GAUSSIAN PROCESSES ASSOCIATED TO INFINITE BEAD-SPRING NETWORKS II: BEADS WITH MASS AND THE VANISHING MASS LIMIT*

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Abstract. We construct families of Gaussian processes $x_{\varepsilon}(t,n)$, $t \in [0,\infty)$, $n \in \mathbb{Z}$, modeling a class of infinite networks of stochastically fluctuating, interacting beads, of small mass, proportional to ε . We examine covariances $\mathbb{E}(x_{\varepsilon}(t_1,n_1)x_{\varepsilon}(t_2,n_2))$ and draw conclusions about the subdiffusive nature of these processes, with particular attention to the behavior as $\varepsilon \to 0$. This complements previous work of the author, which in turn was influenced by work of McKinley, Yao, and Forest.

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1. Introduction

In [7], following earlier work of [5] and [4], we studied the behavior of Gaussian processes that can be described as follows. Let $\ell^2(\mathbb{Z})$ denote the space of functions $a:\mathbb{Z}\to\mathbb{C}$ such that $\sum |a(n)|^2 < \infty$ (here \mathbb{Z} denotes the set of integers and \mathbb{C} the set of complex numbers), and let L be a negative semidefinite, self adjoint operator on $\ell^2(\mathbb{Z})$. We assume finitely supported elements of $\ell^2(\mathbb{Z})$ belong to the domain of L, so we can write

$$Ly(n) = \sum_{m \in \mathbb{Z}} \lambda(n, m) y(m).$$
(1.1)

Self adjointness implies $\lambda(n,m) = \overline{\lambda(m,n)}$. We assume

$$\lambda(n,m) \in \mathbb{R}, \quad \text{hence} \quad \lambda(m,n) = \lambda(n,m).$$
 (1.2)

The process x(t) = (x(t,n)) studied in [7] solves the infinite system of stochastic differential equations

$$dx(t,n) = Lx(t,n)dt + \sigma dW_n(t), \quad x(0,n) = 0,$$
(1.3)

for $n \in \mathbb{Z}$, $t \ge 0$. Here W_n are independent, identically distributed Wiener processes. The system (1.3) provides a model for the motion of a polymer, pictured as a network of beads that interact and are also independently randomly jittered, as in Brownian motion. The particular case

$$Ly(n) = y(n-1) - 2y(n) + y(n+1)$$
(1.4)

gives rise to what is called the Rouse chain model; see [5] and [4] for further details and references to the literature.

In [7], the solution to (1.3) was constructed in the form

$$x(t,n) = \sigma \int_0^t \sum_{m \in \mathbb{Z}} h(t-s,n,m) dW_m(s), \qquad (1.5)$$

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where

$$e^{tL}y(n) = \sum_{m} h(t, n, m)y(m).$$
 (1.6)

It was shown that the series in (1.5) converges and defines a Gaussian process, with mean 0. Formulas were derived for $\mathbb{E}(x(t_1, n_1)x(t_2, n_2))$, with special consideration of

$$\mathbb{E}(x(t,n)^2)$$
, and $\mathbb{E}(|x(t,n_1) - x(t,n_2)|^2)$. (1.7)

The analysis of the first expectation in (1.7) recovered results of [5] and [4] on subdiffusivity of x(t,n), and the analysis of the second expectation in (1.7), and also of $\mathbb{E}(x(t,n_1)x(t,n_2))$, provided information on the joint distribution of $x(t,n_1)$ and $x(t,n_2)$.

As pointed out in [5] and [4], the system (1.3) is the $\varepsilon = 0$ case of the second order system

$$\varepsilon x_{\varepsilon}^{\prime\prime}(t,n) + x_{\varepsilon}^{\prime}(t,n) = L x_{\varepsilon}(t,n) + \sigma W_{n}^{\prime}(t), \qquad (1.8)$$

with prime denoting the t-derivative. Here ε is proportional to the mass of each bead. It is reasonable to consider ε to be positive but quite small. Thus it is of interest to study the solution $x_{\varepsilon}(t,n)$ to (1.8), with particular interest in the behavior as $\varepsilon \searrow 0$. This paper addresses that task. We take initial data

$$x_{\varepsilon}(0,n) = 0, \quad x_{\varepsilon}'(0,n) = 0, \quad \forall n \in \mathbb{Z}.$$

$$(1.9)$$

Since (1.8) changes type when ε reaches 0, this is a singular perturbation problem. We first tackle it under an additional condition on L, namely that it be a bounded operator on $\ell^2(\mathbb{Z})$, with operator norm $||L|| < \infty$. This condition holds for (1.4) and for many (arguably, for most) other examples arising in the bead-spring setting. Other examples include graph Laplacians, shown to be bounded in [7], in the case of infinite graphs, following results exposed for finite graphs in [1]. We produce a formula for the solution to (1.8)–(1.9) valid for

$$0 < \varepsilon < \frac{1}{4\|L\|},\tag{1.10}$$

and study its behavior as $\varepsilon \searrow 0$. (In §6 we drop the hypothesis that L be bounded and allow arbitrary $\varepsilon > 0$.)

To see how such a formula arises, let us rewrite (1.5) as

$$x(t) = \sigma \int_0^t e^{(t-s)L} dW(s),$$
 (1.11)

to celebrate how it comes from Duhamel's formula. To obtain an analogue for (1.8), we set $v_{\varepsilon}(t) = x'_{\varepsilon}(t)$, i.e., $v_{\varepsilon}(t,n) = x'_{\varepsilon}(t,n)$, and rewrite (1.8) as a first order system

$$\frac{d}{dt} \begin{pmatrix} x_{\varepsilon} \\ v_{\varepsilon} \end{pmatrix} = X_{\varepsilon} \begin{pmatrix} x_{\varepsilon} \\ v_{\varepsilon} \end{pmatrix} + \begin{pmatrix} 0 \\ \beta \sigma W'(t) \end{pmatrix}, \tag{1.12}$$

where

$$X_{\varepsilon} = \begin{pmatrix} 0 & I \\ \beta L & -\beta I \end{pmatrix}.$$
 (1.13)

Here and below, we set

$$\beta = \frac{1}{\varepsilon}.\tag{1.14}$$

In (1.8) and (1.12), we use the "white noise" formalism W'(t). The system (1.12) is of course a Wiener-Itô stochastic differential equation, which can be written

$$d\binom{x_{\varepsilon}}{v_{\varepsilon}} = X_{\varepsilon}\binom{x_{\varepsilon}}{v_{\varepsilon}}dt + \binom{0}{\beta\sigma \, dW(t)}.$$

Taking into account the initial data (1.9), the Duhamel formula gives

$$\begin{pmatrix} x_{\varepsilon}(t) \\ v_{\varepsilon}(t) \end{pmatrix} = \sigma \int_{0}^{t} e^{(t-s)X_{\varepsilon}} \begin{pmatrix} 0 \\ \beta W'(s) \end{pmatrix} ds$$

$$= \sigma \int_{0}^{t} e^{(t-s)X_{\varepsilon}} \begin{pmatrix} 0 \\ \beta dW(s) \end{pmatrix}.$$

$$(1.15)$$

To compute $e^{sX_{\varepsilon}}$, we note that by the spectral theorem (cf. [6], Chapter 7) we can treat L as a real number and X_{ε} as a real 2×2 matrix, with "eigenvalues"

$$\lambda_{\pm}(\beta,L) = -\frac{\beta}{2}I \pm \frac{\beta}{2}(I + 4\varepsilon L)^{1/2}, \qquad (1.16)$$

and "eigenvectors"

$$\binom{1}{\lambda_{\pm}(\beta,L)}.$$
(1.17)

One then calculates

$$e^{tX_{\varepsilon}} \begin{pmatrix} 0\\1 \end{pmatrix} = \begin{pmatrix} (e^{t\lambda_{+}} - e^{t\lambda_{-}})/(\lambda_{+} - \lambda_{-})\\ (\lambda_{+}e^{t\lambda_{+}} - \lambda_{-}e^{t\lambda_{-}})/(\lambda_{+} - \lambda_{-}) \end{pmatrix}.$$
 (1.18)

Thus (1.15) yields

$$x_{\varepsilon}(t) = \sigma \int_{0}^{t} [A_{\beta}^{+}(t-s) - A_{\beta}^{-}(t-s)] dW(s), \qquad (1.19)$$

where

$$A^{\pm}_{\beta}(s) = (I + 4\varepsilon L)^{-1/2} e^{s\lambda_{\pm}(\beta,L)}, \qquad (1.20)$$

and $\lambda_{\pm}(\beta, L)$, given by (1.16), are bounded, negative semidefinite, self adjoint operators on $\ell^2(\mathbb{Z})$, as long as (1.10) holds. We have the task to show that the right side of (1.19) is a well defined Gaussian process and to investigate its properties, with particular attention to the behavior as $\varepsilon \searrow 0$, i.e., as $\beta \nearrow \infty$.

For use in subsequent sections, in $\S2$ we collect some results on a class of vector stochastic integrals of the form

$$x(t) = \int_0^t A(t-s) \, dW(s), \tag{1.21}$$

where $\{A(s), A(s)^* : s \ge 0\}$ are strongly continuous families of bounded linear operators on $\ell^2(\mathbb{Z})$. Here, $x(t) = (x(t,n), n \in \mathbb{Z})$. We show that for each n, x(t,n) is well defined and is a continuous function of $t \in [0, \infty)$ with values in $L^2(X, \nu)$, where (X, ν) is a naturally constructed probability space (see §2 for details). Also, for each $t \ge 0$, $n \in \mathbb{Z}$, x(t,n) is a Gaussian random variable with mean zero. These results can be established via material in Chapter 4 of [2], but the setting here is more elementary. For the convenience of readers not familiar with infinite dimensional stochastic analysis, we give short, direct demonstrations of the needed formulas, as a consequence of classical work of Paley, Wiener, and Zygmund. Formulas established in §2 include

$$\mathbb{E}(|x(t,n)|^2) = \int_0^t ||A(s)^* \delta_n||_{\ell^2}^2 \, ds, \qquad (1.22)$$

and more generally

674

$$\mathbb{E}(x(t,n_1)\overline{x(t,n_2)}) = \int_0^t (A(s)^* \delta_{n_2}, A(s)^* \delta_{n_1}) ds, \qquad (1.23)$$

where $\{\delta_n : n \in \mathbb{Z}\}\$ is the orthonormal basis of $\ell^2(\mathbb{Z})$ given by $\delta_n(m) = 1$ if m = n, 0 otherwise. If A(s) and $A(s)^*$ commute for all s, one can erase the asterisks in (1.22) and (1.23).

In §3 we apply the results of §2 to $A(s) = A_{\beta}^{\pm}(s)$, given by (1.20), and construct

$$x_{\varepsilon}(t,n) = x_{\varepsilon}^{+}(t,n) - x_{\varepsilon}^{-}(t,n), \qquad (1.24)$$

when (1.10) holds. Here $x_{\varepsilon}^{\pm}(t,n)$ is the *n*th component of

$$x_{\varepsilon}^{\pm}(t) = \sigma \int_0^t A_{\beta}^{\pm}(t-s) \, dW(s),$$

with A_{β}^{\pm} as in (1.20). We compare $x_{\varepsilon}(t,n)$ to the solution to (1.3), given by (1.5), which we now denote $x_0(t,n)$. We show that

$$\mathbb{E}(|x_{\varepsilon}^{+}(t,n) - x_{0}(t,n)|^{2}) \leq C \varepsilon \mathbb{E}(x_{0}(t,n)^{2}), \qquad (1.25)$$

and

$$\mathbb{E}(x_{\varepsilon}^{-}(t,n)^{2}) \leq C\sigma^{2}(1-e^{-t/\varepsilon})\varepsilon, \qquad (1.26)$$

provided $0 < \varepsilon \leq a/\|L\|$, with a < 1/4; see Theorem 3.1. These estimates imply that whenever $x_0(t,n)$ is subdiffusive, i.e.,

$$\frac{1}{t}\mathbb{E}(x_0(t,n)^2) \longrightarrow 0 \quad \text{as} \quad t \nearrow \infty, \tag{1.27}$$

the processes $x_{\varepsilon}(t,n)$ are uniformly subdiffusive, for ε in such an interval.

In §4 we note that the processes $x_{\varepsilon}^{\pm}(t,n)$ are differentiable (as functions of t with values in $L^2(X,\nu)$), for ε satisfying (1.10), and study

$$v_{\varepsilon}^{\pm}(t,n) = \frac{d}{dt} x_{\varepsilon}^{\pm}(t,n).$$
(1.28)

At least one of these must blow up as $\varepsilon \searrow 0$, since $x_0(t,n)$ is not differentiable; as it turns out, $v_{\varepsilon}^-(t,n)$ blows up. We show that

$$\mathbb{E}(v_{\varepsilon}^{+}(t,n)^{2}) \leq C \mathbb{E}(x_{0}(t,n)^{2}), \qquad (1.29)$$

but

$$\mathbb{E}(v_{\varepsilon}^{-}(t,n)^{2}) \ge \frac{\sigma^{2}}{4\varepsilon}(1 - e^{-t/\varepsilon}).$$
(1.30)

In §5 we convert formulas for $\mathbb{E}(x_{\varepsilon}(t,n)^2)$ into integral formulas, arising from a spectral representation of L, and examine the asymptotic behavior as $t \nearrow \infty$, including more precise versions of the subdiffusivity result (1.27) and their counterparts for $\mathbb{E}(x_{\varepsilon}(t,n)^2)$; see Theorems 5.1–5.2.

Results of §§3–5 use the hypothesis (1.10). We obtain estimates valid uniformly for $0 < \varepsilon \le a/\|L\|$, given a < 1/4. In §6 we extend the scope of our investigation, in two ways. First, we replace (1.10) by

$$0 < \varepsilon < \infty. \tag{1.31}$$

Second, we remove the hypothesis that L be bounded. In this more general setting, frequently $-1/4\varepsilon$ belongs to the spectrum of L and represents a transition from overdamping to underdamping in the system (1.8). The operators $A_{\beta}^{\pm}(s)$ in (1.20) are then not bounded, and the processes $x_{\varepsilon}^{\pm}(t,n)$ do not exist. However,

$$A_{\beta}(s) = A_{\beta}^{+}(s) - A_{\beta}^{-}(s) \tag{1.32}$$

is bounded. In fact, from (1.20) we obtain

$$A_{\beta}(s) = s\beta e^{-s\beta/2} H\left(\frac{s\beta}{2} (I + 4\varepsilon L)^{1/2}\right), \qquad (1.33)$$

where H is the entire holomorphic, even function on \mathbb{C} given by

$$H(z) = \frac{\sinh z}{z}, \quad H(0) = 1.$$
 (1.34)

Using this, we show that the processes $x_{\varepsilon}(t,n)$ exist. We obtain formulas for $\mathbb{E}(x_{\varepsilon}(t,n)^2)$, etc., extending those obtained earlier for ε satisfying (1.10). Making use of these results, we extend the scope of results of §5. Our main results in this section are given in Theorems 6.1–6.2.

2. A class of vector stochastic integrals

In this section we provide some useful formulas for vector stochastic integrals of the form

$$x(t) = \int_0^t A(t-s) \, dW(s), \tag{2.1}$$

where, for each $s \ge 0$,

$$A(s): \ell^2(\mathbb{Z}) \longrightarrow \ell^2(\mathbb{Z}) \tag{2.2}$$

is a bounded linear operator. For simplicity we assume

A(s) and $A(s)^*$ are strongly continuous in $s \in [0, \infty)$, (2.3)

though the calculations below will make it clear that we can relax this hypothesis. Written out more fully, (2.1) takes the form

$$x(t,n) = \int_{0}^{t} \sum_{m \in \mathbb{Z}} a(t-s,n,m) dW_{m}(s), \qquad (2.4)$$

where, for $y \in \ell^2(\mathbb{Z})$,

$$A(s)y(n) = \sum_{m} a(s,n,m)y(m).$$

$$(2.5)$$

The operators arising in (1.19) are self adjoint (for $0 < \varepsilon < 1/4 ||L||$) and reality preserving, but we do not need these properties for the development here. Consequently, the processes (2.4) might be complex valued. Note the adjoint $A(s)^*$ of A(s) satisfies

$$A(s)^* y(m) = \sum_m a^*(s, n, m) y(m), \quad a^*(s, n, m) = \overline{a(s, m, n)},$$
(2.6)

and that

$$a(s,n,m) = A(s)\delta_m(n), \quad a^*(s,n,m) = A(s)^*\delta_m(n),$$
 (2.7)

where $\delta_n \in \ell^2(\mathbb{Z})$ is given by

$$\delta_n(m) = \delta_{n,m} = \begin{cases} 1 & \text{if } m = n, \\ 0 & \text{otherwise.} \end{cases}$$
(2.8)

As stated in the introduction, $\{W_n : n \in \mathbb{Z}\}$ is a collection of independent, identically distributed Wiener processes. In more detail, let B(t) be the Wiener process (Brownian motion), which is a continuous family $B(t) \in L^2(\Omega, \mu)$, where Ω is path space and μ is Wiener measure. Then set $\Omega_n = \Omega$, $\mu_n = \mu$, for $n \in \mathbb{Z}$, and take the product space (with product measure)

$$(X,\nu) = \prod_{n \in \mathbb{Z}} (\Omega_n, \mu_n).$$
(2.9)

We obtain (2.4) as

$$x(t,n) = \lim_{K \to \infty} \sum_{m=-K}^{K} \xi_m(t,n),$$
 (2.10)

where

$$\xi_m(t,n) = \int_0^t a(t-s,n,m) \, dW_m(s). \tag{2.11}$$

Our first task is to establish convergence in $L^2(X,\nu)$ of the right side of (2.10). Note that

$$m \neq m' \Longrightarrow \xi_m(t,n) \perp \xi_{m'}(t,n) \quad \text{in} \quad L^2(X,\nu),$$
(2.12)

so it suffices to bound $\sum_m \mathbb{E}(|\xi_m(t,n)|^2)$. To get this, note that

$$\mathbb{E}(|\xi_m(t,n)|^2) = \int_0^t a(s,n,m)a^*(s,m,n)\,ds,$$
(2.13)

which is the classical Paley-Wiener-Zygmund identity (cf. [3], §2.1). Hence

$$\sum_{m} \mathbb{E}(|\xi_{m}(t,n)|^{2}) = \sum_{m} \int_{0}^{t} a(s,n,m) a^{*}(s,m,n) \, ds$$

$$=\sum_{m} \int_{0}^{t} a(s,n,m) A(s)^{*} \delta_{n}(m) ds$$

= $\int_{0}^{t} A(s) A(s)^{*} \delta_{n}(n) ds$
= $\int_{0}^{t} (A(s) A(s)^{*} \delta_{n}, \delta_{n}) ds$
= $\int_{0}^{t} \|A(s)^{*} \delta_{n}\|_{\ell^{2}}^{2} ds.$ (2.14)

Here $\delta_n \in \ell^2(\mathbb{Z})$ is given by (2.8). Thus we have convergence in (2.10), and

$$\mathbb{E}(|x(t,n)|^2) = \int_0^t ||A(s)^* \delta_n||_{\ell^2}^2 \, ds.$$
(2.15)

The nature of the convergence implies that for each $n \in \mathbb{Z}$, $t \ge 0$, x(t,n) is a Gaussian random variable on (X, ν) with mean 0.

We next aim to show that, under the hypotheses in (2.3), x(t,n) is a continuous function of $t \in [0,\infty)$ with values in $L^2(X,\nu)$, for each n. In preparation for this, we note that

$$\mathbb{E}(x(t_1,n)\overline{x(t_2,n)}) = \sum_k \int_0^{t_1 \wedge t_2} a(t_1 - s, n, k) a^*(t_2 - s, k, n) \, ds.$$
(2.16)

Here $t_1 \wedge t_2 = \min(t_1, t_2)$. We have

$$\sum_{k} a(t_1 - s, n, k) a^*(t_2 - s, k, n) = \sum_{k} a(t_1 - s, n, k) A(t_2 - s)^* \delta_n(k)$$

= $A(t_1 - a) A(t_2 - s)^* \delta_n(n)$
= $(A(t_2 - s)^* \delta_n, A(t_1 - s)^* \delta_n),$ (2.17)

 \mathbf{SO}

$$\mathbb{E}(x(t_1,n),\overline{x(t_2,n)}) = \int_0^{t_1 \wedge t_2} (A(t_2-s)^* \delta_n, A(t_1-s)^* \delta_n) \, ds.$$
(2.18)

Now

$$\mathbb{E}(|x(t_1,n) - x(t_2,n)|^2) = \mathbb{E}(x(t_1,n)^2) + \mathbb{E}(x(t_2,n)^2) - 2\operatorname{Re}\mathbb{E}(x(t_1,n)\overline{x(t_2,n)}),$$
(2.19)

so (2.18) gives (say if $0 \le t_1 \le t_2$)

$$\mathbb{E}(|x(t_1,n)-x(t_2,n)|^2) = \int_{t_1}^{t_2} ||A(t_2-s)^*\delta_n||_{\ell^2}^2 ds + \int_0^{t_1\wedge t_2} \left\{ (A(t_1-s)^*\delta_n, A(t_1-s)^*\delta_n) + (A(t_2-s)^*\delta_n, A(t_2-s)^*\delta_n) - 2\operatorname{Re}(A(t_2-s)^*\delta_n, A(t_1-s)^*\delta_n) \right\} ds. \quad (2.20)$$

The first integral on the right side of (2.20) is $\leq C|t_1 - t_2|$. We can write the second integral as

$$\operatorname{Re} \int_{0}^{t_{1} \wedge t_{2}} \left\{ \left([A(t_{1}-s)^{*} - A(t_{2}-s)^{*}]\delta_{n}, A(t_{1}-s)^{*}\delta_{n} \right) + \left(A(t_{2}-s)^{*}\delta_{n}, [A(t_{2}-s)^{*} - A(t_{1}-s)^{*}]\delta_{n} \right) \right\} ds.$$

$$(2.21)$$

Given (2.3), the fact that (2.21) tends to 0 as $t_1 \rightarrow t_2$ follows readily from the Lebesgue dominated convergence theorem. Let us summarize what we have established.

PROPOSITION 2.1. Under the hypotheses (2.2)-(2.3), the formula (2.4) gives for each $n \in \mathbb{Z}$ a well defined x(t,n), a continuous function of $t \in [0,\infty)$ with values in $L^2(X,\nu)$, for each t,n a Gaussian random variable with mean 0, satisfying the identities (2.15) and (2.18).

We complement (2.15) with the following computation, derived similarly.

$$\mathbb{E}(x(t,n_1)\overline{x(t,n_2)}) = \sum_k \int_0^t a(t-s,n_1,k)\overline{a(t-s,n_2,k)} ds$$

$$= \sum_k \int_0^t a(t-s,n_1,k)a^*(t-s,k,n_2) ds.$$
(2.22)

Parallel to (2.17), we then get

$$\mathbb{E}(x(t,n_1)\overline{x(t,n_2)}) = \int_0^t (A(t-s)^* \delta_{n_2}, A(t-s)^* \delta_{n_1}) ds$$

= $\int_0^t (A(s)^* \delta_{n_2}, A(s)^* \delta_{n_1}) ds.$ (2.23)

Combining (2.15) and (2.23), we have

$$\mathbb{E}(|x(t,n_1) - x(t,n_2)|^2) = \mathbb{E}(|x(t,n_1)|^2) + \mathbb{E}(|x(t,n_2)|^2) - 2\operatorname{Re}\mathbb{E}(x(t,n_1)\overline{x(t,n_2)}) = \int_0^t \left\{ \|A(s)^*\delta_{n_1}\|_{\ell^2}^2 + \|A(s)^*\delta_{n_2}\|_{\ell^2}^2 - 2\operatorname{Re}(A(s)^*\delta_{n_2}, A(s)^*\delta_{n_1}) \right\} ds = \int_0^t \|A(s)^*(\delta_{n_1} - \delta_{n_2})\|_{\ell^2}^2 ds.$$
(2.24)

We now give a condition under which the components x(t,n) of the process (2.1) are differentiable, as functions of t with values in $L^2(X,\nu)$. Let us add to (2.3) the hypothesis

$$A'(s)$$
 and $A'(s)^*$ are strongly continuous in $s \in [0, \infty)$. (2.25)

Then, as in the scalar case, Wiener's integration by parts formula holds for (2.1):

$$x(t) = \int_0^t A'(t-s)W(s)\,ds + A(0)W(t). \tag{2.26}$$

We have the following.

PROPOSITION 2.2. In the setting of Proposition 2.1, if also (2.25) holds and A(0) = 0, then x(t,n) is differentiable for each $n \in \mathbb{Z}$, and x'(t,n) is a continuous function of $t \in [0,\infty)$ with values in $L^2(X,\nu)$.

Proof. Let us temporarily assume that (2.25) also holds for A''(s). Then we differentiate (2.26) and get (provided A(0) = 0)

$$x'(t) = \int_0^t A''(t-s)W(s)\,ds + A'(0)W(t). \tag{2.27}$$

Applying (2.26) with A replaced by A' then gives

$$x'(t) = \int_0^t A'(t-s) \, dW(s). \tag{2.28}$$

A mollification and approximation argument gives (2.28) without the additional assumption on A''.

Returning to the computations (2.13)–(2.23), note that if A(s) is self adjoint for all s, all the asterisks can be removed, and if these operators are reality preserving, all the overlines can be removed. Furthermore,

$$A(s)^* A(s) = A(s)A(s)^* \implies ||A(s)^* \delta_n||_{\ell^2} = ||A(s)\delta_n||_{\ell^2}, \text{ and} (A(s)^* \delta_{n_2}, A(s)^* \delta_{n_1}) = (A(s)\delta_{n_2}, A(s)\delta_{n_1}),$$
(2.29)

so we have the following.

PROPOSITION 2.3. In the setting of Proposition 2.1, if also A(s) is normal for all $s \ge 0$, then

$$\mathbb{E}(|x(t,n)|^2) = \int_0^t ||A(s)\delta_n||_{\ell^2}^2 ds, \qquad (2.30)$$

and more generally

$$\mathbb{E}(x(t,n_1)\overline{x(t,n_2)}) = \int_0^t (A(s)\delta_{n_2}, A(s)\delta_{n_1}) \, ds.$$
(2.31)

3. The processes $x_{\varepsilon}(t,n)$

Results of §2 on the vector stochastic integral $\int_0^t A(t-s) dW(s)$ apply to (1.19) with

$$A(s) = \sigma [A_{\beta}^{+}(s) - A_{\beta}^{-}(s)],$$

$$A_{\beta}^{\pm}(s) = (I + 4\varepsilon L)^{-1/2} e^{s\lambda_{\pm}(\beta,L)},$$

$$\lambda_{\pm}(\beta,L) = -\frac{\beta}{2} I \pm \frac{\beta}{2} (I + 4\varepsilon L)^{1/2}.$$
(3.1)

In the current setting, L is a bounded, reality preserving, negative semidefinite, self adjoint operator on $\ell^2(\mathbb{Z})$, $0 < \varepsilon < 1/4 ||L||$, and $\beta = 1/\varepsilon$. Hence $\lambda_{\pm}(\beta, L)$ are negative semidefinite, self adjoint operators on $\ell^2(\mathbb{Z})$. Thus, for each such ε , $x_{\varepsilon}(t) = (x_{\varepsilon}(t,n), n \in \mathbb{Z})$ has the property that, for each $n \in \mathbb{Z}$, $x_{\varepsilon}(t,n)$ is a continuous function of $t \in [0, \infty)$ with values in $L^2(X, \nu)$, and for each $t \ge 0$ is a real valued Gaussian random variable with mean 0. For further analysis, it is convenient (using (1.19)) to write

$$x_{\varepsilon}(t,n) = x_{\varepsilon}^{+}(t,n) - x_{\varepsilon}^{-}(t,n), \qquad (3.2)$$

where

$$x_{\varepsilon}^{\pm}(t) = \sigma \int_0^t A_{\beta}^{\pm}(t-s) \, dW(s). \tag{3.3}$$

The formula (2.15) gives

$$\mathbb{E}(x_{\varepsilon}^{\pm}(t,n)^{2}) = \sigma^{2} \int_{0}^{t} \|A_{\beta}^{\pm}(s)\delta_{n}\|_{\ell^{2}}^{2} ds.$$
(3.4)

Note that Spec $\lambda_{-}(\beta, L) \subset (-\infty, -\beta/2]$, so we have the operator norm estimate

$$\|A_{\beta}^{-}(s)\| \leq \|(I+4\varepsilon L)^{-1/2}\|e^{-s\beta/2}, \tag{3.5}$$

and we get

$$\mathbb{E}(x_{\varepsilon}^{-}(t,n)^{2}) \leq C\sigma^{2}\varepsilon(1-e^{-\beta t}), \quad 0 < \varepsilon \leq \frac{a}{\|L\|}, \quad a < \frac{1}{4}, \tag{3.6}$$

with C independent of $t \in [0,\infty)$.

In order to analyze $x_{\varepsilon}^{+}(t,n)$, note that, as long as (1.10) holds,

$$(I+4\varepsilon L)^{1/2} = I + 2\varepsilon L\Phi(4\varepsilon L), \qquad (3.7)$$

with $\Phi(\lambda)$ given by

$$(1+\lambda)^{1/2} = 1 + \frac{1}{2}\lambda - \frac{1}{8}\lambda^2 + \cdots$$

= $1 + \frac{1}{2}\lambda \left(1 - \frac{1}{4}\lambda + \cdots\right)$
= $1 + \frac{1}{2}\lambda \Phi(\lambda).$ (3.8)

Note that $\Phi(\lambda)$ is holomorphic on $\{\lambda \in \mathbb{C} : |\lambda| < 1\}$ and

 $\Phi(0) = 1, \quad \Phi(\lambda) > 0 \ \text{ for } \ \lambda \in (-1, 1).$ (3.9)

Hence

$$\lambda_{+}(\beta, L) = L\Phi(4\varepsilon L), \qquad (3.10)$$

 \mathbf{SO}

$$A_{\beta}^{+}(s) = (I + 4\varepsilon L)^{-1/2} e^{sL\Phi(4\varepsilon L)}, \qquad (3.11)$$

and

$$x_{\varepsilon}^{+}(t) = \sigma \int_{0}^{t} (I + 4\varepsilon L)^{-1/2} e^{(t-s)L\Phi(4\varepsilon L)} dW(s).$$
(3.12)

Hence

$$\mathbb{E}(x_{\varepsilon}^{+}(t,n)^{2}) = \sigma^{2} \int_{0}^{t} \|(I+4\varepsilon L)^{-1/2} e^{sL\Phi(4\varepsilon L)} \delta_{n}\|_{\ell^{2}}^{2} ds$$

$$= \sigma^{2} \int_{0}^{t} (e^{2sL\Phi(4\varepsilon L)} (I+4\varepsilon L)^{-1} \delta_{n}, \delta_{n}) ds.$$
(3.13)

If we set

$$G(\lambda) = \int_0^1 e^{-s\lambda} ds = \begin{cases} \frac{1-e^{-\lambda}}{\lambda} & , \ \lambda > 0, \\ 1 & , \ \lambda = 0, \end{cases}$$
(3.14)

we can write (3.13) as

$$\mathbb{E}(x_{\varepsilon}^{+}(t,n)^{2}) = \sigma^{2}t \big(G(-2tL\Phi(4\varepsilon L))(I+4\varepsilon L)^{-1}\delta_{n},\delta_{n} \big).$$
(3.15)

In $\S5$ we will investigate large t behavior of this.

At this point, it is natural to compare $x_{\varepsilon}^+(t)$ with the solution $x_0(t)$ to (1.3), given by (1.11), i.e.,

$$x_0(t) = \sigma \int_0^t e^{(t-s)L} dW(s).$$
(3.16)

Note that, parallel to (3.13)-(3.15),

$$\mathbb{E}(x_0(t,n)^2) = \sigma^2 \int_0^t \|e^{sL}\delta_n\|_{\ell^2}^2 ds$$

= $\sigma^2 t \big(G(-2tL)\delta_n, \delta_n \big).$ (3.17)

Applying (2.15) to the difference of (3.12) and (3.16) gives

$$\mathbb{E}(|x_{\varepsilon}^{+}(t,n) - x_{0}(t,n)|^{2})$$

$$= \sigma^{2} \int_{0}^{t} \|[A_{\beta}^{+}(t-s) - e^{(t-s)L}]\delta_{n}\|_{\ell^{2}}^{2} ds$$

$$= \sigma^{2} \int_{0}^{t} \|[(I+4\varepsilon L)^{-1/2}e^{sL\Phi(4\varepsilon L)} - e^{sL}]\delta_{n}\|_{\ell^{2}}^{2} ds.$$
(3.18)

Using $(a+b)^2 \le 2a^2 + 2b^2$, we can write

$$\mathbb{E}(|x_{\varepsilon}^{+}(t,n) - x_{0}(t,n)|^{2}) \leq 2\sigma^{2}(A_{\varepsilon}(t,n) + B_{\varepsilon}(t,n)), \qquad (3.19)$$

where

$$A_{\varepsilon}(t,n) = \int_{0}^{t} \|(I+4\varepsilon L)^{-1/2} [e^{sL\Phi(4\varepsilon L)} - e^{sL}] \delta_{n} \|_{\ell^{2}}^{2} ds,$$

$$B_{\varepsilon}(t,n) = \int_{0}^{t} \|[(I+4\varepsilon L)^{-1/2} - I] e^{sL} \delta_{n} \|_{\ell^{2}}^{2} ds.$$
(3.20)

Noting that

$$\|(I+4\varepsilon L)^{-1/2} - I\| \le C\varepsilon, \quad \text{for} \quad 0 < \varepsilon \le \frac{a}{\|L\|}, \quad a < \frac{1}{4}, \tag{3.21}$$

and comparing (3.17), we have

$$\sigma^2 B_{\varepsilon}(t,n) \le C \varepsilon \mathbb{E}(x_0(t,n)^2). \tag{3.22}$$

We also have

$$A_{\varepsilon}(t,n) \le C\widetilde{A}_{\varepsilon}(t,n), \qquad (3.23)$$

where

$$\widetilde{A}_{\varepsilon}(t,n) = \int_0^t \|[e^{sL\Phi(4\varepsilon L)} - e^{sL}]\delta_n\|_{\ell^2}^2 ds.$$
(3.24)

To proceed, recall from (3.8) that

$$(1+\lambda)^{1/2} = 1 + \frac{1}{2}\lambda\Phi(\lambda),$$
 (3.25)

where $\Phi(\lambda)$ is given by

$$\Phi(\lambda) = 1 - \frac{1}{4}\lambda + \dots = 1 - \frac{\lambda}{4}\psi(\lambda), \quad \psi(0) = 1,$$
(3.26)

with $\psi(\lambda)$ holomorphic in $\{\lambda \in \mathbb{C} : |\lambda| < 1\}$, real and positive for $\lambda \in (-1,1)$. The positivity can be seen from the concavity of $(1+\lambda)^{1/2}$, which implies $(1+\lambda)^{1/2} \le 1+\lambda/2$ on (-1,1), hence $\Phi(\lambda) \le 1$ on [0,1) and ≥ 1 on (-1,0]. Hence

$$e^{sL\Phi(4\varepsilon L)} - e^{sL} = \left(e^{-s\varepsilon L^2\psi(4\varepsilon L)} - I\right)e^{sL},\tag{3.27}$$

and we have

$$\widetilde{A}_{\varepsilon}(t,n) = \int_0^t \| (I - e^{-s\varepsilon L^2 \psi(4\varepsilon L)}) e^{sL} \delta_n \|_{\ell^2}^2 ds, \qquad (3.28)$$

which gives

$$\sigma^2 \widetilde{A}_{\varepsilon}(t,n) \leq \sup_{0 < s < t} \|I - e^{-s\varepsilon L^2 \psi(4\varepsilon L)}\|^2 \mathbb{E}(x_0(t,n)^2).$$
(3.29)

If we take $a \in (0, 1/4)$ and set

$$\alpha = \sup_{0 < \varepsilon \le a/\|L\|} \|L^2 \psi(4\varepsilon L)\|, \qquad (3.30)$$

then, since $L^2\psi(4\varepsilon L)$ is positive semidefinite, we have

$$\sigma^2 \widetilde{A}_{\varepsilon}(t,n) \le (1 - e^{-\alpha \varepsilon t})^2 \mathbb{E}(x_0(t,n)^2), \tag{3.31}$$

provided

$$0 < \varepsilon \le \frac{a}{\|L\|}, \quad a < \frac{1}{4}. \tag{3.32}$$

The factor in front of $\mathbb{E}(x_0(t,n)^2)$ in (3.31) is $O(\varepsilon)$ for t in each bounded interval in $[0,\infty)$, but one loses uniformity as $t \nearrow \infty$. In fact, (3.31) is not optimal. We proceed to derive a stronger estimate. Writing

$$\left(I - e^{-s\varepsilon L^2\psi(4\varepsilon L)}\right)e^{sL} = \left(e^{sL/2} - e^{sL/2 - s\varepsilon L^2\psi(4\varepsilon L)}\right)e^{sL/2},\tag{3.33}$$

we have

$$\widetilde{A}_{\varepsilon}(t,n) \le A_{\varepsilon}^{\#}(t)^2 \int_0^t \|e^{sL/2} \delta_n\|_{\ell^2}^2 ds, \qquad (3.34)$$

with

$$A_{\varepsilon}^{\#}(t) = \sup_{0 < s < t} \|e^{sL/2} - e^{sL/2 - s\varepsilon L^{2}\psi(4\varepsilon L)}\|$$

$$\leq \sup_{0 < s < t, 0 \le \Lambda \le \|L\|} |e^{-s\Lambda/2} - e^{-s\Lambda/2 - s\varepsilon\Lambda^{2}\psi(-4\varepsilon\Lambda)}|, \qquad (3.35)$$

the latter inequality by the spectral theorem. Now, over the range $0 \le \Lambda \le ||L||$,

$$\varphi = \Lambda \psi(-4\varepsilon \Lambda) \Longrightarrow 0 \le \varphi \le \Lambda_0, \tag{3.36}$$

as long as (3.32) holds, where $\Lambda_0 = \|L\| \sup_{(-1,1)} \psi(\lambda)$, and we obtain

$$A_{\varepsilon}^{\#}(t) \leq \sup_{0 < s < t, \Lambda \ge 0, 0 \le \varphi \le \Lambda_0} \left| e^{-s\Lambda/2} - e^{-s\Lambda/2 - s\varepsilon\Lambda\varphi} \right|.$$
(3.37)

Taking $s\Lambda \mapsto \Lambda$, we get

$$\begin{aligned} A_{\varepsilon}^{\#}(t) &\leq \sup_{\Lambda \geq 0, 0 \leq \varphi \leq \Lambda_{0}} \left| e^{-\Lambda/2} - e^{-\Lambda/2 - \varepsilon \varphi \Lambda} \right| \\ &\leq \sup_{\Lambda \geq 0, 0 \leq \varphi \leq \Lambda_{0}} \varepsilon \varphi \Lambda e^{-\Lambda/2} \\ &\leq \varepsilon \Lambda_{0}, \end{aligned}$$
(3.38)

since $\sup \Lambda e^{-\Lambda/2} = 2/e < 1$. Note that this estimate is independent of t. Meanwhile,

$$\sigma^{2} \int_{0}^{t} \|e^{sL/2} \delta_{n}\|_{\ell^{2}}^{2} ds = 2\sigma^{2} \int_{0}^{t/2} \|e^{sL} \delta_{n}\|_{\ell^{2}}^{2} ds$$
$$= 2\mathbb{E} \left(x_{0} \left(\frac{t}{2}, n \right)^{2} \right)$$
$$\leq 2\mathbb{E} (x_{0}(t, n)^{2}), \qquad (3.39)$$

so (3.34) and (3.38) yield

$$\sigma^2 \widetilde{A}_{\varepsilon}(t,n) \le 2\Lambda_0 \varepsilon \mathbb{E}(x_0(t,n)^2).$$
(3.40)

Let us collect the main results established above.

THEOREM 3.1. As long as (3.32) holds, the formulas (3.2)–(3.3) give, for each $n \in \mathbb{Z}$, a mean zero Gaussian process $t \mapsto x_{\varepsilon}(t,n) = x_{\varepsilon}^+(t,n) - x_{\varepsilon}^-(t,n)$, a continuous function of $t \in [0,\infty)$ with values in $L^2(X,\nu)$. Furthermore, there exist $C, \alpha \in (0,\infty)$ such that when (3.32) holds and $x_0(t,n)$ is given by (3.16), then for all $n \in \mathbb{Z}$, $t \ge 0$,

$$\mathbb{E}(|x_{\varepsilon}^{+}(t,n) - x_{0}(t,n)|^{2}) \leq C\varepsilon \mathbb{E}(x_{0}(t,n)^{2}), \qquad (3.41)$$

and

$$\mathbb{E}(x_{\varepsilon}^{-}(t,n)^{2}) \leq C\sigma^{2}(1-e^{-t/\varepsilon})\varepsilon.$$
(3.42)

We record formulas for the covariance of $x_{\varepsilon}^{\pm}(t,n_1)$ and $x_{\varepsilon}^{\pm}(t,n_2)$. By (2.23), we have (with coherent choice of signs)

$$\mathbb{E}(x_{\varepsilon}^{\pm}(t,n_{1})x_{\varepsilon}^{\pm}(t,n_{2}))$$

$$=\sigma^{2}\int_{0}^{t}(A_{\beta}^{\pm}(s)\delta_{n_{1}},A_{\beta}^{\pm}(s)\delta_{n_{2}})$$

$$=\sigma^{2}\int_{0}^{t}((I+4\varepsilon L)^{-1}e^{2s\lambda_{\pm}(\beta,L)}\delta_{n_{1}},\delta_{n_{2}})ds.$$
(3.43)

In particular, using (3.10),

$$\mathbb{E}(x_{\varepsilon}^{+}(t,n_{1})x_{\varepsilon}^{+}(t,n_{2}))$$

$$=\sigma^{2}\int_{0}^{t}((I+4\varepsilon L)^{-1}e^{2sL\Phi(4\varepsilon L)}\delta_{n_{1}},\delta_{n_{2}})ds$$

$$=\sigma^{2}t((I+4\varepsilon L)^{-1}G(-2tL\Phi(4\varepsilon L))\delta_{n_{1}},\delta_{n_{2}}).$$
(3.44)

4. The processes $v_{\varepsilon}(t,n)$

From (3.1) we see that $A_{\beta}^{\pm'}(0) = (I + 4\varepsilon L)^{-1/2}$ for $0 < \varepsilon < 1/4 ||L||$, so A(0) = 0 and, by Proposition 2.2, $x_{\varepsilon}(t,n)$ is differentiable, as a function of t, with values in $L^2(X,\nu)$, for each $n \in \mathbb{Z}$. By (2.28),

$$x_{\varepsilon}'(t,n) = v_{\varepsilon}(t,n) = v_{\varepsilon}^{+}(t,n) - v_{\varepsilon}^{-}(t,n), \qquad (4.1)$$

with

$$v_{\varepsilon}^{\pm}(t) = \sigma \int_{0}^{t} V_{\beta}^{\pm}(t-s) dW(s), \qquad (4.2)$$

where

$$V_{\beta}^{\pm}(s) = \frac{d}{ds} A_{\beta}^{\pm}(s)$$

$$= (I + 4\varepsilon L)^{-1/2} \lambda_{\pm}(\beta, L) e^{s\lambda_{\pm}(\beta, L)}.$$
(4.3)

As before, $\beta = 1/\varepsilon$. We will compute square expectations and verify, as one should expect, that $\mathbb{E}(v_{\varepsilon}(t,n)^2) \to \infty$ as $\varepsilon \searrow 0$. In fact, we separately examine $\mathbb{E}(v_{\varepsilon}^+(t,n)^2)$ and $\mathbb{E}(v_{\varepsilon}^-(t,n)^2)$, and see that only the latter blows up as $\varepsilon \searrow 0$.

To begin, we have

$$\mathbb{E}(v_{\varepsilon}^{\pm}(t,n)^{2})$$

$$=\sigma^{2}\int_{0}^{t} \|V_{\beta}^{\pm}(s)\delta_{n}\|_{\ell^{2}}^{2} ds$$

$$=\sigma^{2}\int_{0}^{t} \|(I+4\varepsilon L)^{-1/2}\lambda_{\pm}(\beta,L)e^{s\lambda_{\pm}(\beta,L)}\delta_{n}\|_{\ell^{2}}^{2} ds.$$
(4.4)

Recalling from (3.10) that $\lambda_{+}(\beta, L) = L\Phi(4\varepsilon L)$, we have, for

$$0 < \varepsilon \le \frac{a}{\|L\|}, \quad a < \frac{1}{4}, \tag{4.5}$$

that

$$\mathbb{E}(v_{\varepsilon}^{+}(t,n)^{2})$$

$$= \sigma^{2} \int_{0}^{t} \|L\Phi(4\varepsilon L)(I+4\varepsilon L)^{-1/2} e^{sL\Phi(4\varepsilon L)} \delta_{n}\|_{\ell^{2}}^{2} ds$$

$$\leq C \mathbb{E}(x_{\varepsilon}^{+}(t,n)^{2})$$

$$\leq C' \mathbb{E}(x_{0}(t,n)^{2}), \qquad (4.6)$$

the first inequality by (3.13), given the operator norm bound $||L\Phi(4\varepsilon L)|| \leq C$, and the second by (3.41).

To proceed, we have

$$\mathbb{E}(v_{\varepsilon}^{-}(t,n)^{2})$$

$$=\sigma^{2}\int_{0}^{t}((I+4\varepsilon L)^{-1}\lambda_{-}(\beta,L)^{2}e^{2s\lambda_{-}(\beta,L)}\delta_{n},\delta_{n})ds$$

$$=-\frac{\sigma^{2}}{2}((I+4\varepsilon L)^{-1}\lambda_{-}(\beta,L)(I-e^{2t\lambda_{-}(\beta,L)})\delta_{n},\delta_{n}).$$
(4.7)

Now, as long as (4.5) holds, we have, via the spectral theorem,

$$Spec(I + 4\varepsilon L)^{-1} \subset [1, \infty),$$

$$Spec\lambda_{-}(\beta, L) \subset [-\beta, -\beta/2],$$

$$Spec(I - e^{2t\lambda_{-}(\beta, L)}) \subset [1 - e^{-\beta t}, 1],$$

(4.8)

and hence

$$\operatorname{Spec} - (I + 4\varepsilon L)^{-1} \lambda_{-}(\beta, L) (I - e^{2t\lambda_{-}(\beta, L)}) \subset \left[\frac{\beta}{2} (1 - e^{-t\beta}), \infty\right).$$
(4.9)

The variational characterization of the bottom of the spectrum for a positive definite, self adjoint operator, applied to the last inner product in (4.7), then gives

$$\mathbb{E}(v_{\varepsilon}^{-}(t,n)^{2}) \ge \frac{\sigma^{2}}{4\varepsilon}(1 - e^{-t/\varepsilon}), \qquad (4.10)$$

as long as (4.5) holds. The right side of (4.10) clearly blows up as $\varepsilon \searrow 0$, for each t > 0.

5. Spectral representation, asymptotics, and subdiffusivity

Let $L: \ell^2(\mathbb{Z}) \to \ell^2(\mathbb{Z})$ be a bounded, negative, self adjoint operator, as described in §1. The spectral theorem (cf. [6], Theorem VII.3) implies there is a measure space (S, γ) , a unitary map

$$\mathcal{F}: \ell^2(\mathbb{Z}) \longrightarrow L^2(S, \gamma), \tag{5.1}$$

and a function

$$\Lambda \in L^{\infty}(S,\gamma), \quad \Lambda \ge 0, \quad \|\Lambda\|_{L^{\infty}} = \|L\|, \tag{5.2}$$

such that for each $y \in \ell^2(\mathbb{Z}), t \ge 0$,

$$\mathcal{F}Ly(\theta) = -\Lambda(\theta)\mathcal{F}y(\theta), \quad \theta \in S.$$
(5.3)

Consequently,

$$\mathcal{F}e^{tL}y(\theta) = e^{-t\Lambda(\theta)}\mathcal{F}y(\theta),$$

$$\mathcal{F}\Phi(4\varepsilon L)y(\theta) = \Phi(-4\varepsilon\Lambda(\theta))\mathcal{F}y(\theta),$$
(5.4)

etc. The orthonormal basis $\{\delta_n\}$ of $\ell^2(\mathbb{Z})$ gives rise to the orthonormal basis $\{e_n\}$ of $L^2(S,\gamma)$,

$$e_n = \mathcal{F}\delta_n. \tag{5.5}$$

Using these ingredients, we can rewrite the formula (3.17) for the square expectation of $x_0(t,n)$ as

$$\mathbb{E}(x_0(t,n)^2) = \sigma^2 \int_S \int_0^t e^{-2s\Lambda(\theta)} |e_n(\theta)|^2 ds d\gamma(\theta)$$

= $\sigma^2 t \int_S G(2t\Lambda(\theta)) |e_n(\theta)|^2 d\gamma(\theta).$ (5.6)

Similarly, (3.13)–(3.15) yield

$$\mathbb{E}(x_{\varepsilon}^{+}(t,n)^{2}) = \sigma^{2} \int_{S} \int_{0}^{t} (1 - 4\varepsilon \Lambda(\theta))^{-1} e^{-2s\Lambda_{\varepsilon}(\theta)} |e_{n}(\theta)|^{2} ds d\gamma(\theta)$$
$$= \sigma^{2} t \int_{S} (1 - 4\varepsilon \Lambda(\theta))^{-1} G(2t\Lambda_{\varepsilon}(\theta)) |e_{n}(\theta)|^{2} d\gamma(\theta),$$
(5.7)

 for

$$0 < \varepsilon < \frac{1}{4\|L\|}.\tag{5.8}$$

Here,

$$\Lambda_{\varepsilon}(\theta) = \Lambda(\theta) \Phi(-4\varepsilon \Lambda(\theta)) = \Lambda(\theta) (1 + \varepsilon \Lambda(\theta) \psi(-4\varepsilon \Lambda(\theta))),$$
(5.9)

with Φ as in (3.7)–(3.10) and ψ as in (3.26). Note that, as long as (5.8) holds, $\psi(-4\varepsilon\Lambda(\theta)) \ge 0$. More generally, by (3.44),

$$\mathbb{E}(x_{\varepsilon}^{+}(t,n_{1})x_{\varepsilon}^{+}(t,n_{2})) = \sigma^{2} \int_{S} \int_{0}^{t} (1-4\varepsilon\Lambda(\theta))^{-1} e^{-2s\Lambda_{\varepsilon}(\theta)} e_{n_{1}}(\theta) \overline{e_{n_{2}}(\theta)} \, ds \, d\gamma(\theta) = \sigma^{2} t \int_{S} (1-4\varepsilon\Lambda(\theta))^{-1} G(2t\Lambda_{\varepsilon}(\theta)) e_{n_{1}}(\theta) \overline{e_{n_{2}}(\theta)} \, d\gamma(\theta).$$
(5.10)

Let us specialize to the case that L is of convolution type:

$$Ly(n) = \sum_{m} \lambda(n-m)y(m).$$
(5.11)

A special case is given in (1.4), for the Rouse chain model. The convolution case was also emphasized in [5] and [4]. In this case, we can take

$$S = S^{1} = \mathbb{R}/(2\pi\mathbb{Z}), \quad d\gamma(\theta) = d\theta/2\pi,$$

$$\mathcal{F}y(\theta) = \hat{y}(\theta) = \sum_{n} y(n)e^{in\theta}, \quad e_{n}(\theta) = e^{in\theta}, \quad \Lambda(\theta) = -\hat{\lambda}(\theta).$$
(5.12)

In such a case, (5.6)–(5.10) become

$$\mathbb{E}(x_0(t,n)^2) = \frac{\sigma^2 t}{2\pi} \int_{S^1} G(2t\Lambda(\theta)) d\theta, \qquad (5.13)$$

$$\mathbb{E}(x_{\varepsilon}^{+}(t,n)^{2}) = \frac{\sigma^{2}t}{2\pi} \int_{S^{1}} (1 - 4\varepsilon \Lambda(\theta))^{-1} G(2t\Lambda_{\varepsilon}(\theta)) d\theta, \qquad (5.14)$$

and

$$\mathbb{E}(x_{\varepsilon}^{+}(t,n_{1})x_{\varepsilon}^{+}(t,n_{2})) = \frac{\sigma^{2}t}{2\pi} \int_{S^{1}} (1 - 4\varepsilon\Lambda(\theta))^{-1} G(2t\Lambda_{\varepsilon}(\theta)) e^{i(n_{1} - n_{2})\theta} d\theta.$$
(5.15)

Note that the reality condition (1.2) implies

$$\lambda(-n) = \lambda(n), \text{ hence } \Lambda(-\theta) = \Lambda(\theta).$$
 (5.16)

Taking this into account, a short computation yields

$$\mathbb{E}(|x_{\varepsilon}^{+}(t,n_{1})-x_{\varepsilon}^{+}(t,n_{2})|^{2}) = \mathbb{E}(x_{\varepsilon}^{+}(t,n_{1})^{2}) + \mathbb{E}(x_{\varepsilon}^{+}(t,n_{2})^{2}) - 2\mathbb{E}(x_{\varepsilon}^{+}(t,n_{1})x_{\varepsilon}^{+}(t,n_{2})) \\ = \frac{4\sigma^{2}t}{2\pi} \int_{S^{1}} (1-4\varepsilon\Lambda(\theta))^{-1}G(2t\Lambda_{\varepsilon}(\theta))\sin^{2}\frac{(n_{1}-n_{2})\theta}{2}d\theta.$$
(5.17)

Similarly (as seen in [7]), we have

$$\mathbb{E}(|x_0(t,n_1) - x_0(t,n_2)|^2) = \frac{4\sigma^2 t}{2\pi} \int_{S^1} G(2t\Lambda(\theta)) \sin^2 \frac{(n_1 - n_2)\theta}{2} d\theta.$$
(5.18)

Note that for the Rouse chain model, where L is given by (1.4), we have (5.11) with

$$\lambda(n) = \begin{cases} -2, & n = 0, \\ 1, & n = \pm 1, \\ 0, & \text{otherwise}, \end{cases}$$
(5.19)

and hence

$$\Lambda(\theta) = 2 - e^{i\theta} - e^{-i\theta} = 4\sin^2\frac{\theta}{2}.$$
(5.20)

In (2.16) of [4] it was shown that if $\Lambda(\theta)$ is smooth and >0 on $S^1 \setminus \{0\}$ and

$$\Lambda(\theta) \sim |\theta|^{\rho} \sum_{k \ge 0} a_k |\theta|^k, \quad \theta \to 0,$$
(5.21)

with $a_0 \neq 0$, then

$$\frac{t}{2\pi} \int_{S^1} G(2t\Lambda(\theta)) \, d\theta \sim \begin{cases} Ct^{1-1/\rho}, & \rho > 1, \\ C\log t, & \rho = 1, \\ C, & \rho \in (0,1), \end{cases}$$
(5.22)

as $t \to \infty$, and consequently, by (5.13),

$$\mathbb{E}(x_0(t,n)^2) \sim \sigma^2 \times \text{ right side of } (5.22), \text{ as } t \to \infty.$$
(5.23)

This applies to (5.20) with $\rho = 2$. This large t behavior is to be contrasted with that of the Wiener process:

$$\mathbb{E}(W_n(t)^2) = t. \tag{5.24}$$

Because (5.23) is significantly smaller than (5.24) for large t, one says the process $x_0(t,n)$ is subdiffusive. This subdiffusivity result was supplemented in [7] by the following (Propositions 4.1 and 6.1 of [7]), whose proof follows readily from (5.6) and the Lebesgue dominated convergence theorem.

PROPOSITION 5.1. In the general setting of (5.1)-(5.6), if

$$\Lambda(\theta) > 0 \quad for \quad \gamma \text{-}a.e. \quad \theta \in S, \tag{5.25}$$

then, for each $n \in \mathbb{Z}$,

$$\mathbb{E}(x_0(t,n)^2) = o(t) \quad as \quad t \to \infty.$$
(5.26)

Applying Theorem 3.1 immediately leads to the following extension of this result.

THEOREM 5.1. In the general setting of (5.1)–(5.5), if (5.25) holds, then, for each $n \in \mathbb{Z}$,

$$\mathbb{E}(x_{\varepsilon}(t,n)^2) = o(t) \quad as \quad t \to \infty, \tag{5.27}$$

uniformly in $\varepsilon \in (0, a/||L||]$, for each a < 1/4.

Similarly, Theorem 3.1 yields the following extension of the subdiffusivity results for $x_0(t,n)$ discussed above.

THEOREM 5.2. In the setting of (5.11)–(5.12), if $\Lambda(\theta)$ is smooth and >0 on $S^1 \setminus \{0\}$, and satisfies (5.21), then

$$\mathbb{E}(x_{\varepsilon}(t,n)^2) \leq \begin{cases} Ct^{1-1/\rho}, & \rho > 1, \\ C\log t, & \rho = 1, \\ C, & \rho \in (0,1), \end{cases}$$
(5.28)

uniformly for

$$0 < \varepsilon \le \frac{a}{\|L\|}, \quad a < \frac{1}{4}. \tag{5.29}$$

The condition (5.29) will be relaxed in §6.

REMARK 5.3. In (5.28) we have estimates, as opposed to the asymptotic result in (5.23). To obtain a uniform asymptotic analysis of $\mathbb{E}(x_{\varepsilon}(t,n)^2)$ is an intriguing problem, which we hope to take up in future work.

6. Extension of the scope

In this section, we discard the restriction (1.10) on ε and allow arbitrary $\varepsilon > 0$. As always, L is a negative semidefinite, self adjoint operator on $\ell^2(\mathbb{Z})$, but here we do not require L to be bounded. We will assume that finitely supported elements of $\ell^2(\mathbb{Z})$ belong to the domain of L. As mentioned in the introduction, the operators $A_{\beta}^{\pm}(s)$, given by (1.20), need not be bounded. On the other hand, we have

$$\begin{aligned} A_{\beta}(s) &= A_{\beta}^{+}(s) - A_{\beta}^{-}(s) \\ &= (I + 4\varepsilon L)^{-1/2} \left[e^{s\lambda_{+}(\beta,L)} - e^{s\lambda_{-}(\beta,L)} \right] \\ &= (I + 4\varepsilon L)^{-1/2} \left[e^{(s\beta/2)(I + 4\varepsilon L)^{1/2}} - e^{-(s\beta/2)(I + 4\varepsilon L)^{1/2}} \right] e^{-s\beta/2} \\ &= s\beta e^{-s\beta/2} H \left(\frac{s\beta}{2} (I + 4\varepsilon L)^{1/2} \right), \end{aligned}$$
(6.1)

where

$$H(z) = \frac{\sinh z}{z}, \quad H(0) = 1.$$
 (6.2)

Note that H(z) is an entire function, even in z. There can be some ambiguity in specifying $(I+4\varepsilon L)^{1/2}$, but the fact that H(z) is even in z makes such ambiguity harmless. We have $\operatorname{Spec}(I+4\varepsilon L)^{1/2} \subset (0,1]$ if (1.10) holds, while if we merely have $\varepsilon > 0$, we can say

$$\operatorname{Spec}(I + 4\varepsilon L)^{1/2} \subset [0, 1] \cup i\mathbb{R}.$$
(6.3)

Note that for $x, y \in \mathbb{R}$,

$$H(x) = \frac{\sinh x}{x}$$
 and $H(iy) = \frac{\sin y}{y}$ (6.4)

are real. Hence, for $A_{\beta}(s)$ as in (6.1), we have

$$A_{\beta}(s)^* = A_{\beta}(s). \tag{6.5}$$

To estimate the operator norm of $A_{\beta}(s)$, note that $|\sin y| \leq |y|$ for $y \in \mathbb{R}$, and a calculation gives H'(x) > 0 for $x \in [0, \infty)$, so

$$\sup\left\{|H(z)|:z\in\left[0,\frac{s\beta}{2}\right]\cup i\mathbb{R}\right\}=H\left(\frac{s\beta}{2}\right).$$
(6.6)

Consequently,

$$\|A_{\beta}(s)\| \leq s\beta e^{-s\beta/2} \frac{\sinh(s\beta/2)}{s\beta/2}$$

$$= 1 - e^{-s\beta}, \qquad (6.7)$$

with equality if (as happens in the interesting cases) $0 \in \operatorname{Spec} L$.

Results of §2 imply the processes $x_{\varepsilon}(t) = (x_{\varepsilon}(t,n): n \in \mathbb{Z})$ given by (1.19) are well defined for all $\varepsilon > 0$. If we allow L to be unbounded, we need to note that (6.1) gives a strongly continuous family of operators on $\ell^2(\mathbb{Z})$.

Note that

$$A_{\beta}'(s) = -\frac{\beta}{2}A_{\beta}(s) + \beta e^{-s\beta/2} \cosh\left(\frac{s\beta}{2}(I+4\varepsilon L)^{1/2}\right),\tag{6.8}$$

so, extending results of §4, we have $x_{\varepsilon}(t,n)$ differentiable for all $\varepsilon > 0$, as a function of t with values in $L^2(X,\nu)$.

As in $\S5$, the spectral theorem produces a unitary map

$$\mathcal{F}: \ell^2(\mathbb{Z}) \longrightarrow L^2(S, \gamma), \tag{6.9}$$

and a measurable function

$$\Lambda: S \longrightarrow [0,\infty) \tag{6.10}$$

(not bounded if L is not bounded), such that

$$\mathcal{F}e^{tL}y(\theta) = e^{-t\Lambda(\theta)}\mathcal{F}y(\theta). \tag{6.11}$$

In place of (5.7), we have

$$\mathbb{E}(x_{\varepsilon}(t,n)^{2})$$

$$=\sigma^{2}\int_{0}^{t}\|A_{\beta}(s)\delta_{n}\|_{\ell^{2}}^{2}ds$$

$$=\sigma^{2}\int_{0}^{t}\int_{S}(s\beta)^{2}e^{-s\beta}H\left(\frac{s\beta}{2}(1-4\varepsilon\Lambda(\theta))^{1/2}\right)^{2}|e_{n}(\theta)|^{2}d\gamma(\theta)ds, \qquad (6.12)$$

where $e_n = \mathcal{F}\delta_n$. Similarly,

$$\mathbb{E}(|x_{\varepsilon}(t,n) - x_{0}(t,n)|^{2}) = \sigma^{2} \int_{0}^{t} \int_{S} \left[e^{-s\Lambda(\theta)} - s\beta e^{-s\beta/2} H\left(\frac{s\beta}{2} (1 - 4\varepsilon\Lambda(\theta))^{1/2}\right) \right]^{2} |e_{n}(\theta)|^{2} d\gamma(\theta) ds.$$
(6.13)

Calculations parallel to those done in §3 establish that

$$\theta \in S, \ \Lambda(\theta) < \infty, \ s \in [0, \infty), \ \beta = \varepsilon^{-1}$$

$$\Longrightarrow \lim_{\varepsilon \searrow 0} s\beta e^{-s\beta/2} H\left(\frac{s\beta}{2}(1 - 4\varepsilon\Lambda(\theta))^{1/2}\right) = e^{-s\Lambda(\theta)}.$$

$$(6.14)$$

Also, by (6.6)–(6.7), the integrand in (6.13) is dominated in absolute value by $4|e_n(\theta)|^2$, so the Lebesgue dominated convergence theorem establishes the following.

PROPOSITION 6.1. In the current setting, for each $t \in [0,\infty)$, $n \in \mathbb{Z}$,

$$\lim_{\varepsilon \searrow 0} \mathbb{E}(|x_{\varepsilon}(t,n) - x_0(t,n)|^2) = 0.$$
(6.15)

This is a partial extension of Theorem 3.1, though it lacks the punch of the estimates (3.41)–(3.42). We aim to sharpen this up.

To proceed let us fix $M \in [1,\infty)$, take

$$\varepsilon \in \left(0, \frac{M}{4}\right],\tag{6.16}$$

and set

$$S_a = \left\{ \theta \in S : \Lambda(\theta) \le \frac{1}{2M} \right\}, \quad S_b = S \setminus S_a.$$
(6.17)

Thus $\varepsilon \Lambda(\theta) \ge \varepsilon/2M$ on S_b , so

$$\theta \in S_b \Rightarrow (1 - 4\varepsilon \Lambda(\theta))^{1/2} \begin{cases} \leq 1 - \frac{\varepsilon}{M} & \text{if } 4\varepsilon \Lambda(\theta) \leq 1, \\ \text{is purely imaginary} & \text{if } 4\varepsilon \Lambda(\theta) \geq 1. \end{cases}$$
(6.18)

Making use of (6.6), with $s\beta/2$ replaced by $(s\beta/2)(1-\varepsilon/M)$, we have

$$\begin{aligned} \theta \in S_b \Rightarrow \quad s\beta e^{-s\beta/2} \left| H\left(\frac{s\beta}{2}(1-4\varepsilon\Lambda(\theta))^{1/2}\right) \right| \\ &\leq s\beta e^{-s\beta/2} H\left(\frac{s\beta}{2}\left(1-\frac{\varepsilon}{M}\right)\right) \\ &= \frac{2e^{-s\beta/2}}{1-\varepsilon/M} \left(e^{(s\beta/2)(1-\varepsilon/M)} - e^{-(s\beta/2)(1-\varepsilon/M)}\right) \\ &= \frac{2}{1-\varepsilon/M} (e^{-s/2M} - e^{-s\beta(1-\varepsilon/2M)}) \\ &\leq 4e^{-s/2M}, \end{aligned}$$
(6.19)

the second identity via $\varepsilon\beta = 1$. In addition,

$$\theta \in S_b \Longrightarrow e^{-s\Lambda(\theta)} \le e^{-s/2M}, \tag{6.20}$$

so, if $I(s,\varepsilon,\theta)$ denotes the integrand in (6.13), we have

$$I(s,\varepsilon,\theta) \le 25e^{-s/M} |e_n(\theta)|^2, \quad \forall \theta \in S_b,$$
(6.21)

and hence

$$\sigma^{2} \int_{0}^{t} \int_{S_{b}} I(s,\varepsilon,\theta) d\gamma(\theta) ds \leq 25M\sigma^{2} \int_{S_{b}} |e_{n}(\theta)|^{2} d\gamma(\theta)$$

$$\leq 25M\sigma^{2},$$
(6.22)

so as $\varepsilon \to 0$ this contribution to (6.13) converges to 0 with uniform bounds, independent of t.

Next, for $\theta \in S_a$, write

$$s\beta e^{-s\beta/2} H\left(\frac{s\beta}{2}(1-4\varepsilon\Lambda(\theta))^{1/2}\right)$$

=
$$\frac{e^{-(s\beta/2)+(s\beta/2)(1-4\varepsilon\Lambda(\theta))^{1/2}}}{(1-4\varepsilon\Lambda(\theta))^{1/2}} - \frac{e^{-(s\beta/2)-(s\beta/2)(1-4\varepsilon\Lambda(\theta))^{1/2}}}{(1-4\varepsilon\Lambda(\theta))^{1/2}}.$$
(6.23)

We have $4\varepsilon\Lambda(\theta) \le 2\varepsilon/M \le 1/2$ on S_a , given that ε satisfies (6.16), so the last term in (6.23) is

$$\leq \sqrt{2}e^{-s\beta/2}$$
 on S_a . (6.24)

Thus, with $I(s,\varepsilon,\theta)$ as in (6.21)–(6.22), we have

$$\sigma^2 \int_0^t \int\limits_{S_a} I(s,\varepsilon,\theta) \, d\gamma(\theta) \, ds$$

$$=\sigma^{2} \int_{0}^{t} \int_{S_{a}} \left[e^{-s\Lambda(\theta)} - \frac{e^{-(s\beta/2) + (s\beta/2)(1 - 4\varepsilon\Lambda(\theta))^{1/2}}}{(1 - 4\varepsilon\Lambda(\theta))^{1/2}} \right]^{2} |e_{n}(\theta)|^{2} d\gamma(\theta) ds$$
$$+ R(t,\varepsilon), \tag{6.25}$$

where

$$|R(t,\varepsilon)| \le C \int_0^t e^{-s\beta} \, ds \le C\varepsilon, \tag{6.26}$$

with C independent of t. Next, estimates parallel to (3.18)–(3.40) apply to the main term on the right side of (6.25), given that $\varepsilon \leq M/4$ and $\Lambda(\theta) \leq 1/2M$. We have the main term

$$\leq C\sigma^2 \mathbb{E}(x_0(t,n)^2)\varepsilon. \tag{6.27}$$

Putting together these estimates, we have the following.

THEOREM 6.2. For each $M \in [1, \infty)$, we have $C < \infty$ such that, as long as $0 < \varepsilon \leq M/4$,

$$\mathbb{E}(|x_{\varepsilon}(t,n) - x_0(t,n)|^2) \le C\sigma^2 \mathbb{E}(x_0(t,n)^2)\varepsilon + C\varepsilon + R_b(\varepsilon,t),$$
(6.28)

with

$$R_b(\varepsilon, t) \le 25M\sigma^2, \quad \forall t \ge 0, \tag{6.29}$$

and

$$\lim_{\varepsilon \to 0} R_b(\varepsilon, t) = 0. \tag{6.30}$$

Using Theorem 6.2 in place of Theorem 3.1, we have the following extension of Theorem 5.2.

THEOREM 6.3. In the setting of (5.11)-(5.12), if $\Lambda(\theta)$ is smooth and >0 on $S^1 \setminus \{0\}$, and satisfies (5.21), then (5.28) holds, uniformly for $\varepsilon \in (0, K]$, for each $K < \infty$.

REFERENCES

- F. Chung, Spectral Graph Theory, CBMS Reg. Conf. Ser. Math. 92, American Mathematical Society, Providence, RI, 1997.
- [2] G. DaPrato and J. Zabczyk, Stochastic Equations in Infinite Dimensions, Encl. Math. Appl., Cambridge Univ. Press, Cambridge, 45, 1992.
- [3] H. McKean, Stochastic Integrals, Academic Press, New York, 1969.
- S. McKinley, Anomalous diffusion of distinguished particles in bead-spring networks, preprint, 2009, arXiv:0911.4293.
- [5] S. McKinley, L. Yao, and M. G. Forest, Transient anomalous diffusion of tracer particles in soft matter, J. Rheology, 53, 2009.
- [6] M. Reed and B. Simon, Methods of Mathematical Physics, Academic Press, New York, 1, 1980.
- [7] M. Taylor, Gaussian processes associated to infinite bead-spring networks, Commun. Math. Sci., 9, 517–534, 2011.