

# A sharp sup+inf inequality for a nonlinear elliptic equation in $\mathbb{R}^2$

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## 1. Introduction.

In this paper we are concerned with the equation

$$(1.1) \quad \Delta u + K(x)e^u = 0 \text{ in } \mathbb{R}^2,$$

where  $\Omega$  is a domain in  $\mathbb{R}^2$  and  $K$  satisfies

$$(1.2) \quad a \leq K(x) \leq b$$

for some positive constants  $a$  and  $b$ . This equation appears in the problem of finding a metric conformal to the standard Euclidean metric in  $\mathbb{R}^2$  such that  $\frac{1}{2}K(x)$  is the Gaussian curvature of the new metric. For a solution  $u$  of (1.1), the total curvature is defined by

$$(1.3) \quad \alpha = \int_{\mathbb{R}^2} K(x)e^{u(x)} dx$$

Suppose  $K(x)$  satisfies (1.2), one interesting question is how to estimate the total curvature  $\alpha$  in terms of the constants  $a$  and  $b$ . Our first result concerning the equation (1.1) is

**Theorem 1.1.** *Assume  $K$  satisfies (1.2) and  $u$  is a solution of (1.1). Then*

- (1) *the total curvature  $\alpha \geq 4\pi(1 + \sqrt{\frac{a}{b}})$ ; and,*
- (2) *if  $\alpha = 4\pi(1 + \sqrt{\frac{a}{b}})$ , then, after a translation,*

$$(1.4) \quad K(x) = \begin{cases} b & \text{if } |x| \leq r_0, \\ a & \text{if } |x| > r_0, \end{cases}$$

and

(1.5)

$$u(x) = \begin{cases} 2 \log \frac{2\sqrt{\frac{a}{b}}}{1+(\frac{|x|}{r_0})^2} - 2 \log r_0 + \log 2 & \text{if } |x| \leq r_0 \\ 2 \log \left( \frac{2\sqrt{\frac{a}{b}}(\frac{|x|}{r_0})\sqrt{\frac{a}{b}-1}}{1+(\frac{|x|}{r_0})^{2\sqrt{\frac{a}{b}}}} \right) - 2 \log r_0 + \log 2 & \text{if } |x| > r_0 \end{cases}$$

hold for almost everywhere  $x$  and for some  $r_0 > 0$ .

If we consider radial solutions only, then we can estimate the upper bound of the total curvature provided that the total curvature is finite.

**Theorem 1.2.** *Assume that  $K$  satisfies (1.2) and both  $K(x)$  and  $u(x)$  are radially symmetric with respect to 0. Suppose the total curvature of  $u$  is finite, then*

$$(1) \alpha \leq 4\pi \left( 1 + \sqrt{\frac{b}{a}} \right); \text{ and}$$

$$(2) \text{ if } \alpha = 4\pi \left( 1 + \sqrt{\frac{b}{a}} \right), \text{ then}$$

$$(1.6) \quad K(x) = \begin{cases} a & \text{if } |x| \leq r_0 \\ b & \text{if } |x| > r_0, \end{cases}$$

and

(1.7)

$$u(x) = \begin{cases} 2 \log \left( \frac{2\sqrt{\frac{b}{a}}}{1+(\frac{|x|}{r_0})^2} \right) - 2 \log r_0 + \log 2 & \text{if } |x| \leq r_0, \\ 2 \log \left( \frac{2\sqrt{\frac{b}{a}}(\frac{|x|}{r_0})\sqrt{\frac{b}{a}-1}}{1+(\frac{|x|}{r_0})^{2\sqrt{\frac{b}{a}}}} \right) - 2 \log r_0 + \log 2 & \text{if } |x| \geq r_0 \end{cases}$$

for some  $r_0 > 0$ .

As an application of Theorem 1.1, we can derive an interior estimate for a solution  $u$  of

$$(1.8) \quad \Delta u + K(x)e^{u(x)} = 0 \quad \text{in } \Omega$$

where  $\Omega$  is a bounded domain in  $\mathbb{R}^2$  and  $K$  satisfies (1.2). Let  $A$  be a compact set of  $\Omega$ . It was proved in [BM] that for any solution  $u$  of (1.8), the inequality

$$\sup_A u \leq C(a, b, A, \Omega, \inf_\Omega u)$$

holds. It was conjectured in [BM] that the dependence of  $C$  in terms of  $\inf_\Omega u$  is linear. In [S], the conjecture was proved to be true, that is, there exist constants  $C_1 \leq 1$  and  $C_2$  such that

$$(1.9) \quad C_1 \sup_A u + \inf_\Omega u \leq C_2$$

where  $C_2 = C_2(a, b, A, \Omega)$  and  $C_1$  depends on  $\frac{a}{b}$  only.

From (1.5), we know that for (1.9) to hold true,  $C_1$  must be less than or equal to  $\sqrt{\frac{a}{b}}$ . The main result of this paper is

**Theorem 1.3.** *Assume that  $K$  satisfies (1.2) and  $u$  is a solution of (1.8). Then, for any compact subset  $A$  of  $\Omega$ , there exists a constant  $C = C(a, b, A, \Omega)$  such that the inequality,*

$$(1.10) \quad \sqrt{\frac{a}{b}} \sup_A u + \inf_\Omega u \leq C,$$

holds. More generally, if we assume that there are  $\frac{1}{2} \geq \rho > 0$ ,  $\sigma \geq 1$  and  $B \geq 0$  such that for  $|x - y| \leq \rho$ ,

$$(1.11) \quad \frac{K(y)}{K(x)} \leq \sigma + \frac{B}{|\log|x - y||}$$

holds, then any solution  $u$  of (1.8) satisfies

$$(1.12) \quad \sqrt{\frac{1}{\sigma}} \sup_A u + \inf_\Omega u \leq C,$$

where  $C$  depends on  $a, b, A, \Omega, \rho, \sigma$ , and  $B$ .

Obviously, if  $K$  is Hölder continuous, that is,

$$(1.13) \quad |K(x) - K(y)| \leq B|x - y|^\beta$$

for  $x, y \in \bar{\Omega}$  and for some constants  $B, 0 < \beta < 1$ , then  $K$  satisfies (1.11) with  $\sigma = 1$ . Hence, as a corollary of Theorem 1.3, we answer a question asked in [BLS].

**Corollary 1.4.** *Suppose that  $K$  satisfies (1.2) and (1.13). Then any solution  $u$  of (1.8) satisfies*

$$\sup_A u + \inf_\Omega u \leq C$$

where  $C = C(a, b, A, \Omega, B, \beta)$ .

We will first prove both Theorem 1.1 and Theorem 1.2 in Section 2. The proof of Theorem 1.3 will be given in Section 3.

## 2. Proofs of Theorem 1.1 and Theorem 1.2.

In this section, we begin with a proof of Theorem 1.1.

*Proof of Theorem 1.1.* If  $\alpha = +\infty$ , the theorem holds trivially. Hence we may assume  $\alpha < +\infty$ . By a result in [CL] (See Theorem 1.2 in [CL]), for any  $\epsilon > 0$ , there  $R(\epsilon) > 0$  such that

$$(2.1) \quad \frac{-\alpha}{2\pi} \log |x| - c \leq u(x) \leq \left( \frac{-\alpha}{2\pi} + \epsilon \right) \log |x|$$

holds for  $|x| \geq R(\epsilon)$ , where  $c$  is a constant independent of  $\epsilon$ . In particular, we have

$$(2.2) \quad \lim_{|x| \rightarrow +\infty} u(x) = -\infty,$$

and

$$(2.3) \quad \alpha > 4\pi.$$

For each  $t \in \mathbf{R}$ , let  $\Omega_t = \{x \in \mathbb{R}^2 \mid u(x) > t\}$ . By (2.2),  $|\Omega_t|$  is finite, where  $|E|$  denotes the area of a measurable set  $E$  in  $\mathbb{R}^2$ . Let  $u^*$  be the Schwartz symmetrization of  $u$ , that is,  $u^*(x) = u^*(|x|)$  is nonincreasing in  $|x|$  and  $\Omega_t^* = \{x \mid u^*(x) > t\}$  is the ball  $B_r(0)$  with the radius  $r = (\frac{1}{\pi} |\Omega_t|)^{\frac{1}{2}}$ . Since  $u$  is locally Lipschitz, we have  $u^*$  is locally Lipschitz also. We also note that, since  $u$  satisfies equation (1.1),  $|\{x \mid u(x) = t\}| = 0$  for any  $t \in \mathbb{R}$ . Therefore,  $|\Omega_t|$  is strictly increasing in  $t$  and then  $u^*(x)$  is strictly decreasing in  $|x|$ .

Set

$$(2.4) \quad F(r) = \int_{\Omega_{u^*(r)}} K(x) e^{u(x)} dx,$$

and

$$(2.5) \quad \bar{K}(r) = \frac{F'(r)}{2\pi r} e^{-u^*(r)},$$

where  $F'(r)$  denotes the derivative of  $F$  with respect to  $r$ . Since  $|\Omega_{u^*(r)}| = \pi r^2$ , we have, for  $r \geq s$ ,

$$|\Omega_{u^*(r)} \setminus \Omega_{u^*(s)}| = \pi(r^2 - s^2) = \pi(r+s)(r-s).$$

Hence

$$F(r) - F(s) = \int_{\Omega_{u^*(r)} \setminus \Omega_{u^*(s)}} K(x)e^{u(x)} dx \leq c(r-s)$$

for some constant  $c$  depending on  $r$ . Thus  $F'(r)$  exists for almost everywhere  $r$  and  $\bar{K}(r)$  is defined for almost everywhere  $r$ . For such  $r$ , we have

(2.6)

$$\begin{aligned} 2\pi r a e^{u^*(r)} &\leq \lim_{h \rightarrow 0^+} a e^{u^*(r+h)} \frac{\pi(r+h)^2 - \pi r^2}{h} \\ &\leq \lim_{h \rightarrow 0^+} \frac{1}{h} \int_{\Omega_{u^*(r+h)} \setminus \Omega_{u^*(r)}} K(x)e^{u(x)} dx = F'(r) \\ &\leq 2\pi r b e^{u^*(r)}. \end{aligned}$$

By (2.5), we have

$$(2.7) \quad a \leq \bar{K}(r) \leq b$$

for almost everywhere  $r$ .

We want to derive a differential inequality for  $F(r)$ . First, we let  $J_1 = \{\rho \geq 0 \mid \frac{du^*}{dr}$  does not exist at  $r = \rho\}$ ,  $J_2 = \{\rho \geq 0 \mid \frac{du^*}{dr} = 0 \text{ at } r = \rho\}$ , and  $E = u^*(J_1 \cup J_2) = \{t \mid t = u^*(r) \text{ for some } r \in J_1 \cup J_2\}$ . Since  $u^*$  is locally Lipschitz, it is not difficult to see  $H_1(E) = 0$ , where  $H_1$  denotes the one dimensional Hausdorff measure of  $E$ . For  $t \notin E$  and  $r$  satisfying  $u^*(r) = t$ , we have  $\frac{du^*(r)}{dr} \neq 0$  and

$$(2.8) \quad -\frac{d}{dt}|\Omega_t| = \frac{d}{dr}|\Omega_t| \cdot \left(-\frac{du^*}{dr}\right)^{-1} = 2\pi r \left(-\frac{du^*}{dr}\right)^{-1}.$$

It is well-known that for almost everywhere  $t$ , the inequality

$$(2.9) \quad -\frac{d}{dt}|\Omega_t| \geq \int_{\{u=t\}} |\nabla u|^{-1} dH_1,$$

holds. (See the Lemma in §2.3 in [BZ]). For such  $t$ , by (2.8) and the isoperimetric inequality, we have

$$\begin{aligned}
 (2\pi r)^2 &= H_1(\partial\Omega_t^*)^2 \leq \left( \int_{\partial\Omega_t} dH_1 \right)^2 \\
 (2.10) \quad &\leq \left( \int_{\partial\Omega_t} \frac{1}{|\nabla u|} dH_1 \right) \left( \int_{\partial\Omega_t} |\nabla u| dH_1 \right) \\
 &\leq - \left( \frac{du^*(r)}{dr} \right)^{-1} (2\pi r) \int_{\partial\Omega_t} |\nabla u| dH_1.
 \end{aligned}$$

Thus

$$\begin{aligned}
 (2.11) \quad -\frac{du^*(r)}{dr}(2\pi r) &\leq \int_{\partial\Omega_t} |\nabla u| dH_1 = - \int_{\partial\Omega_t} \frac{\partial u}{\partial \nu} dH_1 = - \int_{\Omega_t} \Delta u \, dx \\
 &= \int_{\Omega_t} K e^u \, dx = F(r),
 \end{aligned}$$

where  $t = u^*(r)$ . Let  $\tilde{E} = \{t \notin E \mid (2.9) \text{ does not hold}\}$ , and  $J_3 = \{r \mid u^*(r) \in \tilde{E}\}$ . Since  $H_1(\tilde{E}) = 0$  and  $\frac{du^*}{dr}(r) \neq 0$  for any  $r \in J_3$ , we have  $H_1(J_3) = 0$ . If  $r \in J_2$ , (2.11) holds trivially. Therefore, we conclude that (2.11) holds for almost everywhere  $r$ . By (2.5) and (2.11), we have for almost everywhere  $r$ , that is,  $r \notin J_1 \cup J_3$ ,

$$\begin{aligned}
 (2.12) \quad \frac{d}{dr} \left( \frac{rF'(r)}{K(r)} \right) &= 4\pi r e^{u^*(r)} + 2\pi r^2 e^{u^*} \frac{du^*}{dr} \\
 &\geq \frac{2F'(r)}{K(r)} - \frac{F'(r)F(r)}{2\pi K(r)} \\
 &= \frac{2F'}{K} \left( 1 - \frac{F}{4\pi} \right) \\
 &\geq \begin{cases} \frac{2F'}{a} \left( 1 - \frac{F(r)}{4\pi} \right) & \text{if } F(r) > 4\pi, \\ \frac{2F'}{b} \left( 1 - \frac{F(r)}{4\pi} \right) & \text{if } F(r) \leq 4\pi. \end{cases}
 \end{aligned}$$

Since  $\alpha = F(\infty) > 4\pi$ ,  $F(r_0) = 4\pi$  for some  $r_0$ . Therefore, by (2.12) and noting that  $\frac{rF'(r)}{K(r)}$  is Lipschitz, we have for  $r \geq r_0$ ,

$$(2.13) \quad \begin{aligned} \frac{rF'(r)}{K(r)} &\geq \frac{1}{b} \int_0^{r_0} 2F'(s) \left(1 - \frac{F(s)}{4\pi}\right) ds + \frac{1}{a} \int_{r_0}^r 2F'(s) \left(1 - \frac{F(s)}{4\pi}\right) ds \\ &= \frac{1}{4\pi a} \left( F(r) - 4\pi \left(1 - \sqrt{\frac{a}{b}}\right) \right) \left( 4\pi \left(1 + \sqrt{\frac{a}{b}}\right) - F(r) \right). \end{aligned}$$

Since we assume  $F(\infty) < +\infty$ , there is  $r_k \rightarrow +\infty$  such that  $\lim_{k \rightarrow +\infty} r_k F'(r_k) = 0$ . By (2.13), we have

$$F(\infty) \geq 4\pi \left(1 + \sqrt{\frac{a}{b}}\right)$$

since  $F(\infty) > 4\pi \geq 4\pi(1 - \sqrt{\frac{a}{b}})$ . The proof of part (i) is complete.

Suppose  $\alpha = 4\pi \left(1 + \sqrt{\frac{a}{b}}\right)$ . Then inequalities in (2.9), (2.10), (2.11) and (2.12) must become equalities for almost everywhere  $r \geq 0$ . In particular, we have for almost everywhere  $r \geq 0$ ,

$$(2.14) \quad \frac{du^*(r)}{dr} \neq 0$$

and

$$(2.15) \quad \frac{du^*(r)}{dr} = \frac{F(r)}{2\pi r}$$

Since both  $u^*$  and  $F(r)$  are Lipschitz, we have  $\frac{du^*(r)}{dr} \equiv \frac{F(r)}{2\pi r}$  for  $r > 0$ . From (2.10), we also have for almost everywhere  $t$ ,

$$(2.16) \quad \Omega_t \text{ is a ball and } |\nabla u| \equiv -\frac{du^*(r)}{dr} \text{ on } \partial\Omega_t$$

where  $t = u^*(r)$ .

Applying the coarea formula for  $u$ , we have

$$\int_{\mathbb{R}^2} |\nabla u|^2 = \int_{-\infty}^{M_0} \int_{u^{-1}(s)} |\nabla u| dH_1 ds$$

where  $M_0 = \max_{\mathbb{R}^2} u$  and  $u^{-1}(s) = \{x | u(x) = s\}$ . Let  $s = u^*(\rho)$ . By (2.16), the coarea formula implies

$$\begin{aligned}
 \int_{\mathbb{R}^2} |\nabla u|^2 &= \int_{-\infty}^{M_0} \left( \frac{-du^*(\rho)}{d\rho} \right) (2\pi\rho) ds \\
 (2.17) \qquad &= 2\pi \int_0^\infty \left( \frac{du^*(\rho)}{d\rho} \right)^2 \rho \, d\rho \\
 &= \int_{\mathbb{R}^2} |\nabla u^*(x)|^2 dx .
 \end{aligned}$$

By (2.14), (2.17) and Theorem 1.1 in [BZ], we conclude that  $u(x) = u^*(x + x_0)$  for some  $x_0 \in \mathbb{R}^2$ . Without loss of generality, we may assume  $x_0 = 0$ . By the equation (1.1),  $K(x) = -\Delta u e^{-u}$  is also radially symmetric with respect to 0. By the equality in (2.12), we have

$$K(x) = \overline{K}(|x|) = \begin{cases} b & \text{if } 0 \leq |x| \leq r_0 \\ a & \text{if } |x| \geq r_0 \end{cases}$$

holds for a.e  $x$ , where  $r_0$  satisfies  $F(r_0) = 4\pi$ . Hence (1.1) reduces to the ordinary differential equation

$$(2.18) \quad \begin{cases} u''(r) + \frac{u'(r)}{r} + be^{u(r)} = 0 & \text{for } 0 \leq r \leq r_0, \\ u''(r) + \frac{u'(r)}{r} + ae^{u(r)} = 0 & \text{for } r \geq r_0 \end{cases}$$

with  $F(r_0) = 4\pi$ . By elementary calculations, we can show that  $u$  has the form of (1.5). Thus the proof of Theorem 1.1 is completely finished.  $\square$

*Proof of Theorem 1.2.* Let  $u^*$  and  $\overline{K}$  be defined as in the proof of Theorem 1.1. Since  $u$  and  $K$  are radially symmetric we have  $u = u^*$  and  $K = \overline{K}$  for a.e.  $x$ . By (1.1)

$$-2\pi r \frac{du}{dr} = - \int_{B_r(0)} \Delta u \, dx = F(r) .$$

The inequality (2.12) now becomes

$$\begin{aligned}
 \frac{d}{dr} \left( \frac{rF'(r)}{\overline{K}(r)} \right) &= \frac{2F'}{\overline{K}} \left( 1 - \frac{F}{4\pi} \right) \\
 (2.19) \qquad &\leq \begin{cases} \frac{2F'}{b} \left( 1 - \frac{F}{4\pi} \right) & \text{if } F > 4\pi , \\ \frac{2F'}{a} \left( 1 - \frac{F}{4\pi} \right) & \text{if } F \leq 4\pi . \end{cases}
 \end{aligned}$$



Let  $r_0$  satisfy  $F(r_0) = 4\pi$ . Again for  $r > r_0$ , we integrate to obtain

$$(2.20) \quad \frac{rF'(r)}{K(r)} \leq \frac{1}{b} \left[ \left( \frac{b}{a} - 1 \right) \cdot 4\pi + 2F(r) - \frac{F^2(r)}{4\pi} \right].$$

Since  $F(\infty) < \infty$ , there is  $\{r_k\}$  such that  $\lim_{k \rightarrow \infty} r_k = \infty$  and  $\lim_{k \rightarrow \infty} r_k F'(r_k) = 0$ . Let  $r = r_k$  in (2.20) and let  $k \rightarrow \infty$ , we conclude

$$0 \leq b^{-1} \left[ \left( \frac{b}{a} - 1 \right) \cdot 4\pi + 2F(\infty) - \frac{F^2(\infty)}{4\pi} \right]$$

and

$$F(\infty) \leq 4\pi \left( 1 + \sqrt{\frac{b}{a}} \right).$$

If  $\alpha = 4\pi \left( 1 + \sqrt{\frac{b}{a}} \right)$ , then the equality in (2.19) must hold. Thus  $K(x) = a$  for  $0 \leq |x| \leq r_0$  and  $K(x) = b$  for  $|x| > r_0$  hold for a.e.  $x$ . The form (1.7) then follows from the corresponding ordinary differential equation to (1.1).  $\square$

### 3. The sup+inf Type Inequality.

For the proof of Theorem 1.3, we need the following lemma, which was proved in [Su].

**Lemma 3.1.** *Let  $u$  be a Lipschitz function defined in  $B_R(0)$  and satisfy  $\Delta u + \lambda e^u \leq 0$  in  $B_R(0)$ . Then*

$$(3.1) \quad u(0) \leq \frac{1}{2\pi r} \int_{\partial B_r(0)} u \, ds - 2 \log \left\{ 1 - \frac{1}{8\pi} \int_{B_r(0)} \lambda e^u \, dx \right\}_+$$

holds for  $0 < r < R$  where  $\{\cdot\}_+ = \max\{\cdot, 0\}$ .

*Proof of Theorem 1.3.* Since (1.10) is a special case of (1.12), it suffices to prove (1.12). Assume the conclusion does not hold, that is, there are  $u_i$  and  $K_i$  which satisfy (1.2), (1.8) and (1.11) for some  $\rho, \sigma$  and  $B$  such that

$$(3.2) \quad \frac{1}{\sqrt{\sigma}} \sup_A u_i + \inf_\Omega u_i \rightarrow \infty$$

as  $i \rightarrow \infty$ . Then  $\lim_{i \rightarrow \infty} (\sup_A u_i) = \infty$  and there is  $\{x_i\} \subset A$  such that  $u_i(x_i) = \sup_A u_i$  and  $\lim_{i \rightarrow \infty} u_i(x_i) = \infty$ .

Step 1. We will employ a blow up argument to show that if we rescale the functions  $u_i$ , then there is a subsequence of these functions which converges to a solution  $w$  on entire  $\mathbb{R}^2$  with a minimal total curvature  $\alpha_w$ . That is,  $\alpha_w = 4\pi(1 + \sqrt{\frac{a}{b}})$  and  $w$  has the form (1.5).

Let  $d(S, T)$  denote the distance between two sets  $S$  and  $T$ . When  $S$  is a one-element set  $\{c\}$ , we will also use  $d(c, T)$  to denote  $d(\{c\}, T)$ . Let  $d_0 = \frac{1}{4}d(A, \partial\Omega)$ ,  $\Omega_1 = \{x \in \Omega \mid d(x, \partial\Omega) > d_0\}$  and

$$g_i(x) = u_i(x) + 2 \log(d(x, \partial\Omega_1)) .$$

Then  $\lim_{i \rightarrow \infty} g_i(x_i) = \infty$  and the maximum of  $g_i$  occurs at some point  $\bar{x}_i \in \Omega_1$ . We have

$$u_i(\bar{x}_i) \geq u_i(x_i) - C(\Omega)$$

with  $C(\Omega)$  depending on the diameter of  $\Omega$ . Let  $M_i = u_i(\bar{x}_i)$ ,  $\bar{L}_i = \frac{1}{2}d(\bar{x}_i, \partial\Omega_1)e^{\frac{M_i}{2}}$ , and

$$v_i(y) = u_i(e^{-\frac{M_i}{2}}y + \bar{x}_i) - M_i .$$

Then

$$(3.3) \quad \Delta v_i + \bar{K}_i(y)e^{v_i} = 0 \quad \text{for } |y| \leq \bar{L}_i ,$$

where  $\bar{K}_i(y) = K_i(e^{-\frac{M_i}{2}}y + \bar{x}_i)$  and  $\lim_{i \rightarrow \infty} \bar{L}_i \geq \frac{1}{2} \lim_{i \rightarrow \infty} \exp[\frac{1}{2}g_i(\bar{x}_i)] = \infty$ . For  $|y| \leq \bar{L}_i$ ,

$$\begin{aligned} v_i(y) &\leq g_i\left(\bar{x}_i + e^{-\frac{M_i}{2}}y\right) - g_i(\bar{x}_i) + 2 \log \frac{d(\bar{x}_i, \partial\Omega_1)}{d(\bar{x}_i + e^{-\frac{M_i}{2}}y, \partial\Omega_1)} \\ &\leq 2 \log 2 . \end{aligned}$$

Let  $(r, \theta)$  be the polar coordinate in  $\mathbb{R}^2$ , then

$$\left| \int_0^{2\pi} \frac{\partial v_i}{\partial r} r d\theta \right| = \left| \int_{|y| \leq r} \bar{K}_i(y) e^{v_i} dy \right| \leq cr^2$$

and

$$\begin{aligned} \int_0^{2\pi} |v_i(r, \theta)| d\theta &\leq \left| \int_0^{2\pi} v_i(r, \theta) d\theta \right| + 4\pi(2 \log 2) \\ &\leq \int_0^r \left| \int_0^{2\pi} \frac{\partial v_i}{\partial s}(s, \theta) d\theta \right| ds + 4\pi(2 \log 2) \\ &\leq c(1 + r^2) \end{aligned}$$

Thus for  $|y| \leq l$  with any fixed  $l$

$$\begin{aligned} v_i(y) &= \int_{|\eta| \leq 2l} G(y, \eta) \overline{K}_i(\eta) e^{v_i(\eta)} d\eta - \int_{|\eta|=2l} \frac{\partial G}{\partial r} v_i ds \\ &\geq -c \left[ \int_{|\eta| \leq 2l} G(y, \eta) d\eta + \int_0^{2\pi} |v_i(2l, \theta)| d\theta \right] \\ &\geq -c(1 + l^2) \end{aligned}$$

where  $G(y, \eta)$  is the Green function of  $-\Delta$  on  $B_{2l}(0)$ . By the elliptic theory, we can pass to subsequences of  $\{v_i\}$  and  $\{\overline{K}_i\}$  and assume on any compact set of  $\mathbb{R}^2$ ,

$$\begin{cases} v_i \rightarrow w & \text{in } C^{1,s} \text{ for any } 0 < s < 1 \\ \overline{K}_i \overset{*}{\rightarrow} K_0 & \text{weakly* in } L^\infty \end{cases}$$

as  $i \rightarrow \infty$ . Moreover,  $a \leq K_0(y) \leq b$  and

$$\Delta w + K_0(y) e^w = 0 \quad \text{in } \mathbb{R}^2.$$

By (1.11)

$$\frac{\overline{K}_i(y)}{\overline{K}_i(z)} \leq \sigma + \frac{B}{\left| \log \left( e^{-\frac{M_i}{2}} |y - z| \right) \right|} \rightarrow \sigma$$

as  $i \rightarrow \infty$  if  $|y|, |z| \leq l$ . Thus the essential supremum and infimum of  $K_0$  satisfy

$$(3.4) \quad \frac{\text{ess. sup } K_0}{\text{ess. inf } K_0} \leq \sigma.$$

By theorem 1.1,

$$\alpha_w \equiv \int_{\mathbb{R}^2} K_0(y) e^w dy \geq 4\pi \left( 1 + \sqrt{\frac{\text{ess. inf } K_0}{\text{ess. sup } K_0}} \right) \geq 4\pi \left( 1 + \frac{1}{\sqrt{\sigma}} \right)$$

Recall  $\rho$  is the number in (1.11). Let  $\bar{\rho} = \frac{1}{2} \min(\rho, d(\partial\Omega_1, \partial\Omega))$ ,  $L_i = \bar{\rho} e^{\frac{M_i}{2}}$  and  $G(x, z)$  be the Green function of  $-\Delta$  on  $B_{\bar{\rho}}(\bar{x}_i)$ . Then

$$(3.5) \quad \begin{aligned} M_i &= \int_{|z-\bar{x}_i| \leq \bar{\rho}} K_i(z) e^{u_i(z)} G(\bar{x}_i, z) dz + s_i \\ &\geq \int_{|y| \leq L_i} \left( \frac{M_i}{4\pi} - \frac{\log \left| \frac{y}{\bar{\rho}} \right|}{2\pi} \right) \bar{K}_i(y) e^{v_i(y)} dy + s_i, \end{aligned}$$

where  $s_i = \frac{1}{2\pi\bar{\rho}} \int_{|z-\bar{x}_i|=\bar{\rho}} u_i(z) ds$ . First, we claim that

$$\alpha_w = 4\pi \left( 1 + \sqrt{\frac{1}{\sigma}} \right).$$

For if  $\alpha_w > 4\pi \left( 1 + \sqrt{\frac{1}{\sigma}} \right)$ , then there are  $\epsilon_1 > 0$  and large  $l$  such that

$$\int_{|y| \leq l} K_0(y) e^w dy \geq (1 + 2\epsilon_1) 4\pi \left( 1 + \sqrt{\frac{1}{\sigma}} \right).$$

By (3.5)

$$(3.6) \quad \begin{aligned} M_i &\geq \int_{|y| \leq l} \left( \frac{M_i}{4\pi} - \frac{\log \left| \frac{y}{\bar{\rho}} \right|}{2\pi} \right) \bar{K}_i(y) e^{v_i} dy + s_i \\ &\geq \left( 1 - \frac{\epsilon_1}{1 + 2\epsilon_1} \right) \frac{M_i}{4\pi} \int_{|y| \leq l} \bar{K}_i e^{v_i} dy + s_i \\ &\geq (1 + \epsilon_1) 4\pi \left( 1 + \sqrt{\frac{1}{\sigma}} \right) \cdot \frac{M_i}{4\pi} + s_i \end{aligned}$$

for large  $i$ . Thus

$$0 \geq \sqrt{\frac{1}{\sigma}} M_i + s_i$$

for large  $i$ . Since  $s_i \geq \inf_{\Omega} u_i$  and  $M_i \geq u_i(x_i) - C(\Omega)$ , it contradicts to (3.2). Hence  $\alpha_w = 4\pi \left( 1 + \sqrt{\frac{1}{\sigma}} \right)$  must hold. By Theorem 1.1,  $w$  is radially symmetric with respect to some point  $y_0 \in \mathbb{R}^2$ . After a translation, we may assume  $y_0 = 0$ .

Step 2. We will find a sequence  $l_i \leq L_i$  such that

$$(3.7) \quad \int_{|y| \leq l_i} \bar{K}_i e^{v_i} dy \geq 4\pi \left(1 + \sqrt{\frac{1}{\sigma}}\right) - C_1 M_i^{-1}$$

and

$$(3.8) \quad \int_{|y| \leq l_i} \log \left| \frac{y}{\rho} \right| \bar{K}_i e^{v_i} dy \leq C_2$$

for  $C_1$  and  $C_2$  independent of  $i$  as  $i \rightarrow \infty$ .

Let

$$(3.9) \quad l_i = \sup \left\{ l \leq L_i \mid \int_{|y| \leq l} \bar{K}_i(y) e^{v_i} dy \leq 4\pi \left(1 + \frac{1}{\sqrt{\sigma}}\right) \right\}.$$

Since  $v_i$  uniformly converges to  $w$  in any compact set of  $\mathbb{R}^2$  and  $w$  satisfies

$$\int_{\mathbb{R}^2} K_0(y) e^w dy = 4\pi \left(1 + \frac{1}{\sqrt{\sigma}}\right),$$

we have

$$\lim_{i \rightarrow +\infty} l_i = +\infty.$$

Let  $\epsilon$  be a small positive number which will be chosen later. Then for large  $i$ , there exist  $r_1^i$  and  $r_1$  such that

$$\int_{|y| \leq r_1} K_0(y) e^w dy = 4\pi \left(1 + \frac{1}{\sqrt{\sigma}}\right) - \epsilon,$$

and

$$(3.10) \quad \int_{|y| \leq r_1^i} \bar{K}_i(y) e^{v_i(y)} dy = 4\pi \left(1 + \frac{1}{\sqrt{\sigma}}\right) - \epsilon.$$

Obviously,  $\lim_{i \rightarrow +\infty} r_1^i = r$ . Hence,

$$(3.11) \quad \int_{r_1^i \leq |y| \leq l_i} \bar{K}_i e^{v_i} dy \leq \epsilon.$$

Since  $0 < a \leq K_i \leq b$ , we have  $\Delta v_i + b e^{v_i} \geq 0$  in  $B_{l_i}(0)$ . Thus, if  $\epsilon$  is small enough, we can apply Lemma 3.1 to  $v_i$ . By Lemma 3.1, the inequality

$$\begin{aligned} v_i(x) &\leq \frac{1}{2\pi r} \int_{\partial B_r(x)} v_i ds - 2 \log \left\{ 1 - \frac{1}{8\pi} \int_{B_r(x)} b e^u dx \right\}_+ \\ &\leq \frac{1}{2\pi r} \int_{\partial B_r(x)} v_i ds + \log 4. \end{aligned}$$

holds for  $r \leq \frac{|x|}{2}$  and  $2r_1^i \leq |x| \leq \frac{l_i}{2}$  if  $\epsilon$  is sufficiently small in (3.11) such that

$$(3.12) \quad b \int_{r_1^i \leq |y| \leq l_i} e^{v_i(y)} dy \leq 4\pi .$$

Thus,

$$v_i(x) \leq \frac{1}{\pi r^2} \int_{B_r(x)} v_i(y) dy + \log 4 .$$

By Jensen inequality and letting  $r = \frac{|x|}{2}$ , we have

$$(3.13) \quad \begin{aligned} e^{v_i(x)} &\leq 4 \exp \left( \frac{1}{\pi r^2} \int_{B_r(x)} v_i \right) \leq \frac{4}{\pi r^2} \int_{B_r(x)} e^{v_i(y)} dy \\ &\leq \frac{4}{a\pi r^2} \int_{r_1^i \leq |y| \leq l_i} \bar{K}_i(y) e^{v_i(y)} dy \\ &\leq \frac{16\epsilon}{\pi a} |x|^{-2} = C_1 \epsilon |x|^{-2} . \end{aligned}$$

for  $2r_1^i \leq |x| \leq \frac{l_i}{2}$ . Let

$$\bar{v}_i(r) = \frac{1}{2\pi r} \int_{|y|=r} v_i(y) dy .$$

Then for  $r \leq l_i$ ,

$$\begin{aligned} \frac{d}{dr} \bar{v}_i(r) &= -\frac{1}{2\pi r} \int_{|y| \leq r} \Delta v_i dy = -\frac{1}{2\pi r} \int_{|y| \leq r} \bar{K}_i e^{v_i} dy \\ &\geq -2 \left( 1 + \frac{1}{\sqrt{\sigma}} \right) \frac{1}{r} . \end{aligned}$$

Integrating the inequality above gives

$$(3.14) \quad \bar{v}_i(r) \geq -2 \left( 1 + \frac{1}{\sqrt{\sigma}} \right) \log r + C_2$$

for some constant  $C_2$  and  $r \leq l_i$ . We want to apply the Harnack inequality to obtain a lower bound for  $v_i$ . To see this, we employ (3.13) and have for  $4r_1^i \leq r \leq \frac{l_i}{4}$ ,

$$(3.15) \quad v_i(x) + 2 \log r \leq \log C_1 + \log \epsilon \leq \frac{1}{2} \log \epsilon < 0 \quad \text{for } \frac{r}{2} \leq |x| \leq 2r$$

provided  $\epsilon$  is small enough. From now on  $\epsilon$  is a fixed small positive number such that (3.12) and (3.15) hold. Let

$$\tilde{v}_i(x) = v_i(rx) + 2 \log r$$

for  $\frac{1}{2} \leq |x| \leq 2$ . Then  $\tilde{v}_i(x) \leq \frac{1}{2} \log \epsilon$  in  $\frac{1}{2} \leq |x| \leq 2$  and  $\tilde{v}_i$  satisfies

$$\Delta \tilde{v}_i + \overline{K}_i(rx) e^{\tilde{v}_i} = 0 \quad \text{in } \frac{1}{2} \leq |x| \leq 2 .$$

Since  $\frac{\overline{K}_i(rx) e^{\tilde{v}_i}}{|\tilde{v}_i|} \leq C_1 \frac{\epsilon}{|\log \epsilon|}$  by (3.13) and (3.15), the Harnack inequality can be applied to  $-\tilde{v}_i$ . Hence there exists a constant  $C \geq 1$  such that for  $|x| = 1$

$$-\tilde{v}_i(x) \leq -C \int_{|x|=1} \tilde{v}_i .$$

Going back to the function  $v_i$ , we have from (3.14)

$$v_i(x) \geq -2 \log |x| - \frac{2C}{\sqrt{\sigma}} \log |x| + C_3$$

for  $4r_1^i \leq |x| \leq \frac{l_i}{4}$ . Therefore, letting  $\delta = \left(2 + \frac{2C}{\sqrt{\sigma}}\right)^{-1}$ , we have for  $4r_1^i \leq |x| = R^\delta \leq \frac{l_i}{4}$

$$(3.16) \quad \begin{aligned} v_i(x) &\geq -2 \log R + C_3 \\ &\geq \sup_{\frac{l_i}{2} \geq |y| \geq R} v_i(y) \end{aligned}$$

Obviously,  $v_i(x) \geq \sup_{\frac{l_i}{2} \geq |y| \geq R} v_i(y)$  holds true also for  $|x| = R^\delta \leq 4r_1^i$  since

$w(x) = w(|x|)$  is strictly decreasing in  $|x|$  and  $v_i$  converges to  $w$  on any compact set.

Let  $m_i(r) = \max_{|y|=r} v_i(y)$  and  $t_0 = m_i\left(\frac{l_i}{2}\right)$ . For  $t > t_0$ , let  $\Omega_t^i = \{y \mid |y| \leq \frac{l_i}{2}$

and  $v_i(y) > t\}$ . Obviously, the closure  $\overline{\Omega}_t^i$  of  $\Omega_t^i$  is always contained in the open ball  $B_{\frac{l_i}{2}}(0)$ . Let  $v_i^*(x) = v_i^*(|x|)$  denote the Schwartz symmetrization, that is,  $v_i^*(|x|)$  is nonincreasing in  $|x|$  and  $\{x \mid v_i^*(x) > t\}$  is the ball  $B_r(0)$  with  $r = (\pi^{-1} |\Omega_t^i|)^{\frac{1}{2}}$  for  $t > t_0$ . Since  $v_i$  satisfies the equation (3.3),  $|\{y \mid |y| \leq l_i, v_i(y) = t > t_0\}| = 0$ . Hence  $|\Omega_t^i|$  is strictly decreasing in  $t$  and then  $v_i^*(|x|)$  is strictly decreasing in  $r$ . We also have  $v_i^*(|x|)$  is locally Lipschitz. As in Section 2, we let

$$(3.17) \quad F_i(r) = \int_{\Omega_{v_i^*(r)}} \overline{K}_i(y) e^{v_i} dy ,$$

and

$$(3.18) \quad \hat{K}_i(r) = \frac{F'_i(r)}{2\pi r e^{v_i^*(r)}}$$

for all  $r$  such that  $v_i^*(r)$  is well-defined. As before, we can prove that  $F_i(r)$  is locally Lipschitz, therefore,  $\hat{K}_i(r)$  is defined for almost everywhere  $r$ .

We note that an immediate consequence of (3.16) is, for  $R \leq \frac{l_i}{4}$ :

$$(3.19) \quad B_{R\delta^2}(0) \subseteq \Omega_m^i(R\delta) \subseteq B_R(0) .$$

The first part of (3.19) is easily seen from (3.16). For the second part of (3.19), let  $z \in \Omega_m(R\delta)$ . Then by (3.16), we have

$$v_i(z) > \max_{|y|=R\delta} v_i(y) \geq \max_{\frac{l_i}{2} \geq |y| \geq R} v_i(y) .$$

Hence  $z \in B_R(0)$ , and (3.19) is proved. From (3.19), we immediately have for  $R \leq \frac{l_i}{4}$ ,

$$(3.20) \quad \int_{|y| \leq R} \bar{K}_i(y) e^{v_i(y)} dy \geq F \left( \left( \pi^{-1} \left| \Omega_m(R\delta) \right| \right)^{\frac{1}{2}} \right) \geq F \left( R\delta^2 \right) .$$

Let  $R_i = \left( \pi^{-1} \left| \Omega_m \left( \left( \frac{l_i}{4} \right)^{\frac{\delta}{2}} \right) \right| \right)^{\frac{1}{2}}$ . Then

$$(3.21) \quad \left( \frac{l_i}{4} \right)^{\frac{\delta^2}{2}} \leq R_i \leq \left( \frac{l_i}{4} \right)^{\frac{1}{2}} .$$

Thus,  $\lim_{i \rightarrow +\infty} R_i = +\infty$ . Obviously,  $v_i^*(r)$  is defined for all  $r \leq R_i$ . By (3.18) and (2.6), (3.19) implies for  $r \leq R_i$ ,

$$(3.22) \quad a_i \leq \hat{K}_i(r) \leq b_i$$

where

$$(3.23) \quad a_i = \text{ess. inf}_{|y| \leq \left( \frac{l_i}{4} \right)^{\frac{1}{2}}} \bar{K}_i(y) ,$$

and

$$(3.24) \quad b_i = \text{ess. sup}_{|y| \leq \left( \frac{l_i}{4} \right)^{\frac{1}{2}}} \bar{K}_i(y) .$$



Let  $r_0^i < r_1^i$  satisfy  $F_i(r_0^i) = 4\pi$ . As in section 2, we can derive a similar differential inequality for  $F_i(r)$  as in (2.13), that is, the inequality

$$(3.25) \quad \frac{rF_i'(r)}{\tilde{K}_i(r)} \geq \frac{-1}{4\pi a_i} \left[ F(r) - 4\pi \left( 1 - \sqrt{\frac{a_i}{b_i}} \right) \right] \left[ F(r) - 4\pi \left( 1 + \sqrt{\frac{a_i}{b_i}} \right) \right]$$

holds true for  $r_0^i \leq r \leq R_i$  almost everywhere. Let  $\tilde{R}_i = \sup \left\{ r \leq R_i \mid F_i(r) \leq 4\pi \left( 1 + \sqrt{\frac{a_i}{b_i}} \right) \right\}$ . By (3.25), we have

$$\frac{F_i'(r)}{F_i(r) - 4\pi \left( 1 - \sqrt{\frac{a_i}{b_i}} \right)} + \frac{F_i'(r)}{4\pi \left( 1 + \sqrt{\frac{a_i}{b_i}} \right) - F(r)} \geq 2\sqrt{\frac{a_i}{b_i}} r^{-1}$$

for  $r_0^i \leq r \leq \tilde{R}_i$ . Integrating the differential inequality gives

$$\log \frac{4\pi \left( 1 + \sqrt{\frac{a_i}{b_i}} \right) - F_i(r)}{F_i(r) - 4\pi \left( 1 - \sqrt{\frac{a_i}{b_i}} \right)} \leq -2\sqrt{\frac{a_i}{b_i}} \log \frac{r}{r_0^i},$$

which implies

$$(3.26) \quad F_i(r) \geq 4\pi \left( 1 + \sqrt{\frac{a_i}{b_i}} \right) - Cr^{-2\sqrt{\frac{a_i}{b_i}}}$$

for  $r_0^i \leq r \leq \tilde{R}_i$ , where  $C$  is a positive constant independent of  $i$ . Trivially, (3.26) holds for  $\tilde{R}_i \leq r \leq R_i$  also.

By (3.23) and (3.24), we have

$$\frac{b_i}{a_i} \leq \sup_{|y|, |z| \leq (\frac{1}{4})^{\frac{1}{2}}} \frac{\overline{K}_i(y)}{\overline{K}_i(z)} \leq \sigma + C_0 \left| \log \left( l_i^{\frac{1}{2}} e^{-\frac{M_i}{2}} \right) \right|^{-1} \leq \sigma + C_1 M_i^{-1}.$$

Therefore,

$$(3.27) \quad \sqrt{\frac{a_i}{b_i}} \geq \frac{1}{\sqrt{\sigma}} - C_2 M_i^{-1}.$$

Combined with (3.26), it yields

$$(3.28) \quad F_i(r) \geq 4\pi \left( 1 + \frac{1}{\sqrt{\sigma}} \right) - Cr^{-2\sqrt{\frac{1}{\sigma}}} - C_2 M_i^{-1}$$

for  $r_0^i \leq r \leq R_i$ . By (3.20) and (3.28), we have

$$(3.29) \quad \int_{|y| \leq R} \bar{K}_i e^{v_i} dy \geq F(R^{\delta^2}) \geq 4\pi \left(1 + \sqrt{\frac{1}{\sigma}}\right) - CR^{-2\sqrt{\frac{1}{\sigma}}\delta^2} - C_2 M_i^{-1}$$

for  $R \leq \left(\frac{l_i}{4}\right)^{\frac{1}{2}}$ . Hence,

$$(3.30) \quad \int_{R \leq |y| \leq l_i} \bar{K}_i e^{v_i} dy \leq CR^{-2\sqrt{\frac{1}{\sigma}}\delta^2} + C_2 M_i^{-1}.$$

Applying Lemma 3.1 and (3.13) again, there exists  $C_3 > 0$  such that

$$(3.31) \quad e^{v_i(x)} \leq \frac{4}{a\pi r^2} \int_{\frac{1}{2}|x| \leq y \leq l_i} \bar{K}_i e^{v_i} dy \leq C_3 \left(|x|^{-2-2\sqrt{\frac{1}{\sigma}}\delta^2} + M_i^{-1}|x|^{-2}\right)$$

for  $|x| \leq \frac{1}{2} \left(\frac{l_i}{4}\right)^{\frac{1}{2}}$ . Choose  $l_i^*$  satisfy  $(\log l_i^*)^2 = \log l_i$ . Obviously,  $l_i^* \leq \frac{1}{2} \left(\frac{l_i}{4}\right)^{\frac{1}{2}}$  for large  $i$ . Then, by (3.30) and (3.31), we have

$$\begin{aligned} & \int_{|y| \leq l_i} \log \frac{|y|}{\rho} \bar{K}_i(y) e^{v_i(y)} dy \\ & \leq \int_{|y| \leq l_i^*} \log \frac{|y|}{\rho} \bar{K}_i(y) e^{v_i(y)} dy + \log \frac{l_i}{\rho} \int_{l_i^* \leq |y| \leq l_i} \bar{K}_i(y) e^{v_i(y)} dy \\ & \leq C_4 \left[ \left(1 + M_i^{-1} (\log l_i^*)^2\right) + (\log l_i) \left(l_i^{*-2\sqrt{\frac{1}{\sigma}}\delta^2} + M_i^{-1}\right) \right] \\ & \leq C_5, \end{aligned}$$

which obviously yields (3.8). If  $l_i < L_i$ , then (3.7) holds trivially. If  $l_i = L_i$ , then by letting  $R = \left(\frac{l_i}{4}\right)^{\frac{1}{2}}$  in (3.29), it yields

$$\int_{|y| \leq L_i} \bar{K}_i e^{v_i} dy \geq 4\pi \left(1 + \frac{1}{\sqrt{\sigma}}\right) - C_3 l_i^{-\delta^2\sqrt{\frac{1}{\sigma}}} - C_2 M_i^{-1}$$

which (3.7) follows immediately.

Step 3. To obtain a contradiction to (3.2), we note that by (3.5), (3.7) and (3.8),

$$\begin{aligned} M_i & \geq \int_{|y| \leq l_i} \left( \frac{M_i}{4\pi} - \frac{\log \frac{|y|}{\rho}}{2\pi} \right) \bar{K}_i(y) e^{v_i(y)} dy + s_i \\ & \geq \left(1 + \frac{1}{\sqrt{\sigma}}\right) M_i + \inf_{\Omega} u_i - C_6, \end{aligned}$$

where  $C_6$  is a constant. Thus, we have

$$\frac{1}{\sqrt{\sigma}} M_i + \inf_{\Omega} u_i \leq C_6 ,$$

which obviously leads to a contradiction to (3.2). The proof of Theorem 1.3 is completely finished.  $\square$

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