x'_{\varepsilon}(t,n) = v_{\varepsilon}(t,n) = v_{\varepsilon}^{+}(t,n) - v_{\varepsilon}^{-}(t,n),
$$
\n(4.1)

with

$$
v_{\varepsilon}^{\pm}(t) = \sigma \int_0^t V_{\beta}^{\pm}(t-s)dW(s), \tag{4.2}
$$

where

$$
V_{\beta}^{\pm}(s) = \frac{d}{ds} A_{\beta}^{\pm}(s)
$$

= $(I + 4\varepsilon L)^{-1/2} \lambda_{\pm}(\beta, L) e^{s\lambda_{\pm}(\beta, L)}.$ (4.3)

As before, $\beta = 1/\varepsilon$. We will compute square expectations and verify, as one should expect, that $\mathbb{E}(v_{\varepsilon}(t,n)^2) \to \infty$ as $\varepsilon \searrow 0$. In fact, we separately examine $\mathbb{E}(v_{\varepsilon}(t,n)^2)$ and $\mathbb{E}(v_{\varepsilon}^-(t,n)^2)$, and see that only the latter blows up as $\varepsilon \searrow 0$.

To begin, we have

$$
\mathbb{E}(v_{\varepsilon}^{\pm}(t,n)^{2})
$$
\n
$$
= \sigma^{2} \int_{0}^{t} \|V_{\beta}^{\pm}(s)\delta_{n}\|_{\ell^{2}}^{2} ds
$$
\n
$$
= \sigma^{2} \int_{0}^{t} \|(I + 4\varepsilon L)^{-1/2}\lambda_{\pm}(\beta, L)e^{s\lambda_{\pm}(\beta, L)}\delta_{n}\|_{\ell^{2}}^{2} ds.
$$
\n(4.4)

Recalling from (3.10) that $\lambda_+(\beta, L) = L\Phi(4\varepsilon L)$, we have, for

$$
0 < \varepsilon \le \frac{a}{\|L\|}, \quad a < \frac{1}{4},\tag{4.5}
$$

that

$$
\mathbb{E}(v_{\varepsilon}^{+}(t,n)^{2})
$$
\n
$$
= \sigma^{2} \int_{0}^{t} \|L\Phi(4\varepsilon L)(I + 4\varepsilon L)^{-1/2} e^{sL\Phi(4\varepsilon L)} \delta_{n} \|_{\ell^{2}}^{2} ds
$$
\n
$$
\leq C \mathbb{E}(x_{\varepsilon}^{+}(t,n)^{2})
$$
\n
$$
\leq C' \mathbb{E}(x_{0}(t,n)^{2}),
$$
\n(4.6)

the first inequality by (3.13), given the operator norm bound $||L\Phi(4\varepsilon L)|| \leq C$, and the second by (3.41) .

To proceed, we have

$$
\mathbb{E}(v_{\varepsilon}^{-}(t,n)^{2})
$$
\n
$$
= \sigma^{2} \int_{0}^{t} ((I + 4\varepsilon L)^{-1} \lambda_{-}(\beta, L)^{2} e^{2s\lambda_{-}(\beta, L)} \delta_{n}, \delta_{n}) ds
$$
\n
$$
= -\frac{\sigma^{2}}{2} ((I + 4\varepsilon L)^{-1} \lambda_{-}(\beta, L)(I - e^{2t\lambda_{-}(\beta, L)}) \delta_{n}, \delta_{n}). \tag{4.7}
$$

Now, as long as (4.5) holds, we have, via the spectral theorem,

$$
\operatorname{Spec}(I + 4\varepsilon L)^{-1} \subset [1, \infty),
$$

\n
$$
\operatorname{Spec}\lambda_{-}(\beta, L) \subset [-\beta, -\beta/2],
$$

\n
$$
\operatorname{Spec}(I - e^{2t\lambda_{-}(\beta, L)}) \subset [1 - e^{-\beta t}, 1],
$$
\n(4.8)

and hence

$$
\text{Spec} - (I + 4\varepsilon L)^{-1} \lambda_{-}(\beta, L)(I - e^{2t\lambda_{-}(\beta, L)}) \subset \left[\frac{\beta}{2}(1 - e^{-t\beta}), \infty\right). \tag{4.9}
$$

The variational characterization of the bottom of the spectrum for a positive definite, self adjoint operator, applied to the last inner product in (4.7), then gives

$$
\mathbb{E}(v_{\varepsilon}^-(t,n)^2) \ge \frac{\sigma^2}{4\varepsilon}(1 - e^{-t/\varepsilon}),\tag{4.10}
$$

as long as (4.5) holds. The right side of (4.10) clearly blows up as $\varepsilon \searrow 0$, for each $t > 0$.

5. Spectral representation, asymptotics, and subdiffusivity

Let $L:\ell^2(\mathbb{Z})\to\ell^2(\mathbb{Z})$ be a bounded, negative, self adjoint operator, as described in §1. The spectral theorem (cf. [6], Theorem VII.3) implies there is a measure space (S,γ) , a unitary map

$$
\mathcal{F}: \ell^2(\mathbb{Z}) \longrightarrow L^2(S, \gamma), \tag{5.1}
$$

and a function

$$
\Lambda\!\in\! L^\infty(S,\gamma),\quad \Lambda\!\geq\! 0,\quad \|\Lambda\|_{L^\infty}\!=\!\|L\|,\tag{5.2}
$$

such that for each $y \in \ell^2(\mathbb{Z}), t \geq 0$,

$$
\mathcal{F}Ly(\theta) = -\Lambda(\theta)\mathcal{F}y(\theta), \quad \theta \in S.
$$
\n(5.3)

Consequently,

$$
\mathcal{F}e^{tL}y(\theta) = e^{-t\Lambda(\theta)}\mathcal{F}y(\theta),
$$

\n
$$
\mathcal{F}\Phi(4\varepsilon L)y(\theta) = \Phi(-4\varepsilon\Lambda(\theta))\mathcal{F}y(\theta),
$$
\n(5.4)

etc. The orthonormal basis $\{\delta_n\}$ of $\ell^2(\mathbb{Z})$ gives rise to the orthonormal basis $\{e_n\}$ of $L^2(S,\gamma),$

$$
e_n = \mathcal{F}\delta_n. \tag{5.5}
$$

Using these ingredients, we can rewrite the formula (3.17) for the square expectation of $x_0(t,n)$ as

$$
\mathbb{E}(x_0(t,n)^2) = \sigma^2 \int_S \int_0^t e^{-2s\Lambda(\theta)} |e_n(\theta)|^2 ds d\gamma(\theta)
$$

= $\sigma^2 t \int_S G(2t\Lambda(\theta)) |e_n(\theta)|^2 d\gamma(\theta).$ (5.6)

Similarly, (3.13) – (3.15) yield

$$
\mathbb{E}(x_{\varepsilon}^{+}(t,n)^{2})
$$
\n
$$
= \sigma^{2} \int_{S} \int_{0}^{t} (1 - 4\varepsilon \Lambda(\theta))^{-1} e^{-2s\Lambda_{\varepsilon}(\theta)} |e_{n}(\theta)|^{2} ds d\gamma(\theta)
$$
\n
$$
= \sigma^{2} t \int_{S} (1 - 4\varepsilon \Lambda(\theta))^{-1} G(2t\Lambda_{\varepsilon}(\theta)) |e_{n}(\theta)|^{2} d\gamma(\theta), \qquad (5.7)
$$

for

$$
0 < \varepsilon < \frac{1}{4||L||}.\tag{5.8}
$$

Here,

$$
\Lambda_{\varepsilon}(\theta) = \Lambda(\theta)\Phi(-4\varepsilon\Lambda(\theta)) \n= \Lambda(\theta)\left(1 + \varepsilon\Lambda(\theta)\psi(-4\varepsilon\Lambda(\theta))\right),
$$
\n(5.9)

with Φ as in (3.7)–(3.10) and ψ as in (3.26). Note that, as long as (5.8) holds, $\psi(-4\varepsilon\Lambda(\theta))\geq 0$. More generally, by (3.44),

$$
\mathbb{E}(x_{\varepsilon}^{+}(t, n_{1})x_{\varepsilon}^{+}(t, n_{2}))
$$
\n
$$
= \sigma^{2} \int_{S} \int_{0}^{t} (1 - 4\varepsilon \Lambda(\theta))^{-1} e^{-2s\Lambda_{\varepsilon}(\theta)} e_{n_{1}}(\theta) \overline{e_{n_{2}}(\theta)} ds d\gamma(\theta)
$$
\n
$$
= \sigma^{2} t \int_{S} (1 - 4\varepsilon \Lambda(\theta))^{-1} G(2t\Lambda_{\varepsilon}(\theta)) e_{n_{1}}(\theta) \overline{e_{n_{2}}(\theta)} d\gamma(\theta).
$$
\n(5.10)

Let us specialize to the case that L is of convolution type:

$$
Ly(n) = \sum_{m} \lambda(n-m)y(m). \tag{5.11}
$$

A special case is given in (1.4), for the Rouse chain model. The convolution case was also emphasized in [5] and [4]. In this case, we can take

$$
S = S1 = \mathbb{R}/(2\pi\mathbb{Z}), \quad d\gamma(\theta) = d\theta/2\pi,
$$

\n
$$
\mathcal{F}y(\theta) = \hat{y}(\theta) = \sum_{n} y(n)e^{in\theta}, \quad e_n(\theta) = e^{in\theta}, \quad \Lambda(\theta) = -\hat{\lambda}(\theta).
$$
\n(5.12)

In such a case, (5.6) – (5.10) become

$$
\mathbb{E}(x_0(t,n)^2) = \frac{\sigma^2 t}{2\pi} \int_{S^1} G(2t\Lambda(\theta)) d\theta,
$$
\n(5.13)

$$
\mathbb{E}(x_{\varepsilon}^+(t,n)^2) = \frac{\sigma^2 t}{2\pi} \int_{S^1} (1 - 4\varepsilon \Lambda(\theta))^{-1} G\big(2t\Lambda_{\varepsilon}(\theta)\big) d\theta,\tag{5.14}
$$

and

$$
\mathbb{E}(x_{\varepsilon}^{+}(t,n_{1})x_{\varepsilon}^{+}(t,n_{2}))
$$
\n
$$
=\frac{\sigma^{2}t}{2\pi}\int_{S^{1}}(1-4\varepsilon\Lambda(\theta))^{-1}G\big(2t\Lambda_{\varepsilon}(\theta)\big)e^{i(n_{1}-n_{2})\theta}d\theta.
$$
\n(5.15)

Note that the reality condition (1.2) implies

$$
\lambda(-n) = \lambda(n)
$$
, hence $\Lambda(-\theta) = \Lambda(\theta)$. (5.16)

Taking this into account, a short computation yields

$$
\mathbb{E}(|x_{\varepsilon}^{+}(t,n_{1})-x_{\varepsilon}^{+}(t,n_{2})|^{2})
$$
\n
$$
=\mathbb{E}(x_{\varepsilon}^{+}(t,n_{1})^{2})+\mathbb{E}(x_{\varepsilon}^{+}(t,n_{2})^{2})-2\mathbb{E}(x_{\varepsilon}^{+}(t,n_{1})x_{\varepsilon}^{+}(t,n_{2}))
$$
\n
$$
=\frac{4\sigma^{2}t}{2\pi}\int_{S^{1}}(1-4\varepsilon\Lambda(\theta))^{-1}G(2t\Lambda_{\varepsilon}(\theta))\sin^{2}\frac{(n_{1}-n_{2})\theta}{2}d\theta.
$$
\n(5.17)

Similarly (as seen in [7]), we have

$$
\mathbb{E}(|x_0(t, n_1) - x_0(t, n_2)|^2) = \frac{4\sigma^2 t}{2\pi} \int_{S^1} G(2t\Lambda(\theta)) \sin^2 \frac{(n_1 - n_2)\theta}{2} d\theta.
$$
 (5.18)

Note that for the Rouse chain model, where L is given by (1.4) , we have (5.11) with

$$
\lambda(n) = \begin{cases}\n-2, & n = 0, \\
1, & n = \pm 1, \\
0, & \text{otherwise,} \n\end{cases}
$$
\n(5.19)

and hence

$$
\Lambda(\theta) = 2 - e^{i\theta} - e^{-i\theta} = 4\sin^2\frac{\theta}{2}.\tag{5.20}
$$

In (2.16) of [4] it was shown that if $\Lambda(\theta)$ is smooth and >0 on $S^1 \setminus \{0\}$ and

$$
\Lambda(\theta) \sim |\theta|^{\rho} \sum_{k \ge 0} a_k |\theta|^k, \quad \theta \to 0,
$$
\n(5.21)

with $a_0 \neq 0$, then

$$
\frac{t}{2\pi} \int_{S^1} G(2t\Lambda(\theta)) d\theta \sim\n\begin{cases}\nCt^{1-1/\rho}, & \rho > 1, \\
C\log t, & \rho = 1, \\
C, & \rho \in (0,1),\n\end{cases}\n\tag{5.22}
$$

as $t\rightarrow\infty$, and consequently, by (5.13),

$$
\mathbb{E}(x_0(t,n)^2) \sim \sigma^2 \times \text{ right side of (5.22), as } t \to \infty. \tag{5.23}
$$

This applies to (5.20) with $\rho = 2$. This large t behavior is to be contrasted with that of the Wiener process:

$$
\mathbb{E}(W_n(t)^2) = t.\tag{5.24}
$$

Because (5.23) is significantly smaller than (5.24) for large t, one says the process $x_0(t,n)$ is subdiffusive. This subdiffusivity result was supplemented in [7] by the following (Propositions 4.1 and 6.1 of [7]), whose proof follows readily from (5.6) and the Lebesgue dominated convergence theorem.

PROPOSITION 5.1. In the general setting of (5.1) – (5.6) , if

$$
\Lambda(\theta) > 0 \quad \text{for} \quad \gamma \text{-}a.e. \quad \theta \in S,\tag{5.25}
$$

then, for each $n \in \mathbb{Z}$,

$$
\mathbb{E}(x_0(t,n)^2) = o(t) \quad \text{as} \quad t \to \infty. \tag{5.26}
$$

Applying Theorem 3.1 immediately leads to the following extension of this result.

THEOREM 5.1. In the general setting of (5.1) – (5.5) , if (5.25) holds, then, for each $n\in\mathbb{Z}$,

$$
\mathbb{E}(x_{\varepsilon}(t,n)^2) = o(t) \quad \text{as} \quad t \to \infty,
$$
\n(5.27)

uniformly in $\varepsilon \in (0, a/\Vert L \Vert)$, for each $a < 1/4$.

Similarly, Theorem 3.1 yields the following extension of the subdiffusivity results for $x_0(t,n)$ discussed above.

THEOREM 5.2. In the setting of (5.11) - (5.12) , if $\Lambda(\theta)$ is smooth and >0 on $S^1 \setminus \{0\}$, and satisfies (5.21), then

$$
\mathbb{E}(x_{\varepsilon}(t,n)^{2}) \leq \begin{cases} Ct^{1-1/\rho}, & \rho > 1, \\ C\log t, & \rho = 1, \\ C, & \rho \in (0,1), \end{cases}
$$
 (5.28)

uniformly for

$$
0 < \varepsilon \le \frac{a}{\|L\|}, \quad a < \frac{1}{4}.\tag{5.29}
$$

The condition (5.29) will be relaxed in §6.

Remark 5.3. In (5.28) we have estimates, as opposed to the asymptotic result in (5.23). To obtain a uniform asymptotic analysis of $\mathbb{E}(x_{\varepsilon}(t,n)^2)$ is an intriguing problem, which we hope to take up in future work.

6. Extension of the scope

In this section, we discard the restriction (1.10) on ε and allow arbitrary $\varepsilon > 0$. As always, L is a negative semidefinite, self adjoint operator on $\ell^2(\mathbb{Z})$, but here we do not require L to be bounded. We will assume that finitely supported elements of $\ell^2(\mathbb{Z})$ belong to the domain of L. As mentioned in the introduction, the operators $A_{\beta}^{\pm}(s)$, given by (1.20), need not be bounded. On the other hand, we have

$$
A_{\beta}(s) = A_{\beta}^{+}(s) - A_{\beta}^{-}(s)
$$

= $(I + 4\varepsilon L)^{-1/2} [e^{s\lambda_{+}(\beta, L)} - e^{s\lambda_{-}(\beta, L)}]$
= $(I + 4\varepsilon L)^{-1/2} [e^{(s\beta/2)(I + 4\varepsilon L)^{1/2}} - e^{-(s\beta/2)(I + 4\varepsilon L)^{1/2}}]e^{-s\beta/2}$
= $s\beta e^{-s\beta/2} H(\frac{s\beta}{2}(I + 4\varepsilon L)^{1/2}),$ (6.1)

where

$$
H(z) = \frac{\sinh z}{z}, \quad H(0) = 1.
$$
 (6.2)

Note that $H(z)$ is an entire function, even in z. There can be some ambiguity in specifying $(I+4\varepsilon L)^{1/2}$, but the fact that $H(z)$ is even in z makes such ambiguity harmless. We have $Spec(I+4\varepsilon L)^{1/2} \subset (0,1]$ if (1.10) holds, while if we merely have $\varepsilon > 0$, we can say

$$
\operatorname{Spec}(I + 4\varepsilon L)^{1/2} \subset [0,1] \cup i\mathbb{R}.\tag{6.3}
$$

Note that for $x, y \in \mathbb{R}$,

$$
H(x) = \frac{\sinh x}{x} \quad \text{and} \quad H(iy) = \frac{\sin y}{y} \tag{6.4}
$$

are real. Hence, for $A_{\beta}(s)$ as in (6.1), we have

$$
A_{\beta}(s)^{*} = A_{\beta}(s). \tag{6.5}
$$

To estimate the operator norm of $A_{\beta}(s)$, note that $|\sin y| \le |y|$ for $y \in \mathbb{R}$, and a calculation gives $H'(x) > 0$ for $x \in [0, \infty)$, so

$$
\sup\left\{|H(z)|:z\in\left[0,\frac{s\beta}{2}\right]\cup i\mathbb{R}\right\}=H\left(\frac{s\beta}{2}\right).
$$
\n(6.6)

Consequently,

$$
||A_{\beta}(s)|| \le s\beta e^{-s\beta/2} \frac{\sinh(s\beta/2)}{s\beta/2}
$$

= 1 - e^{-s\beta}, \t(6.7)

with equality if (as happens in the interesting cases) $0 \in \text{Spec } L$.

Results of §2 imply the processes $x_{\varepsilon}(t) = (x_{\varepsilon}(t,n): n \in \mathbb{Z})$ given by (1.19) are well defined for all $\varepsilon > 0$. If we allow L to be unbounded, we need to note that (6.1) gives a strongly continuous family of operators on $\ell^2(\mathbb{Z})$.

Note that

$$
A'_{\beta}(s) = -\frac{\beta}{2}A_{\beta}(s) + \beta e^{-s\beta/2} \cosh\left(\frac{s\beta}{2}(I + 4\varepsilon L)^{1/2}\right),\tag{6.8}
$$

so, extending results of §4, we have $x_{\varepsilon}(t,n)$ differentiable for all $\varepsilon > 0$, as a function of t with values in $L^2(X,\nu)$.

As in §5, the spectral theorem produces a unitary map

$$
\mathcal{F}: \ell^2(\mathbb{Z}) \longrightarrow L^2(S, \gamma), \tag{6.9}
$$

and a measurable function

$$
\Lambda: S \longrightarrow [0, \infty) \tag{6.10}
$$

(not bounded if L is not bounded), such that

$$
\mathcal{F}e^{tL}y(\theta) = e^{-t\Lambda(\theta)}\mathcal{F}y(\theta).
$$
\n(6.11)

In place of (5.7), we have

$$
\mathbb{E}(x_{\varepsilon}(t,n)^{2})
$$
\n
$$
= \sigma^{2} \int_{0}^{t} \|A_{\beta}(s)\delta_{n}\|_{\ell^{2}}^{2} ds
$$
\n
$$
= \sigma^{2} \int_{0}^{t} \int_{S} (s\beta)^{2} e^{-s\beta} H\left(\frac{s\beta}{2} (1 - 4\varepsilon \Lambda(\theta))^{1/2}\right)^{2} |e_{n}(\theta)|^{2} d\gamma(\theta) ds, \qquad (6.12)
$$

where $e_n = \mathcal{F} \delta_n$. Similarly,

$$
\mathbb{E}(|x_{\varepsilon}(t,n) - x_0(t,n)|^2)
$$

= $\sigma^2 \int_0^t \int\limits_S \left[e^{-s\Lambda(\theta)} - s\beta e^{-s\beta/2} H\left(\frac{s\beta}{2} (1 - 4\varepsilon\Lambda(\theta))^{1/2}\right) \right]^2 |e_n(\theta)|^2 d\gamma(\theta) ds.$ (6.13)

Calculations parallel to those done in §3 establish that

$$
\theta \in S, \ \Lambda(\theta) < \infty, \ s \in [0, \infty), \ \beta = \varepsilon^{-1}
$$
\n
$$
\implies \lim_{\varepsilon \searrow 0} s \beta e^{-s\beta/2} H\left(\frac{s\beta}{2} (1 - 4\varepsilon \Lambda(\theta))^{1/2}\right) = e^{-s\Lambda(\theta)}.
$$
\n
$$
(6.14)
$$

Also, by (6.6) – (6.7) , the integrand in (6.13) is dominated in absolute value by $4|e_n(\theta)|^2$, so the Lebesgue dominated convergence theorem establishes the following.

PROPOSITION 6.1. In the current setting, for each $t \in [0,\infty)$, $n \in \mathbb{Z}$,

$$
\lim_{\varepsilon \searrow 0} \mathbb{E}(|x_{\varepsilon}(t, n) - x_0(t, n)|^2) = 0.
$$
\n(6.15)

This is a partial extension of Theorem 3.1, though it lacks the punch of the estimates (3.41) – (3.42) . We aim to sharpen this up.

To proceed let us fix $M \in [1,\infty)$, take

$$
\varepsilon \in \left(0, \frac{M}{4}\right],\tag{6.16}
$$

and set

$$
S_a = \left\{ \theta \in S : \Lambda(\theta) \le \frac{1}{2M} \right\}, \quad S_b = S \setminus S_a. \tag{6.17}
$$

Thus $\varepsilon \Lambda(\theta) \geq \varepsilon/2M$ on S_b , so

$$
\theta \in S_b \Rightarrow (1 - 4\varepsilon \Lambda(\theta))^{1/2} \begin{cases} \leq 1 - \frac{\varepsilon}{M} & \text{if } 4\varepsilon \Lambda(\theta) \leq 1, \\ \text{is purely imaginary} & \text{if } 4\varepsilon \Lambda(\theta) \geq 1. \end{cases}
$$
(6.18)

Making use of (6.6), with $s\beta/2$ replaced by $(s\beta/2)(1-\varepsilon/M)$, we have

$$
\theta \in S_b \Rightarrow s\beta e^{-s\beta/2} \Big| H\Big(\frac{s\beta}{2} (1 - 4\varepsilon \Lambda(\theta))^{1/2}\Big)\Big|
$$

\n
$$
\leq s\beta e^{-s\beta/2} H\Big(\frac{s\beta}{2} \Big(1 - \frac{\varepsilon}{M}\Big)\Big)
$$

\n
$$
= \frac{2e^{-s\beta/2}}{1 - \varepsilon/M} \Big(e^{(s\beta/2)(1 - \varepsilon/M)} - e^{-(s\beta/2)(1 - \varepsilon/M)}\Big)
$$

\n
$$
= \frac{2}{1 - \varepsilon/M} (e^{-s/2M} - e^{-s\beta(1 - \varepsilon/2M)})
$$

\n
$$
\leq 4e^{-s/2M}, \tag{6.19}
$$

the second identity via $\varepsilon\beta = 1$. In addition,

$$
\theta \in S_b \Longrightarrow e^{-s\Lambda(\theta)} \le e^{-s/2M},\tag{6.20}
$$

so, if $I(s, \varepsilon, \theta)$ denotes the integrand in (6.13), we have

$$
I(s,\varepsilon,\theta) \le 25e^{-s/M}|e_n(\theta)|^2, \quad \forall \theta \in S_b,
$$
\n(6.21)

and hence

$$
\sigma^2 \int_0^t \int_{S_b} I(s,\varepsilon,\theta) d\gamma(\theta) ds \le 25M\sigma^2 \int_{S_b} |e_n(\theta)|^2 d\gamma(\theta)
$$

$$
\le 25M\sigma^2,
$$
 (6.22)

so as $\varepsilon \to 0$ this contribution to (6.13) converges to 0 with uniform bounds, independent of t.

Next, for $\theta \in S_a$, write

$$
s\beta e^{-s\beta/2} H\left(\frac{s\beta}{2} (1 - 4\varepsilon \Lambda(\theta))^{1/2}\right)
$$

=
$$
\frac{e^{-(s\beta/2) + (s\beta/2)(1 - 4\varepsilon \Lambda(\theta))^{1/2}}}{(1 - 4\varepsilon \Lambda(\theta))^{1/2}} - \frac{e^{-(s\beta/2) - (s\beta/2)(1 - 4\varepsilon \Lambda(\theta))^{1/2}}}{(1 - 4\varepsilon \Lambda(\theta))^{1/2}}.
$$
 (6.23)

We have $4\varepsilon\Lambda(\theta) \leq 2\varepsilon/M \leq 1/2$ on S_a , given that ε satisfies (6.16), so the last term in (6.23) is

$$
\leq \sqrt{2}e^{-s\beta/2} \quad \text{on} \quad S_a. \tag{6.24}
$$

Thus, with $I(s, \varepsilon, \theta)$ as in (6.21) – (6.22) , we have

$$
\sigma^2 \int_0^t \int\limits_{S_a} I(s,\varepsilon,\theta) \, d\gamma(\theta) \, ds
$$

$$
= \sigma^2 \int_0^t \int_{S_a} \left[e^{-s\Lambda(\theta)} - \frac{e^{-(s\beta/2) + (s\beta/2)(1 - 4\varepsilon\Lambda(\theta))^{1/2}}}{(1 - 4\varepsilon\Lambda(\theta))^{1/2}} \right]^2 |e_n(\theta)|^2 d\gamma(\theta) ds
$$

+ $R(t,\varepsilon)$, (6.25)

where

$$
|R(t,\varepsilon)| \le C \int_0^t e^{-s\beta} ds \le C\varepsilon,
$$
\n(6.26)

with C independent of t. Next, estimates parallel to (3.18) – (3.40) apply to the main term on the right side of (6.25), given that $\varepsilon \leq M/4$ and $\Lambda(\theta) \leq 1/2M$. We have the main term

$$
\leq C\sigma^2 \mathbb{E}(x_0(t,n)^2)\varepsilon. \tag{6.27}
$$

Putting together these estimates, we have the following.

THEOREM 6.2. For each $M \in [1,\infty)$, we have $C < \infty$ such that, as long as $0 < \varepsilon \le M/4$,

$$
\mathbb{E}(|x_{\varepsilon}(t,n)-x_0(t,n)|^2) \leq C\sigma^2 \mathbb{E}(x_0(t,n))^2 \varepsilon + C\varepsilon + R_b(\varepsilon,t),\tag{6.28}
$$

with

$$
R_b(\varepsilon, t) \le 25M\sigma^2, \quad \forall t \ge 0,
$$
\n(6.29)

and

$$
\lim_{\varepsilon \to 0} R_b(\varepsilon, t) = 0. \tag{6.30}
$$

Using Theorem 6.2 in place of Theorem 3.1, we have the following extension of Theorem 5.2.

THEOREM 6.3. In the setting of (5.11) – (5.12) , if $\Lambda(\theta)$ is smooth and >0 on $S^1 \setminus \{0\}$, and satisfies (5.21), then (5.28) holds, uniformly for $\varepsilon \in (0,K]$, for each $K < \infty$.

REFERENCES

- [1] F. Chung, Spectral Graph Theory, CBMS Reg. Conf. Ser. Math. 92, American Mathematical Society, Providence, RI, 1997.
- [2] G. DaPrato and J. Zabczyk, Stochastic Equations in Infinite Dimensions, Encl. Math. Appl., Cambridge Univ. Press, Cambridge, 45, 1992.
- [3] H. McKean, Stochastic Integrals, Academic Press, New York, 1969.
- [4] S. McKinley, Anomalous diffusion of distinguished particles in bead-spring networks, preprint, 2009, arXiv:0911.4293.
- [5] S. McKinley, L. Yao, and M. G. Forest, Transient anomalous diffusion of tracer particles in soft matter, J. Rheology, 53, 2009.
- [6] M. Reed and B. Simon, Methods of Mathematical Physics, Academic Press, New York, 1, 1980.
- [7] M. Taylor, Gaussian processes associated to infinite bead-spring networks, Commun. Math. Sci., 9, 517–534, 2011.