

Functorial quantum field theory in the Riemannian setting

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We construct examples of functorial quantum field theories in the Riemannian setting by quantizing free massive bosons.

1. Introduction

There are a number of mathematical approaches to rigorously define Quantum Field Theory. As suggested by the path integral quantization, in the framework pioneered by Atiyah, Kontsevich, Segal and many others, roughly a d -dimensional Riemannian QFT is a rule E that assigns to a d -dimensional closed oriented Riemannian manifold Σ a number $E(\Sigma)$, which depends only on the isomorphism class of Σ . More generally, E assigns to a d -dimensional compact oriented Riemannian manifold Σ with $\partial\Sigma = \bar{Y}_0 \sqcup Y_1$ Hilbert spaces $E(Y_0)$ and $E(Y_1)$ together with a bounded operator $E(\Sigma) : E(Y_0) \rightarrow E(Y_1)$. The main requirement is that E is local. Here, by local we mean if Σ is obtained by gluing Σ_1 and Σ_2 along the common boundary Y , then the contributions from Σ_1 and Σ_2 are sufficient to compute $E(\Sigma)$.

More precisely, let $d\text{-RBord}$ be the d -dimensional Riemannian bordism “category” whose objects are oriented closed $d - 1$ dimensional Riemannian manifolds. We use Y s to denote objects of this category. A morphism $\Sigma : Y_0 \rightarrow Y_1$ in $d\text{-RBord}$ is a d -dimensional compact oriented Riemannian manifold with orientation preserving isometry $\bar{Y}_0 \sqcup Y_1 \rightarrow \partial\Sigma$. Here, \bar{Y}_0 means the intrinsic orientation on Y_0 is reversed. We require a morphism Σ in $d\text{-RBord}$ to have a *product metric* near the boundary so that we have a well defined composition of morphisms by gluing. We remark that $d\text{-RBord}$ is just a semicategory not a category as it does not have identity morphisms.

Let \mathbf{Hilb} be the category whose objects are separable Hilbert spaces and morphisms are continuous linear operators. We note \mathbf{Hilb} is a symmetric monoidal category with the monoidal structure given by the Hilbert space tensor product.

Following Atiyah[1] and Segal [20], the discussion above can be summarized as

Definition 1.1. A $(d, d - 1)$ Riemannian Functorial QFT is a functor $E : d\text{-RBord} \rightarrow \mathbf{Hilb}$ that maps disjoint unions into tensor products.

We will also consider a more general bordism category in Subsection 3.5, namely objects and morphisms come with a fixed vector bundle with a metric and a compatible connection.

Our main goal here is to construct a class of examples of $(d, d - 1)$ Riemannian FQFTs. Naively, one can try to do path integral quantization to construct such examples. Path integral quantization heuristically means construction of “measures” on the space of fields given by a classical field theory.

A classical field theory on a d -dimensional compact oriented Riemannian manifold Σ consists of two pieces of data: space of fields $\mathfrak{F}(\Sigma)$ and an action functional S which is a real valued function on $\mathfrak{F}(\Sigma)$. Furthermore, one requires S to be local which means S depends only on fields and their derivatives.

The space of fields $\mathfrak{F}(\Sigma)$ is space of sections of a fiber bundle P over Σ . An example of space of fields is $C^\infty(\Sigma)$, the space of real valued smooth functions on Σ and an example of action functional is the energy functional

$$(1.1) \quad S(\phi) = \frac{1}{2} \int_{\Sigma} d\phi \wedge *d\phi + m^2 * \phi^2$$

where $*$ is the Hodge star operator associated to the Riemannian metric and m is a positive real number.

The idea behind path integral quantization is to make sense of the integrals

$$E(\Sigma) = \int_{\mathfrak{F}(\Sigma)} e^{-S(\phi)} D\phi$$

where $D\phi$ is a “volume measure” on $\mathfrak{F}(\Sigma)$ and it is very difficult to achieve in general. Fortunately, using ideas from constructive quantum field theory [8], it is possible to construct “measures” in some linear σ -models. However, these measures, which are required to define path integrals, live on the space of distributions not on the space of fields. This does not create serious problems in some nice cases. In [15], Douglas Pickrell follows this point of view to construct a class of examples of 2-dimensional FQFTs called $P(\phi)_2$ -interaction theories. We use ideas from Pickrell[15] and Segal [19] to show

Theorem 1.2. *There is a projective representation of $d\text{-RBord}$ in \mathbf{Hilb} .*

Our main result is to promote this projective representation to get a FQFT when d is even. More precisely, we prove

Theorem 1.3. *There exists d -dimensional Riemannian FQFT when d is even.*

In fact, Theorem 1.3 can be used to construct FQFT in all dimensions as follows. A closed oriented Riemannian 1-manifold S induces a functor $d\text{-RBord} \rightarrow (d + 1)\text{-RBord}$ by taking product with S . When d is odd, composing this functor with the one given by Theorem 1.3 gives a FQFT. This shows that:

Corollary 1.4. *There are FQFTs for d odd as well.*

Acknowledgements

This paper is based on a part of the author’s PhD thesis at the University of Notre Dame. The author is grateful to his advisor Stephan Stolz for the discussions, guidance and support. The author also thanks Brian Hall and Liviu Nicolaescu for their inspiring ideas and comments. The author benefited from conversation with Owen Gwilliam, Douglas Pickrell, Pavel M̃nev and Augusto Stoffel. It is the author’s pleasure to thank Sayonita Ghosh Hajra for feedback. The author acknowledges support from Max Planck Institute for Mathematics, Bonn, where a part of this work was written.

2. Background

2.1. Dirichlet-to-Neumann operator

Let Σ be a d -dimensional compact oriented Riemannian manifold, $C^\infty(\Sigma)$ the space of real valued smooth functions on Σ , Δ_Σ the nonnegative Laplacian on Σ and $m > 0$.

Assume $Y = \partial\Sigma \neq \emptyset$. Let $i : Y \rightarrow \Sigma$ be the inclusion and let ν denote the outward unit normal to Y . We recall a classical fact that the Dirichlet problem

$$(\Delta_\Sigma + m^2)\phi = 0 \quad \text{with} \quad \phi|_Y = \eta$$

has a unique solution ϕ_η in $C^\infty(\Sigma)$ for all $\eta \in C^\infty(Y)$ ([17], p. 193). We call the unique solution ϕ_η to the Dirichlet problem the *Helmholtz extension* of η .

Definition 2.1. The operator on Y defined by

$$D_\Sigma \eta = \frac{\partial \phi_\eta}{\partial \nu}$$

is called the Dirichlet-to-Neumann operator associated to the Helmholtz operator $\Delta_\Sigma + m^2$.

It is well known that D_Σ is a positive elliptic pseudo differential operator of order one as an operator on Y ([21], Chapter 7, Section 11) and it has the same principal as the operator $(\Delta_Y + m^2)^{\frac{1}{2}}$ ([12], Theorem 2.1). In particular, D_Σ is symmetric.

Let $\text{Helm}(\Sigma)$ denote the space of Helmholtz solutions on Σ and $\Phi(Y)$ denote $C^\infty(Y) \oplus C^\infty(Y)$. There is a map $\text{Helm}(\Sigma) \rightarrow \Phi(Y)$ given by $\Psi \mapsto (\Psi|_Y, \frac{\partial \Psi}{\partial \nu}|_Y)$. We also note that $\Phi(Y)$ can be given a canonical symplectic structure $\omega_{\Phi(Y)}$, which is the hyperbolic form induced by the L^2 -inner product on $C^\infty(Y)$.

Remark 2.2. The image of $\text{Helm}(\Sigma)$ under the map above can be identified with the graph of $D_\Sigma : C^\infty(Y) \rightarrow C^\infty(Y)$ inside $\Phi(Y)$ and hence is a Lagrangian subspace of $\Phi(Y)$.

2.1.1. Gluing Dirichlet-to-Neumann operators. Let Σ_1 be a d dimensional oriented compact Riemannian manifold with $\partial \Sigma_1 = Y_2 \sqcup \overline{Y_1}$ where $\overline{Y_1}$ means Y_1 with the induced orientation which is opposite of the intrinsic orientation on Y_1 . We say Y_1 is an incoming boundary and Y_2 an outgoing boundary. Let Σ_2 be another d dimensional oriented compact Riemannian manifold with $\partial \Sigma_2 = Y_3 \sqcup \overline{Y_2}$ and let $\text{Helm}(\Sigma_2) \circ \text{Helm}(\Sigma_1)$ be the fiber product $\text{Helm}(\Sigma_2) \times_{\Phi(Y_2)} \text{Helm}(\Sigma_1)$. More precisely $\text{Helm}(\Sigma_2) \circ \text{Helm}(\Sigma_1)$ contains pairs (ϕ_2, ϕ_1) where $\phi_2(\phi_1)$ is a Helmholtz solution on $\Sigma_2(\Sigma_1)$ such that their values and normal derivatives agree on Y_2 . Let $\Sigma_2 \circ \Sigma_1$ denote the Riemannian manifold glued along Y_2 . It is easy to see that there is a canonical isomorphism between $\text{Helm}(\Sigma_2) \circ \text{Helm}(\Sigma_2)$ and $\text{Helm}(\Sigma_2 \circ \Sigma_1)$. We can rephrase this discussion in terms of Dirichlet-to-Neumann operators.

We can write

$$D_{\Sigma_1} = \begin{bmatrix} A & B \\ B^t & D \end{bmatrix} : C^\infty(Y_2) \oplus C^\infty(Y_1) \rightarrow C^\infty(Y_2) \oplus C^\infty(Y_1)$$

where

$$\begin{aligned} A &: C^\infty(Y_2) \rightarrow C^\infty(Y_2), B : C^\infty(Y_1) \rightarrow C^\infty(Y_2), \\ D &: C^\infty(Y_1) \rightarrow C^\infty(Y_1) \text{ and } B^t : C^\infty(Y_2) \rightarrow C^\infty(Y_1). \end{aligned}$$

Now, let us describe the operators A , B , B^t and D . Given ϕ_2 in $C^\infty(Y_2)$, calculate the Helmholtz solution on Σ_1 with boundary value ϕ_2 on Y_2 and vanishing boundary value on Y_1 ; then $A(\phi_2)$ is the outward normal derivative of the Helmholtz solution along Y_2 . Given ϕ_1 in $C^\infty(Y_1)$, calculate the Helmholtz solution on Σ_1 with boundary value ϕ_1 on Y_1 and vanishing boundary value on Y_2 ; then $D(\phi_1)$ is the inward normal derivative of the Helmholtz solution along Y_1 . Similarly we define B and B^t . The notation B^t is used to remind us that D_{Σ_1} is a symmetric operator.

We write

$$D_{\Sigma_2} = \begin{bmatrix} K & L \\ L^t & M \end{bmatrix} : C^\infty(Y_3) \oplus C^\infty(Y_2) \rightarrow C^\infty(Y_3) \oplus C^\infty(Y_2).$$

Lemma 2.3.

$$D_{\Sigma_2 \circ \Sigma_1} = \begin{bmatrix} K - L(A+M)^{-1}L^t & -L(A+M)^{-1}B \\ -B^t(A+M)^{-1}L^t & D - B^t(A+M)^{-1}B \end{bmatrix}$$

as an operator from $C^\infty(\partial(\Sigma_2 \circ \Sigma_1))$ to itself.

Proof. We will only sketch the proof. For simplicity, we assume $Y_1 = \emptyset$. We need to show

$$D_{\Sigma_2 \circ \Sigma_1} \psi = K - L(D_{\Sigma_1} + M)^{-1}L^t \psi$$

for all $\psi \in C^\infty(Y_3)$. Let Ψ be the Helmholtz solution on $\Sigma_2 \circ \Sigma_1$ such that $\Psi|_{Y_2} = \psi$ and let Ψ_2 be the Helmholtz solution on Σ_2 such that $\Psi_2|_{Y_3} = \psi$ and $\Psi_2|_{Y_2} = 0$. Then, $\Psi - \Psi_2$ is the Helmholtz solution on Σ_2 such that $\Psi - \Psi_2|_{Y_3} = 0$ and $\Psi - \Psi_2|_{Y_2} = \phi$ where $\phi = \psi|_{Y_2}$. This implies

$$B\phi = D_{\Sigma_2 \circ \Sigma_1} \psi - K\psi \text{ and } B^t \psi = -(D_{\Sigma_1} + M)\phi.$$

Now, the fact that $(D_{\Sigma_1} + M)$ is invertible (Theorem 2.1, [4]) proves the lemma. □

Let Y be a closed oriented Riemannian manifold, $m > 0$ and $s \in \mathbb{R}$. The bilinear form on $C^\infty(Y)$ given by

$$\langle f, g \rangle = \int_Y f(\Delta_Y + m^2)^s g \, d\text{vol}(Y)$$

defines an inner product on $C^\infty(Y)$. We define the Sobolev space $W^s(Y)$ as the completion of $C^\infty(Y)$ with respect to this inner product.

Lemma 2.4. *Let Σ be a compact oriented Riemannian manifold and $\partial\Sigma = Y$. Then*

$$\alpha_\Sigma = D_\Sigma(\Delta_Y + m^2)^{-\frac{1}{2}}$$

defines a unique continuous positive operator on $W^{\frac{1}{2}}(Y)$.

Proof. Recall that D_Σ and $(\Delta_Y + m^2)$ are both positive pseudo differential operators of order one and both have same principal symbol. This means

$$D_\Sigma(\Delta_Y + m^2)^{-\frac{1}{2}} = I + R$$

where R is a pseudo differential operator of order at most -1 . This shows $(\Delta_Y + m^2)^{-\frac{1}{2}}D_\Sigma$ is a pseudo differential operator of order zero and, hence, it has a unique extension to a continuous linear operator on $W^{\frac{1}{2}}(Y)$. Positivity of α_Σ is obvious. □

Let Σ_1 and Σ_2 be as above. Recall that α_{Σ_1} , α_{Σ_2} and $\alpha_{\Sigma_2 \circ \Sigma_1}$ are continuous positive operators on $W^{\frac{1}{2}}(Y_3) \oplus W^{\frac{1}{2}}(Y_2)$, $W^{\frac{1}{2}}(Y_2) \oplus W^{\frac{1}{2}}(Y_1)$ and $W^{\frac{1}{2}}(Y_3) \oplus W^{\frac{1}{2}}(Y_1)$ respectively. We want to show α_{Σ_1} , α_{Σ_2} and $\alpha_{\Sigma_2 \circ \Sigma_1}$ are related in the similar fashion as D_{Σ_1} , D_{Σ_2} and $D_{\Sigma_2 \circ \Sigma_1}$. The following definition is useful to describe this relation.

Definition 2.5. Let H_1 , H_2 and H_3 be Hilbert spaces,

$$\mathcal{A}_1 : H_2 \oplus H_1 \rightarrow H_2 \oplus H_1 \text{ and } \mathcal{A}_2 : H_3 \oplus H_2 \rightarrow H_3 \oplus H_2$$

positive continuous operators with continuous inverses. We write

$$\mathcal{A}_1 = \begin{bmatrix} A & B \\ B^t & D \end{bmatrix} \text{ and } \mathcal{A}_2 = \begin{bmatrix} K & L \\ L^t & M \end{bmatrix},$$

and define

$$(2.1) \quad \mathcal{A}_2 \circ \mathcal{A}_1 = \begin{bmatrix} K-L(A+M)^{-1}L^t & -L(A+M)^{-1}B \\ -B^t(A+M)^{-1}L^t & D-B^t(A+M)^{-1}B \end{bmatrix}.$$

As a corollary of Lemma 2.4, we have

Corollary 2.6. $\alpha_{\Sigma_2} \circ \alpha_{\Sigma_1} = \alpha_{\Sigma_2 \circ \Sigma_1}$.

Next, we describe the idea behind the Definition 2.5. We will consider the finite dimensional situation. Informally, we can use \mathcal{A}_1 and \mathcal{A}_2 to construct Lagrangian subspaces in some symplectic vector spaces and the definition of $\mathcal{A}_2 \circ \mathcal{A}_1$ is just a consequence of ‘‘composing’’ Lagrangian subspaces. We refer to Lemma 2.11 for a precise statement.

Definition 2.7. Let V and W be two vector spaces. A linear relation $P : V \rightarrow W$ is a subspace of $W \oplus V$.

Let $P : V \rightarrow W$ and $Q : W \rightarrow U$ be two linear relations. The composition $Q \circ P$ of linear relations is the fiber product $Q \times_W P$ which is defined by

$$Q \circ P = \{(u, v) | \exists w \in W \text{ with } (u, w) \in U \oplus W, (w, v) \in W \oplus V\}.$$

Definition 2.8. Let (V, Ω_V) be a finite dimensional symplectic vector space and $V_{\mathbb{C}}$ be the complexification of V and abusing the notation, Ω_V be the complex bilinear extension of Ω_V . A subspace L of $V_{\mathbb{C}}$ is said to be a positive Lagrangian if L is a Lagrangian subspace and the Hermitian form defined by

$$\langle v, w \rangle = -i\Omega_V(\bar{v}, w)$$

is positive definite on L .

Let (V, Ω_V) be as above. We use the notation \bar{V} to denote the symplectic vector space $(V, -\Omega_V)$.

Lemma 2.9. Let U, V, W be finite dimensional symplectic vector spaces. Let P be a positive Lagrangian subspace of $V_{\mathbb{C}} \oplus \bar{U}_{\mathbb{C}}$ and Q be a positive Lagrangian subspace of $W_{\mathbb{C}} \oplus \bar{V}_{\mathbb{C}}$. Then $Q \circ P$ is a positive Lagrangian subspace of $W_{\mathbb{C}} \oplus \bar{U}_{\mathbb{C}}$.

Proof. Using the well known fact the composition of Lagrangian linear relations is again a Lagrangian linear relation ([13], Theorem 8.3) we see $Q \circ P$ is a Lagrangian linear relation. To complete the proof we have to show $Q \circ P$ is positive Lagrangian subspace which is a simple computation. Let $(w, u) \in Q \circ P$, then there is $v \in V_{\mathbb{C}}$ such that $(w, v) \in Q$ and $(v, u) \in P$. Now,

$$i\Omega_{W_{\mathbb{C}} \oplus \bar{U}_{\mathbb{C}}}((\bar{w}, \bar{u}), (w, u)) = i\Omega_{W_{\mathbb{C}} \oplus \bar{V}_{\mathbb{C}}}((\bar{w}, \bar{v}), (w, v)) + i\Omega_{V_{\mathbb{C}} \oplus \bar{U}_{\mathbb{C}}}((\bar{v}, \bar{u}), (v, u))$$

implies $-i\Omega_{W_{\mathbb{C}} \oplus \bar{U}_{\mathbb{C}}}((\bar{w}, \bar{u}), (w, u))$ is positive definite. □

Let H be a finite dimensional real Hilbert space and define $V = V^+ \oplus V^-$ where $V^+ = V^- = H$. Then (V, Ω_V) is a symplectic vector space with the symplectic form the hyperbolic form defined by

$$\Omega((v_1^+, v_1^-), (v_2^+, v_2^-)) = \langle v_1^+, v_2^- \rangle - \langle v_2^+, v_1^- \rangle .$$

Let $V_{\mathbb{C}}$ be the complexification of V . Let $\mathcal{A} : H \rightarrow H$ be a positive operator. Proof of the following lemma is obvious.

Lemma 2.10. *The graph of $i\mathcal{A}$ is a positive Lagrangian subspace of $V_{\mathbb{C}}$.*

More generally, let H_1, H_2 and H_3 be finite dimensional real Hilbert spaces. Let $\mathcal{A}_1 : H_2 \oplus H_1 \rightarrow H_2 \oplus H_1$ and $\mathcal{A}_2 : H_3 \oplus H_2 \rightarrow H_3 \oplus H_2$ be positive operators. We write

$$\mathcal{A}_1 = \begin{bmatrix} A & B \\ B^t & D \end{bmatrix} \text{ and } \mathcal{A}_2 = \begin{bmatrix} K & L \\ L^t & M \end{bmatrix}.$$

Define

$$\widetilde{\mathcal{A}}_1 = \begin{bmatrix} A & B \\ -B^t & -D \end{bmatrix} \text{ and } \widetilde{\mathcal{A}}_2 = \begin{bmatrix} K & L \\ -L^t & -M \end{bmatrix}.$$

It is easy to check the following lemma.

Lemma 2.11. *(i) The subspace L_1 which is the graph of $i\widetilde{\mathcal{A}}_1$ is a positive Lagrangian subspace of*

$$(H_2 \oplus H_2)_{\mathbb{C}} \oplus \overline{((H_1 \oplus H_1)_{\mathbb{C}})}.$$

(ii) The subspace L_2 which is the graph of $i\widetilde{\mathcal{A}}_2$ is a positive Lagrangian subspace of

$$(H_3 \oplus H_3)_{\mathbb{C}} \oplus \overline{(H_2 \oplus H_2)_{\mathbb{C}}}.$$

(iii) $L_2 \circ L_1$ is the graph of $i\widetilde{\mathcal{A}}_2 \circ \mathcal{A}_1$ where

$$\mathcal{A}_2 \circ \mathcal{A}_1 = \begin{bmatrix} K-L(A+M)^{-1}L^t & -L(A+M)^{-1}B \\ -B^t(A+M)^{-1}L^t & D-B^t(A+M)^{-1}B \end{bmatrix}.$$

2.2. Gaussian measures and Bosonic Fock spaces

Here, we recall a few facts about Gaussian measures on a real nuclear vector space V which we use in this paper. Our main reference for Gaussian measure related facts is reference [2]. A well known fundamental fact is that there is a canonical bijection between the set of nondegenerate centered Gaussian measures on V^{\vee} and the set of inner products on V . Given a nondegenerate centered Gaussian measure μ on V^{\vee} the corresponding inner product on V is given by

$$\langle v, w \rangle = \int_{V^{\vee}} (v, F) \cdot (w, F) d\mu(F).$$

where (v, F) and (w, F) are the natural pairings between elements of V and V^{\vee} respectively. Conversely, given an inner product on V we can use Milnos theorem to construct the corresponding Gaussian measure.

We follow the slogan that “almost all” features of a nondegenerate Gaussian measure μ on V^\vee are controlled by a subspace $H(\mu)$ of V^\vee where $H(\mu)$ carries a Hilbert space structure. This Hilbert space is called the *Cameron-Martin* space of the pair (V^\vee, μ) (See [2], page 44 for precise definition of the Cameron-Martin space). We can think of the Cameron-Martin space as the dual of the Hilbert space which is the completion of the V with respect to the inner product that defines the given Gaussian measure.

Example 2.12. Let Y be an object in $d\text{-RBord}$ which means Y a closed oriented $d - 1$ Riemannian manifold. Consider the inner product on $C^\infty(Y)$ given by

$$\langle f, g \rangle_{W^{\frac{1}{2}}(Y)} = \int_Y f(\Delta_Y + m^2)^{\frac{1}{2}} g \, d\text{vol}(Y)$$

Let μ_Y be the corresponding Gaussian measure on $D'(Y)$ the space of distributions on Y . Then the Sobolev space $W^{\frac{1}{2}}(Y)$ is the Cameron-Martin space of μ_Y .

One key feature of the Cameron-Martin space is that it can be used to detect whether two Gaussian measures are mutually absolutely continuous or not. Let (V^\vee, μ) be as before and let $H(\mu)$ be its Cameron-Martin space. Let ν be another Gaussian measure on V^\vee with the Cameron-Martin space $H(\nu)$, then μ and ν are mutually absolutely continuous if and only if $H(\mu) = H(\nu)$ as vector spaces and there is a symmetric Hilbert-Schmidt operator $S : H(\mu) \rightarrow H(\mu)$ such that

$$\langle u, v \rangle_{H(\nu)} = \langle u, v \rangle_{H(\mu)} + \langle Su, v \rangle_{H(\mu)}$$

for all $u, v \in H(\mu)$. See chapter 9 in [2] for details.

Another feature which interests us is that it completely characterizes the square integrable functions in the sense we explain below.

Let V be a nuclear space, μ a Gaussian measure on V^\vee and $H(\mu)$ the Cameron-Martin space of μ . Then there is an isomorphism of Hilbert spaces

$$S : L^2(V^\vee, \mu) \rightarrow \text{Sym}^* H(\mu)^\vee$$

which is known as *Segal-Ito* isomorphism ([9], Chapter 3). Here $\text{Sym}^* H(\mu)^\vee$ is the Bosonic Fock space (See 4.1) of $H(\mu)^\vee$. If V is a finite dimensional real Hilbert space then S is essentially the forward heat operator. Hence we may think of S as the infinite dimensional version of the forward heat operator.

Remark 2.13. The upshot of the previous paragraph is that we do not really need Gaussian measure to describe the space of square integrable functions. It is enough to have the Cameron-Martin space at hand.

We use remark 2.13 frequently explicitly or implicitly, to turn measure theoretic constructions into constructions in Bosonic Fock spaces.

3. Construction of free theories

3.1. Heuristic computation of path integrals and gluing formula for the zeta regularized determinants

We describe a heuristic path integral quantization of a free massive scalar field theory to motivate the construction of the FQFTs in this paper. We recall that we want to make sense of the integrals of the form

$$(3.1) \quad E(\Sigma) = \int_{\mathfrak{F}(\Sigma)} e^{-S(\phi)} D\phi$$

where $S(\phi)$ is given by (1.1).

Σ is closed: First, we assume Σ is closed. Then, using integration by parts, we can write

$$S(\phi) = \frac{1}{2} \int_{\Sigma} \phi(\Delta_{\Sigma} + m^2)\phi \, dvol(\Sigma)$$

This shows $E(\Sigma)$ is a Gaussian. If $\mathfrak{F}(\Sigma)$ were a finite dimensional vector space, then $E(\Sigma)$ would be simply $\det(\Delta_{\Sigma} + m^2)^{-\frac{1}{2}}$. Hence, we need a generalization of determinant for the infinite dimensional case and zeta regularized determinant [6] is one such generalization. This suggests

$$E(\Sigma) = \det_{\zeta}(\Delta_{\Sigma} + m^2)^{-\frac{1}{2}}.$$

Σ has nonempty boundary: Assume that $\partial\Sigma = Y$. Let $\eta \in C^{\infty}(Y)$ and we want to heuristically make sense of integral

$$(3.2) \quad E(\Sigma)(\eta) = \int_{\{\phi \in \mathfrak{F}(\Sigma) : \phi|_Y = \eta\}} e^{-S(\phi)} D\phi$$

First, we need the following lemma

Lemma 3.1. *If ψ is a Helmholtz solution, then*

$$\int_{\Sigma} d\phi \wedge *d\psi + m^2 * \phi\psi = \int_Y \phi D_{\Sigma}\psi|_Y \, d\text{vol}(Y).$$

Proof. Using $d(\phi * d\psi) = d\phi \wedge *d\psi + \phi d * d\psi$, we have

$$\begin{aligned} & \int_{\Sigma} d\phi \wedge *d\psi + m^2 * \phi\psi \\ &= \int_Y (\phi * d\psi)|_Y + \int_{\Sigma} \phi(-d * d\psi + m^2 * \psi) \\ &= \int_Y (\phi * d\psi)|_Y \end{aligned}$$

Now,

$$\int_Y (\phi * d\psi)|_Y = \int_Y \phi D_{\Sigma}\psi|_Y \, d\text{vol}(Y)$$

follows from the proof of Proposition 4.1.54 of [14]. □

Now, we go back to heuristic computation of $E(\Sigma)(\eta)$. Let $\phi \in \mathfrak{F}(\Sigma)$ such that $\phi|_Y = \eta$. Let us write $\phi = \phi_{\eta} + \hat{\phi}$, where ϕ_{η} is the Helmholtz solution with boundary value η and $\hat{\phi}$ vanishes on the boundary. Using Lemma 3.1 and integration by parts, we see that

$$S(\phi) = S(\phi_{\eta}) + S(\hat{\phi}) = \frac{1}{2} \int_Y \eta D_{\Sigma}\eta \, d\text{vol}(Y) + \frac{1}{2} \int_{\Sigma} \hat{\phi}(\Delta_{\Sigma} + m^2)\hat{\phi} \, d\text{vol}(\Sigma)$$

Hence, heuristically,

$$E(\Sigma)(\eta) = e^{-S(\phi_{\eta})} \det_{\zeta}(\Delta_{\Sigma,D} + m^2)^{-\frac{1}{2}}$$

where $\Delta_{\Sigma,D}$ is the Laplacian with the Dirichlet boundary condition.

Gluing: Now assume that $\Sigma = \Sigma_2 \circ \Sigma_1$ in d-RBord where $\Sigma_1 : \emptyset \rightarrow Y$ and $\overline{\Sigma_2} : Y \rightarrow \emptyset$. In particular, Σ is closed. Then,

$$\begin{aligned} & E(\Sigma_2) \circ E(\Sigma_1) \\ &= \det_{\zeta}(\Delta_{\Sigma_1,D} + m^2)^{-\frac{1}{2}} \det_{\zeta}(\Delta_{\Sigma_2,D} + m^2)^{-\frac{1}{2}} \int_{\eta \in C^{\infty}(Y)} e^{-\frac{1}{2} \langle (D_{\Sigma_2} + D_{\Sigma_1})\eta, \eta \rangle} D\eta \\ &= \det_{\zeta}(\Delta_{\Sigma_1,D} + m^2)^{-\frac{1}{2}} \det_{\zeta}(\Delta_{\Sigma_2,D} + m^2)^{-\frac{1}{2}} \det_{\zeta}(D_{\Sigma_2} + D_{\Sigma_1})^{-\frac{1}{2}}. \end{aligned}$$

In order to have a FQFT, we want $E(\Sigma) = E(\Sigma_2) \circ E(\Sigma_1)$. Hence, we expect to have a gluing formula for the zeta regularized determinants. In

other words, if we compute path integrals heuristically and expect a FQFT, then it predicts the gluing formula for the zeta regularized determinants.

Fortunately, there is a gluing formula for the zeta regularized determinants ([3], Theorem A or [12], Corollary 1.3) when d is even. Hence, heuristic computation of path integrals leads to $E(\Sigma) = E(\Sigma_2) \circ E(\Sigma_1)$.

In summary, heuristic computations of path integrals together with the gluing formula for the zeta regularized determinants suggests the existence of functorial quantum field theories. In the rest of the paper we show this is indeed the case. In particular, when $\partial\Sigma = Y$, we want to give a precise mathematical meaning to $E(\Sigma)$ from above. We want to think of $E(\Sigma)$ as a vector in a Hilbert space associated to Y , and we will see in our approach, that we must modify $E(\Sigma)$ from above.

3.2. A strategy for construction of free theories

In [20], Segal suggests free theories can be constructed composing two functors $\text{d-RBord} \rightarrow T_{pol}$ and $T_{pol} \rightarrow \mathbf{Hilb}$. We heuristically follow this path. Here T_{pol} is the “category” of “polarized” symplectic vector spaces whose objects are polarized symplectic vector spaces (V, J) , where V is a symplectic vector space and J is a polarization. A morphism $(V_0, J_0) \rightarrow (V_1, J_1)$ is a polarization of $V_1 \oplus \overline{V_0}$ which is “close” to $J_1 - J_0$. For an object Y in d-RBord the datum $(\Phi(Y), \sqrt{\Delta_Y + m^2})$, can thought of as a “polarized” symplectic vector space. The functor $\text{d-RBord} \rightarrow T_{pol}$ assigns to an object Y the polarized symplectic vector space $(\Phi(Y), \sqrt{\Delta_Y + m^2})$. If $\Sigma : \emptyset \rightarrow Y$ is a morphism in d-RBord , then the operator D_Σ can be interpreted as a morphism in T_{pol} . The other functor $T_{pol} \rightarrow \mathbf{Hilb}$ can be thought of as passing to the quantum theory.

3.3. Construction of Hilbert space of states

Let Y be an object in d-RBord . There are at least two equivalent ways to construct the Hilbert space associated to Y .

We recall, to Y , one can assign the symplectic vector space $(\Phi(Y), \omega_{\Phi(Y)})$ where $\Phi(Y) = C^\infty(Y) \oplus C^\infty(Y)$ and $\omega_{\Phi(Y)}$ is the hyperbolic form induced by the L^2 inner product on $C^\infty(Y)$. The underlying vector space of $\Phi(Y)$ can be thought of as the Cauchy data for the classical solutions associated to the action functional (1.1) on a very thin cylinder around Y . We need an extra data namely a polarization so that we can use geometric quantization to construct a Hilbert space. We may think of graph of $i\sqrt{\Delta_Y + m^2}$ as a

positive polarization (see Section 9.5 in [16]) on $\Phi(Y)$. Now, we can assign to Y the Hilbert space given by geometric quantization.

Since we are interested in explicit computations, we use the ideas from constructive quantum field theory ([8], [15]) to construct the Hilbert space of states.

To an object Y of $d\text{-RBord}$, we assign the Hilbert space of states

$$E(Y) = \text{Sym}^* W^{\frac{1}{2}}(Y)^\vee.$$

Recall from Section 2.2 that $E(Y)$ is the L^2 space of the Gaussian measure on $D'(Y)$ whose Cameron-Martin space is $W^{\frac{1}{2}}(Y)$.

It is clear from the construction $E(\phi) = \mathbb{R}$, $E(\overline{Y}) = E(Y)^\vee$, and $E(Y_1 \sqcup Y_2) = E(Y_1) \otimes E(Y_2)$ where \otimes is the Hilbert space tensor product.

Remark 3.2. Following [15], $E(Y)$ can be defined as the the space of half densities of the measure class the Gaussian measure μ on $D'(Y)$ with the Cameron-Martin space $W^{\frac{1}{2}}(Y)$.

Remark 3.3. The complexification of $E(Y)$ is essentially an irreducible unitary representation of the “standard” Heisenberg group of $\Phi(Y)$ as explained in Section 9.5 in [16] or Section 5.4 in [7].

3.4. Construction of operators

Let $\Sigma : \emptyset \rightarrow Y$ be a morphism in $d\text{-RBord}$. We want to construct a vector $E(\Sigma) \in E(Y)$. Let us first explain the measure theoretic idea behind the construction. We note there is a Gaussian measure μ_Y on the space of distributions $D'(Y)$ on Y whose Cameron-Martin space is the Sobolev space $W^{\frac{1}{2}}(Y)$. Recall that D_Σ is a positive operator on Y and hence it defines an inner product on $C^\infty(Y)$ given by

$$\langle f, g \rangle_1 = \int_Y f D_\Sigma g \, d\text{vol}(Y).$$

Let ν_Y be the Gaussian measure on $D'(Y)$ corresponding to the inner product \langle, \rangle_1 .

By our assumption Σ has product metric near Y . In this case, it turns out that $D_\Sigma - (\Delta_Y + m^2)^{\frac{1}{2}}$ is a smoothing operator (Theorem 2.1, [12]). As a consequence of this fact we observe the measures μ_Y and ν_Y are mutually absolutely continuous. Now the idea is to define the vector $E(\Sigma)$ is a suitable

multiple of the square root of the Radon-Nikodym derivative ν_Y with respect to μ_Y . We will give this construction in the Fock space setting.

We recall $\alpha_\Sigma = D_\Sigma(\Delta_Y + m^2)^{-\frac{1}{2}}$ and α_Σ is a continuous positive operator on $W^{\frac{1}{2}}(Y)$. Moreover,

Lemma 3.4. $\alpha_\Sigma - I$ is trace class operator on $W^{\frac{1}{2}}(Y)$.

Proof. Theorem 2.1 in [12] implies $D_\Sigma - (\Delta_Y + m^2)^{\frac{1}{2}}$ is a smoothing operator on Y . This means

$$\alpha_\Sigma = I + R$$

where R is a smoothing operator. Since R is smoothing the unique extension of R on $W^{\frac{1}{2}}(Y)$ is trace class. □

Let $C(\alpha_\Sigma) = (I - \alpha_\Sigma)(I + \alpha_\Sigma)^{-1}$ be the Cayley Transform of α_Σ (see Definition 4.6). We observe $C(\alpha_\Sigma)$ is Hilbert-Schmidt symmetric operator on $W^{\frac{1}{2}}(Y)$ and $\|C(\alpha_\Sigma)\| < 1$. From 4.1, we have

$$\mathcal{E}(C(\alpha_\Sigma)) = \text{Exp} \left(\frac{1}{2} C(\alpha_\Sigma) \right) \in \text{Sym}^* W^{\frac{1}{2}}(Y)^\vee.$$

We will use $\mathcal{E}(C(\alpha_\Sigma))$ to define $E(\Sigma)$.

Remark 3.5. The vector $\frac{\mathcal{E}(C(\alpha_\Sigma))}{\|\mathcal{E}(C(\alpha_\Sigma))\|}$ is the image of the square root of the Radon-Nikodym derivative of ν_Y with respect to μ_Y under the *Segal-Ito* isomorphism ([11], Theorem 3.2.24).

More generally if $\Sigma_1 : Y_1 \rightarrow Y_2$ is a morphism in d-RBord , then we interpret Σ as a morphism from ϕ to $Y_2 \sqcup \overline{Y_1}$. Then

$$\mathcal{E}(C(\alpha_{\Sigma_1})) \in E(Y_1)^\vee \otimes E(Y_2).$$

We can identify $\mathcal{E}(C(\alpha_{\Sigma_1}))$ with a Hilbert-Schmidt operator $E(Y_1) \rightarrow E(Y_2)$ which will be again denoted by $\mathcal{E}(C(\alpha_{\Sigma_1}))$.

Proposition 3.6. Let H_1, H_2 and H_2 be real separable Hilbert spaces. Let \mathcal{A}_1 and \mathcal{A}_2 be continuous positive operators respectively on $H_2 \oplus H_1$ and

$H_3 \oplus H_2$ such that $\mathcal{A}_1 - I$ and $\mathcal{A}_2 - I$ are trace class operators. Then,

$$(3.3) \quad \mathcal{E}(C(\mathcal{A}_2)) \circ \mathcal{E}(C(\mathcal{A}_1)) = c(\mathcal{A}_2, \mathcal{A}_1)\mathcal{E}(C(\mathcal{A}_2 \circ \mathcal{A}_1))$$

where

$$c(\mathcal{A}_2, \mathcal{A}_1) = \frac{\det(\mathcal{A}_1)^{\frac{1}{4}} \cdot \det(\mathcal{A}_2)^{\frac{1}{4}}}{\det(\frac{A+M}{2})^{\frac{1}{2}} \det(\mathcal{A}_2 \circ \mathcal{A}_1)^{\frac{1}{4}}} \frac{\|\mathcal{E}(C(\mathcal{A}_1))\| \cdot \|\mathcal{E}(C(\mathcal{A}_2))\|}{\|\mathcal{E}(C(\mathcal{A}_2 \circ \mathcal{A}_1))\|}$$

and the norms are taken in the Fock spaces. Here,

$$\mathcal{A}_1 = \begin{bmatrix} A & B \\ B^t & D \end{bmatrix}, \quad \mathcal{A}_2 = \begin{bmatrix} K & L \\ L^t & M \end{bmatrix} \quad \text{and} \quad \mathcal{A}_2 \circ \mathcal{A}_1 = \begin{bmatrix} K-L(A+M)^{-1}L^t & -L(A+M)^{-1}B \\ -B^t(A+M)^{-1}L^t & D-B^t(A+M)^{-1}B \end{bmatrix}.$$

Proof. First of all we note that $\mathcal{A}_1 - I$ and $\mathcal{A}_2 - I$ trace class implies $\det(\frac{A+M}{2})$ and $\det(\mathcal{A}_2 \circ \mathcal{A}_1)$ exist. Also, we note \mathcal{E} , C and the composition map \circ are continuous. Hence, it is sufficient to show the equation (3.3) holds for \mathcal{A}_1 and \mathcal{A}_2 such that $\mathcal{A}_1 - I$ and $\mathcal{A}_2 - I$ are finite rank operators. This case is done in the Proposition 4.8. □

Definition 3.7. A projective representation T of a category \mathcal{C} in **Hilb** assigns to an object C in \mathcal{C} a Hilbert Space $T(C)$, and to a morphism $P : C \rightarrow D$ a continuous linear operator $T(P)$ such that for any pair of morphisms $P : C \rightarrow D$ and $Q : D \rightarrow E$ we have

$$T(Q \circ P) = \lambda(Q, P)T(Q) \circ T(P)$$

where $\lambda(Q, P)$ is a nonzero complex number.

Theorem 3.8. *The assignment, $Y \rightarrow E(Y)$ where Y is an object in $d\text{-RBord}$ and $\Sigma \rightarrow \mathcal{E}(C(\alpha_\Sigma))$ where $\Sigma : \emptyset \rightarrow Y$ is a morphism in $d\text{-RBord}$, defines a projective representation of $d\text{-RBord}$ in **Hilb**.*

Proof. Let $\Sigma_1 : Y_1 \rightarrow Y_2$ and $\Sigma_2 : Y_2 \rightarrow Y_3$ be morphisms in $d\text{-RBord}$. From corollary, 2.6 we recall

$$\alpha_{\Sigma_2} \circ \alpha_{\Sigma_1} = \alpha_{\Sigma_2 \circ \Sigma_1}.$$

Now applying Proposition 3.6 to α_{Σ_2} , α_{Σ_1} and $\alpha_{\Sigma_2 \circ \Sigma_1}$, we get

$$\mathcal{E}(C(\alpha_{\Sigma_2})) \circ \mathcal{E}(C(\alpha_{\Sigma_1})) = c(\alpha_{\Sigma_2}, \alpha_{\Sigma_1})\mathcal{E}(C(\alpha_{\Sigma_2 \circ \Sigma_1})),$$

where

$$(3.4) \quad c(\alpha_{\Sigma_2}, \alpha_{\Sigma_1}) = \frac{\det(\alpha_{\Sigma_1})^{\frac{1}{4}} \cdot \det(\alpha_{\Sigma_1})^{\frac{1}{4}}}{\det\left(\frac{(A+M)(\Delta_{Y_2} + m^2)^{-\frac{1}{2}}}{2}\right)^{\frac{1}{2}} \det(\mathcal{E}(C(\alpha_{\Sigma_2 \circ \Sigma_1})))^{\frac{1}{4}}} \frac{\|\mathcal{E}(C(\alpha_{\Sigma_2}))\| \cdot \|\mathcal{E}(C(\alpha_{\Sigma_1}))\|}{\|\mathcal{E}(C(\alpha_{\Sigma_2 \circ \Sigma_1}))\|}$$

which is nonzero. □

Next, the idea is to “deprojectivize” the projective representation constructed above to construct Free Scalar Theory.

Let $\Sigma : \emptyset \rightarrow Y$ be a morphism in d-RBord. In order to deprojective the projective representation of Theorem 3.8, we must rescale $\mathcal{E}(C(\alpha_\Sigma))$ properly. Hence, we define

$$E(\Sigma) = \frac{1}{\det_\zeta(\Delta_{\Sigma_1, D} + m^2)^{\frac{1}{2}} \cdot \det_\zeta(2D_\Sigma)^{\frac{1}{4}}} \cdot \frac{\mathcal{E}(C(\alpha_\Sigma))}{\|\mathcal{E}(C(\alpha_\Sigma))\|}$$

Here the notation $\Delta_{\Sigma, D}$ means the operator Δ_Σ with the Dirichlet boundary condition and \det_ζ is for the zeta regularized determinant ([6]).

More generally, if

$$\Sigma : Y_1 \rightarrow Y_2$$

is a morphism in d-RBord, we think of

$$\Sigma : \emptyset \rightarrow \overline{Y_1} \sqcup Y_2$$

and then define $E(\Sigma)$ as before. This means $E(\Sigma) \in E(Y_1)^\vee \otimes E(Y_2)$, and consequently we get a Hilbert-Schmidt operator

$$E(\Sigma) : E(Y_1) \rightarrow E(Y_2).$$

In fact, $E(\Sigma)$ will be a trace class operator as a consequence of 3.10, because we can always write Σ as a composition of two morphisms in d-RBord. Also, by construction we have

$$E(\sqcup_i \Sigma_i) = \otimes_i E(\Sigma_i).$$

Lemma 3.9. *Let $\Sigma : \emptyset \rightarrow Y$ then*

$$\det(\alpha_\Sigma) = \frac{\det_\zeta(2D_\Sigma)}{\det_\zeta(2(\Delta_Y + m^2)^{\frac{1}{2}})}.$$

Proof. We note

$$2(\Delta_Y + m^2)^{\frac{1}{2}}\alpha_\Sigma = 2D_\Sigma.$$

Recall that $\det(\alpha_\Sigma)$ exists as $\alpha_\Sigma - I$ is trace class operator. Using a well known fact (If $A - I$ is trace class then $\det_\zeta(AB) = \det(A) \cdot \det_\zeta B$ Proposition 2.20 [18]), we get

$$\det_\zeta(2(\Delta_Y + m^2)^{\frac{1}{2}}\alpha_\Sigma) = \det_\zeta(2(\Delta_Y + m^2)^{\frac{1}{2}}) \det(\alpha_\Sigma).$$

This gives the desired result $\det(\alpha_\Sigma) = \frac{\det_\zeta(2D_\Sigma)}{\det_\zeta(2(\Delta_Y + m^2)^{\frac{1}{2}})}$. □

Now, we prove the gluing law for the operators.

Theorem 3.10. *Let $\Sigma_1 : Y_1 \rightarrow Y_2$ and $\Sigma_2 : Y_2 \rightarrow Y_3$ be two morphisms in d -RBord.*

(i) *There exists a non zero constant $C(\Sigma_2, \Sigma_1)$ such that*

$$E(\Sigma_2) \circ E(\Sigma_1) = C(\Sigma_2, \Sigma_1)E(\Sigma_2 \circ \Sigma_1).$$

(ii) *When d is even, $C(\Sigma_2, \Sigma_1) = 1$.*

Proof. For $\Sigma_1 : Y_1 \rightarrow Y_2$, we write

$$D_{\Sigma_1} = \begin{bmatrix} A & B \\ B^t & D \end{bmatrix} \text{ where } D_{\Sigma_1} : C^\infty(Y_2) \oplus C^\infty(Y_1) \rightarrow C^\infty(Y_2) \oplus C^\infty(Y_1).$$

Similarly for $\Sigma_2 : Y_2 \rightarrow Y_3$, we write

$$D_{\Sigma_2} = \begin{bmatrix} K & L \\ L^t & M \end{bmatrix} \text{ where } D_{\Sigma_2} : C^\infty(Y_3) \oplus C^\infty(Y_2) \rightarrow C^\infty(Y_3) \oplus C^\infty(Y_2).$$

We note

$$\begin{aligned} E(\Sigma_2) \circ E(\Sigma_1) &= \frac{1}{\det_\zeta(\Delta_{\Sigma_1, D} + m^2)^{\frac{1}{2}} \det_\zeta(2D_{\Sigma_1})^{\frac{1}{4}}} \\ &\cdot \frac{1}{\det_\zeta(\Delta_{\Sigma_2, D} + m^2)^{\frac{1}{2}} \det_\zeta(2D_{\Sigma_2})^{\frac{1}{4}}} \\ &\cdot \frac{1}{\|\mathcal{E}(C(\alpha_{\Sigma_1}))\|} \cdot \frac{1}{\|\mathcal{E}(C(\alpha_{\Sigma_2}))\|} \mathcal{E}(C(\alpha_{\Sigma_2})) \circ \mathcal{E}(C(\alpha_{\Sigma_1})). \end{aligned}$$

By proof of Theorem 3.8, we have

$$\begin{aligned} & \frac{1}{\|\mathcal{E}(C(\alpha_{\Sigma_1}))\|} \cdot \frac{1}{\|\mathcal{E}(C(\alpha_{\Sigma_1}))\|} \mathcal{E}(C(\alpha_{\Sigma_2})) \circ \mathcal{E}(C(\alpha_{\Sigma_1})) \\ &= \frac{\det(\alpha_{\Sigma_1})^{\frac{1}{4}} \det(\alpha_{\Sigma_2})^{\frac{1}{4}}}{\det(\alpha_{\Sigma_2 \circ \Sigma_1})^{\frac{1}{4}} \det(\frac{1}{2}(\Delta_{Y_2} + m^2)^{-\frac{1}{2}} D_{\Sigma_1, \Sigma_2})^{\frac{1}{2}}} \cdot \frac{\mathcal{E}(C(\alpha_{\Sigma_2 \circ \Sigma_1}))}{\|\mathcal{E}(C(\alpha_{\Sigma_2 \circ \Sigma_1}))\|} \end{aligned}$$

where $D_{\Sigma_1, \Sigma_2} = A + M$.

Let us recall from the Lemma 3.9

$$\begin{aligned} \det(\alpha_{\Sigma_1}) &= \frac{\det_{\zeta}(2D_{\Sigma_1})}{\det_{\zeta}(2(\Delta_{Y_1} + m^2)^{\frac{1}{2}}) \det_{\zeta}(2(\Delta_{Y_2} + m^2)^{\frac{1}{2}})} \\ \det(\alpha_{\Sigma_2}) &= \frac{\det_{\zeta}(2D_{\Sigma_2})}{\det_{\zeta}(2(\Delta_{Y_2} + m^2)^{\frac{1}{2}}) \cdot \det_{\zeta}(2(\Delta_{Y_3} + m^2)^{\frac{1}{2}})}, \\ \det(\alpha_{\Sigma_2 \circ \Sigma_1}) &= \frac{\det_{\zeta}(2D_{\Sigma_2 \circ \Sigma_1})}{\det_{\zeta}(2(\Delta_{Y_1} + m^2)^{\frac{1}{2}}) \det_{\zeta}(2(\Delta_{Y_3} + m^2)^{\frac{1}{2}})}, \end{aligned}$$

and using the argument of the same lemma

$$\det(\frac{1}{2}(\Delta_{Y_2} + m^2)^{-\frac{1}{2}} D_{\Sigma_1, \Sigma_2}) = \frac{\det_{\zeta}(D_{\Sigma_1, \Sigma_2})}{\det_{\zeta}(2(\Delta_{Y_1} + m^2)^{\frac{1}{2}})}.$$

Using these relations, we see

$$\frac{\det(\alpha_{\Sigma_1})^{\frac{1}{4}} \cdot \det(\alpha_{\Sigma_2})^{\frac{1}{4}}}{\det(\frac{1}{2}(\Delta_{Y_2} + m^2)^{-\frac{1}{2}} D_{\Sigma_1, \Sigma_2})^{\frac{1}{2}} \cdot \det(\alpha_{\Sigma_2 \circ \Sigma_1})^{\frac{1}{4}}} = \frac{\det_{\zeta}(2D_{\Sigma_1})^{\frac{1}{4}} \cdot \det_{\zeta}(2D_{\Sigma_2})^{\frac{1}{4}}}{\det_{\zeta}(D_{\Sigma_1, \Sigma_2})^{\frac{1}{2}} \cdot \det_{\zeta}(2D_{\Sigma_3})^{\frac{1}{4}}}.$$

This gives

$$\begin{aligned} & \frac{1}{\|\mathcal{E}(C(\alpha_{\Sigma_1}))\|} \cdot \frac{1}{\|\mathcal{E}(C(\alpha_{\Sigma_1}))\|} \mathcal{E}(C(\alpha_{\Sigma_2})) \circ \mathcal{E}(C(\alpha_{\Sigma_1})) \\ &= \frac{\det_{\zeta}(2D_{\Sigma_1})^{\frac{1}{4}} \cdot \det_{\zeta}(2D_{\Sigma_2})^{\frac{1}{4}}}{\det_{\zeta}(D_{\Sigma_1, \Sigma_2})^{\frac{1}{2}} \cdot \det_{\zeta}(2D_{\Sigma_3})^{\frac{1}{4}}} \frac{\mathcal{E}(C(\alpha_{\Sigma_2 \circ \Sigma_1}))}{\|\mathcal{E}(C(\alpha_{\Sigma_2 \circ \Sigma_1}))\|}. \end{aligned}$$

Hence,

$$\begin{aligned} & E(\Sigma_2) \circ E(\Sigma_1) \\ &= \frac{1}{\det_{\zeta}(\Delta_{\Sigma_1, D} + m^2)^{\frac{1}{2}} \cdot \det_{\zeta}(\Delta_{\Sigma_2, D} + m^2)^{\frac{1}{2}} \cdot \det_{\zeta}(D_{\Sigma_1, \Sigma_2})^{\frac{1}{2}} \cdot \det_{\zeta}(2D_{\Sigma_3})^{\frac{1}{4}}} \\ & \cdot \frac{\mathcal{E}(C(\alpha_{\Sigma_2 \circ \Sigma_1}))}{\|\mathcal{E}(C(\alpha_{\Sigma_2 \circ \Sigma_1}))\|}. \end{aligned}$$

Using the gluing formula for the zeta regularized determinants ([3], Theorem A or [12], Corollary 1.3), we have

$$\det_{\zeta}(\Delta_{\Sigma_2 \circ \Sigma_1, D} + m^2) = K(\Sigma_1, \Sigma_2) \det_{\zeta}(\Delta_{\Sigma_1, D} + m^2) \cdot \det_{\zeta}(\Delta_{\Sigma_1, D} + m^2) \cdot \det_{\zeta}(D_{\Sigma_1, \Sigma_2})$$

for some nonzero constant $K(\Sigma_1, \Sigma_2)$. Hence,

$$E(\Sigma_2) \circ E(\Sigma_1) = \frac{1}{\sqrt{K(\Sigma_1, \Sigma_2)}} E(\Sigma_2 \circ \Sigma_1).$$

Furthermore, when d is even $K(\Sigma_1, \Sigma_2) = 1$ ([12], Remark 1 after Corollary 1.3). □

Corollary 3.11. *Suppose that Σ is a closed d -dimensional oriented Riemannian manifold and d is even. Assume that $\Sigma = \Sigma_2 \circ \Sigma_1$ in d -RBord where $\Sigma_1 : \emptyset \rightarrow Y$ and $\Sigma_2 : Y \rightarrow \emptyset$. Then*

$$E(\Sigma) = \det_{\zeta}(\Delta_{\Sigma} + m^2)^{-\frac{1}{2}}.$$

As suggested by the previous corollary, we complete the construction of the functorial free field theory by assigning to a closed oriented Riemannian manifold Σ the number

$$E(\Sigma) = \det_{\zeta}(\Delta_{\Sigma} + m^2)^{-\frac{1}{2}}.$$

3.5. More examples of free theory

We describe an immediate generalization of previous examples which is to quantize a more general classical field theory. We take $\mathfrak{F}(\Sigma)$ to be sections of a real finite dimensional vector bundle F on Σ . Assume that F comes with a metric g^F and a compatible connection ∇^F . Let $d_{\nabla^F} : \Omega^*(M, F) \rightarrow \Omega^{*+1}(M, F)$ be the twisted deRham operator acting on F -valued forms. We are interested in $\Omega^0(M, F)$ where the action functional is defined as

$$S(\phi) = \frac{1}{2} \int_{\Sigma} (d_{\nabla^F} \phi \wedge *d_{\nabla^F} \phi)_{g^F} + m^2(\phi \wedge *\phi)_{g^F}$$

where $*$ is the Hodge star operator.

We denote the Hodge Laplacian acting on sections of F by Δ_{Σ} i.e. $\Delta_{\Sigma} = d_{\nabla^F}^* \circ d_{\nabla^F}$ where $d_{\nabla^F}^*$ is the formal adjoint of d_{∇^F} .

We consider the semicategory \mathcal{C}_d which is a straightforward generalization of d-RBord. An object in \mathcal{C}_d is a quadruple (Y, F, g^F, ∇^F) . Here Y is a $d - 1$ dimensional closed oriented Riemannian manifold, F is a real vector bundle over Y of finite rank, g^F a metric on F and ∇^F is a flat metric connection. Let $(Y_i, F_i, g^{F_i}, \nabla^{F_i})$, $i = 1, 2$ be two objects of \mathcal{C}_d . A morphism from $(Y_1, F_1, g^{F_1}, \nabla^{F_1})$ to $(Y_2, F_2, g^{F_2}, \nabla^{F_2})$ is a quadruple $(\Sigma, F, g^F, \nabla^F)$ where $\Sigma : Y_1 \rightarrow Y_2$ is a bordism in d-RBord, F is a vector bundle on Σ with $F|_{Y_i} = F_i$, g^F is a metric on F with $g^F|_{F_i} = g^{F_i}$ and ∇^F is a flat metric connection compatible with the connections ∇^{F_i} on F_i . We assume the geometry of all the metrics and the connection ∇^F is of *product type* near Y_i . We define the composition of morphisms to be the one given by gluing.

Let (Y, F, g^F, ∇^F) be an object of \mathcal{C}_d . We define $W^{\frac{1}{2}}(Y, F)$ to be completion of the space of smooth sections $C^\infty(Y, F)$ with respect to the inner product

$$(\phi, \psi) = \int_Y (\phi, (\Delta_F + m^2)^{\frac{1}{2}} \psi)_{g^F} \, dvol(Y).$$

Let $(\Sigma, F, g^F, \nabla^F)$ be a morphism in \mathcal{C}_d from ϕ to (Y, F, g^F, ∇^F) . We define D_Σ , α_Σ and $\mathcal{E}(C(\alpha_\Sigma))$ as before (see Section 2 and Section 3).

To an object (Y, F, g^F, ∇^F) in \mathcal{C}_d , assign the Hilbert space

$$E((Y, F, g^F, \nabla^F)) = \text{Sym}^* W^{\frac{1}{2}}(Y, F)^\vee$$

and to a morphism $(\Sigma, F, g^F, \nabla^F)$ from ϕ to (Y, F, g^F, ∇^F) assign the vector

$$\frac{1}{\det_\zeta(\Delta_{\Sigma_1, D} + m^2)^{\frac{1}{2}} \cdot \det_\zeta(2D_\Sigma)^{\frac{1}{4}}} \cdot \frac{\mathcal{E}(C(\alpha_\Sigma))}{\|\mathcal{E}(C(\alpha_\Sigma))\|} \in E((Y, F, g^F, \nabla^F)).$$

This assignment defines a projective Functorial QFT for all d and Functorial QFT when d is even as above.

In particular, if we take the trivial vector bundle \mathbb{R}^N , then we get a Functorial QFT with target \mathbb{R}^N .

4. Appendix

4.1. Bosonic Fock space

Note 4.1. *In what follows we assume the Hilbert spaces in consideration are separable.*

Definition 4.2. The *Bosonic Fock* space of a Hilbert space H is the Hilbert space direct sum

$$\bigoplus_{n=0}^{\infty} \text{Sym}^n(H) = \left\{ (\alpha_n)_{n=0}^{\infty} : \alpha_n \in \text{Sym}^n H \text{ with } \sum_{n=0}^{\infty} \|\alpha_n\|^2 < \infty \right\}$$

where $\text{Sym}^n(H)$ is the closed subspace of $H^{\otimes n}$ that is invariant under the action of permutation group S_n . We note $\text{Sym}^n(H)$ is the closure of the subspace spanned by vectors of the form

$$h_1 \otimes_s \cdots \otimes_s h_n = \frac{1}{\sqrt{n!}} \sum_{\sigma \in S_n} h_{\sigma(1)} \otimes \cdots \otimes h_{\sigma(n)}.$$

We use $\text{Sym}^*(H)$ to denote the Bosonic Fock space of H . We also use the notation $h_1 \cdot h_2 \cdots h_n$ for $h_1 \otimes_s \cdots \otimes_s h_n$.

Let H be a real Hilbert space and let $A : H \rightarrow H$ be a symmetric Hilbert Schmidt linear operator. Define a map $S(A) : H \times H \rightarrow \mathbb{R}$ by

$$S(A)(u \otimes v) = \langle Au, v \rangle .$$

Since A is symmetric and Hilbert Schmidt, we see $S(A) \in \text{Sym}^2(H^\vee)$. Moreover, if $O : H \rightarrow H$ is an orthogonal operator then $S(OAO^{-1}) = S(A)$. By abuse of notation we write A for $S(A)$.

Definition 4.3. Let H be a finite dimensional real Hilbert space and $A : H \rightarrow H$ be a symmetric linear operator. We define

$$\mathcal{E}(A) = \text{Exp} \left(\frac{1}{2} S(A) \right)$$

where

$$\text{Exp}(\alpha) = \sum_{n=0}^{\infty} \frac{\alpha^n}{n!}$$

for all $\alpha \in \text{Sym}^k H^\vee$.

Proposition 4.4. Assume that H is finite dimensional and $A : H \rightarrow H$ a symmetric linear operator. Then $\mathcal{E}(A) \in \text{Sym}^*(H^\vee)$ if and only if $\|A\| < 1$. If $\|A\| < 1$, then

$$\|\mathcal{E}(A)\|^2 = \frac{1}{\sqrt{\det(I - A^2)}}.$$

Proposition 4.4 is a direct consequence of the following elementary lemma.

Lemma 4.5. *Assume that H is one dimensional and $x \in H^\vee$ such that $\|x\| = 1$. Let $a \in \mathbb{R}$. Then $\text{Exp}(\frac{1}{2}ax^2) \in \text{Sym}^*(H^\vee)$ if and only if $|a| < 1$. If $|a| < 1$ then $\|\text{Exp}(\frac{1}{2}ax^2)\|^2 = \frac{1}{\sqrt{1-a^2}}$.*

Proof. We note $\|x^n\|^2 = n!$. Hence,

$$\|\text{Exp}(\frac{1}{2}ax^2)\|^2 = \sum_{n=0}^{\infty} \frac{a^{2n}(2n)!}{2^{2n}(n!)^2}.$$

We know that the series $\sum_{n=0}^{\infty} \frac{a^{2n}(2n)!}{2^{2n}(n!)^2}$ converges if and only if $|a| < 1$ and when it converges the sum is $\frac{1}{\sqrt{1-a^2}}$. □

Assume that H is finite dimension real Hilbert space and let A and B be symmetric linear operators on H of norm smaller than one. A direct computation shows

$$\langle \mathcal{E}(A), \mathcal{E}(B) \rangle = \det(1 - AB)^{-\frac{1}{2}}$$

and

$$\|\mathcal{E}(A) - \mathcal{E}(B)\|^2 = \det(1 - A^2)^{-\frac{1}{2}} + \det(1 - B^2)^{-\frac{1}{2}} - 2(\det(1 - AB)^{-\frac{1}{2}}).$$

As an immediate consequence of this calculation, we get the assignment $A \mapsto \mathcal{E}(A)$ is continuous on the space of symmetric operators in H which have norm smaller than one.

Now we move on to the infinite dimensional case. Let $\mathcal{Z}(H)$ denote the space of symmetric Hilbert-Schmidt linear operators on $A : H \rightarrow H$ such that the operator norm of A is smaller than one. We note $\mathcal{Z}(H)$ is an open subset of the space of Hilbert-Schmidt operators on H . Now, using a well known fact that the space of finite rank operators is dense in the space of Hilbert Schmidt operators [7], we have a continuous map $\mathcal{E} : \mathcal{Z}(H) \rightarrow \text{Sym}^* H^\vee$ given by

$$\mathcal{E}(A) = \text{Exp} \left(\frac{1}{2}S(A) \right).$$

It also follows that

$$\langle \mathcal{E}(A), \mathcal{E}(B) \rangle = \det(1 - AB)^{-\frac{1}{2}}$$

for all $A, B \in \mathcal{Z}(H)$.

Consider the map $C : (0, \infty) \rightarrow (-1, 1)$ given by $x \mapsto (1 - x)(1 + x)^{-1}$. It is easy to see that C is a homeomorphism. Now we use C to define the *Cayley transform* of an operator.

Definition 4.6. Let H be a real Hilbert space and let $A : H \rightarrow H$ be a continuous positive operator. The *Cayley transform* of A is defined by $C(A) = (I - A)(I + A)^{-1}$. We note $C(A) \in \mathcal{Z}(H)$ whenever A is a positive operator on H where $\mathcal{Z}(H)$ is the set of symmetric operators on H with norm smaller than 1.

Let H be a real Hilbert space. The set $\mathbb{S}(H)$ denotes the set of continuous positive linear operators on H such that $A - I$ is Hilbert-Schmidt. We note each $A \in \mathbb{S}(H)$ has a bounded inverse.

There is an obvious map from $\mathbb{S}(H)$ to the Hilbert space of Hilbert-Schmidt operators on H given by $A \mapsto A - I$. We give $\mathbb{S}(H)$ the weakest topology such that the map $A \mapsto A - I$ is a continuous map. In other words, A_n converges to A in $\mathbb{S}(H)$ if and only if $A_n - I$ converges to $A - I$ with respect to the Hilbert Schmidt norm. The proof of the following lemma is obvious.

Lemma 4.7. $C : \mathbb{S}(H) \rightarrow \mathcal{Z}(H)$ is a homeomorphism.

Assume that H_1, H_2 and H_3 are finite dimensional real Hilbert spaces. Let $\mathcal{A}_1 \in \mathbb{S}(H_2 \oplus H_1)$ and $\mathcal{A}_2 \in \mathbb{S}(H_3 \oplus H_2)$. We write

$$\mathcal{A}_1 = \begin{bmatrix} A & B \\ B^t & D \end{bmatrix} \text{ and } \mathcal{A}_2 = \begin{bmatrix} K & L \\ L^t & M \end{bmatrix}$$

and define

$$(4.1) \quad \mathcal{A}_2 \circ \mathcal{A}_1 = \begin{bmatrix} K - L(A+M)^{-1}L^t & -L(A+M)^{-1}B \\ -B^t(A+M)^{-1}L^t & D - B^t(A+M)^{-1}B \end{bmatrix}.$$

In fact $\mathcal{A}_2 \circ \mathcal{A}_1 \in \mathbb{S}(H_3 \oplus H_1)$ and the operation \circ is associative and continuous (Lemma 3.2.11 in [11]).

Now using the Cayley transform we can define $\circ : \mathcal{Z}(H_3 \oplus H_2) \times \mathcal{Z}(H_2 \oplus H_1) \rightarrow \mathcal{Z}(H_3 \oplus H_1)$ that makes the following diagram commutative:

$$\begin{array}{ccc} \mathbb{S}(H_3 \oplus H_2) \times \mathbb{S}(H_2 \oplus H_1) & \xrightarrow{\circ} & \mathbb{S}(H_3 \oplus H_1) \\ \downarrow & & \downarrow \\ \mathcal{Z}(H_3 \oplus H_2) \times \mathcal{Z}(H_2 \oplus H_1) & \xrightarrow{\circ} & \mathcal{Z}(H_3 \oplus H_1) \end{array}$$

Let $\mathcal{A}_1 \in \mathbb{S}(H_2 \oplus H_1)$ then $\mathcal{E}(C(\mathcal{A}_1)) \in \text{Sym}^*(H_2^\vee \oplus H_1^\vee)$. Note that $\mathcal{E}(C(\mathcal{A}_1))$ can be canonically identified with a Hilbert-Schmidt operator

$$\text{Sym}^* H_1^\vee \rightarrow \text{Sym}^* H_2^\vee.$$

We use $\mathcal{E}(C(\mathcal{A}_1))$ to denote this Hilbert-Schmidt operator by abusing the notation. Similarly $\mathcal{A}_2 \in \mathbb{S}(H_3 \oplus H_2)$ induces a Hilbert-Schmidt operator

$$\mathcal{E}(C(\mathcal{A}_2)) : \text{Sym}^* H_2^\vee \rightarrow \text{Sym}^* H_3^\vee.$$

Now, we are ready for a key proposition which was used in Proposition 3.2.

Proposition 4.8. *Let \mathcal{A}_1 and \mathcal{A}_2 as above. Then $\mathcal{E}(C(\mathcal{A}_2)) \circ \mathcal{E}(C(\mathcal{A}_1)) = c(\mathcal{A}_2, \mathcal{A}_1)\mathcal{E}(C(\mathcal{A}_2 \circ \mathcal{A}_1))$ where*

$$c(\mathcal{A}_2, \mathcal{A}_1) = \frac{\det(\mathcal{A}_1)^{\frac{1}{4}} \cdot \det(\mathcal{A}_2)^{\frac{1}{4}}}{\det(\frac{A+M}{2})^{\frac{1}{2}} \det(\mathcal{A}_2 \circ \mathcal{A}_1)^{\frac{1}{4}}} \frac{\|\mathcal{E}(C(\mathcal{A}_1))\| \cdot \|\mathcal{E}(C(\mathcal{A}_2))\|}{\|\mathcal{E}(C(\mathcal{A}_2 \circ \mathcal{A}_1))\|}$$

and the norms are taken in the Fock spaces.

Proof. Using Proposition 4.10 we observe

$$\frac{\mathcal{E}(C(\mathcal{A}_i))}{\|\mathcal{E}(C(\mathcal{A}_i))\|} = \det(\mathcal{A}_i)^{\frac{1}{4}} Q \circ L_{\mathcal{A}_i} \circ Q^{-1}$$

and

$$L_{\mathcal{A}_2} \circ L_{\mathcal{A}_1} = C'(\mathcal{A}_2, \mathcal{A}_1)L_{\mathcal{A}_2 \circ \mathcal{A}_1}$$

where $C'(\mathcal{A}_2, \mathcal{A}_1) = \det(\frac{1}{2}(A + M))^{-\frac{1}{2}}$. Now,

$$\begin{aligned} \mathcal{E}(\mathcal{A}_2) \circ \mathcal{E}(\mathcal{A}_1) &= \det(\mathcal{A}_1)^{\frac{1}{4}} \cdot \det(\mathcal{A}_1)^{\frac{1}{4}} \cdot \|\mathcal{E}(C(\mathcal{A}_2))\| \cdot \|\mathcal{E}(C(\mathcal{A}_1))\| \\ &\quad Q \circ (L_{\mathcal{A}_2} \circ L_{\mathcal{A}_1}) \circ Q^{-1} \\ &= \frac{\det(\mathcal{A}_1)^{\frac{1}{4}} \cdot \det(\mathcal{A}_1)^{\frac{1}{4}}}{\det((\frac{1}{2}A + M))^{\frac{1}{2}}} \cdot \|\mathcal{E}(C(\mathcal{A}_2))\| \cdot \|\mathcal{E}(C(\mathcal{A}_1))\| \\ &\quad Q \circ (L_{\mathcal{A}_2 \circ \mathcal{A}_1}) \circ Q^{-1} \\ &= c(\mathcal{A}_2, \mathcal{A}_1)\mathcal{E}(C(\mathcal{A}_2 \circ \mathcal{A}_1)). \end{aligned}$$

□

4.2. Segal-Ito isomorphism finite dimensional case

Let μ be a standard Gaussian measure on \mathbb{R} . We recall it defines an inner product on \mathbb{R}^\vee . There is a unique isomorphism $Q : \text{Sym}^*(\mathbb{R}^\vee) \rightarrow L^2(\mathbb{R}, d\mu)$

which on $\text{Sym}^n(\mathbb{R}^\vee)$ is given by

$$Q(x^n) = h_n$$

where $x \in \mathbb{R}^\vee$ is dual to $1 \in \mathbb{R}$ and h_n is the degree n Hermite polynomial.

We recall the following which we will use later.

Fact 4.9. $e^{ax - \frac{1}{2}a^2} = \sum_{n=0}^\infty \frac{1}{n!} a^n h_n.$

Proof. We note

$$\begin{aligned} e^{ax - \frac{1}{2}a^2} &= e^{\frac{1}{2}x^2 - \frac{1}{2}(x-a)^2} \\ &= e^{\frac{1}{2}x^2} e^{-\frac{1}{2}(x-a)^2}. \end{aligned}$$

Now we write

$$e^{-\frac{1}{2}(x-a)^2} = \sum_{n=0}^\infty \frac{a^n}{n!} \frac{d^n}{da^n} e^{-\frac{1}{2}(x-a)^2} \Big|_{a=0}$$

which is the Maclaurin series expansion in the variable a . We note

$$\frac{d^n}{da^n} e^{-\frac{1}{2}(x-a)^2} \Big|_{a=0} = (-1)^n \frac{d^n}{dx^n} e^{-\frac{1}{2}x^2}.$$

This gives us

$$e^{ax - \frac{1}{2}a^2} = \sum_{n=0}^\infty \frac{a^n}{n!} (-1)^n e^{\frac{1}{2}x^2} \frac{d^n}{dx^n} e^{-\frac{1}{2}x^2} = \sum_{n=0}^\infty \frac{1}{n!} a^n h_n. \quad \square$$

Hence,

$$Q(\text{Exp}(ax)) = e^{ax - \frac{1}{2}a^2}$$

where $e^{ax - \frac{1}{2}a^2}$ is considered as a function on \mathbb{R} . It also follows

$$Q(\text{Exp } h) = e^{h - \frac{1}{2}\|h\|^2}$$

for all $h \in \mathbb{R}^\vee$ implying Q is unique.

We denote the complex valued L^2 functions by $L^2(\mathbb{R}, \mu, \mathbb{C})$. Then Q induces a unitary map $Q : \text{Sym}^*(\mathbb{C}^\vee) \rightarrow L^2(\mathbb{R}, \mu, \mathbb{C})$.

Let H be a finite dimensional real vector space and μ_H a centered Gaussian measure on H . The Gaussian measure induces an inner product on H .

Then the discussion above for the one dimensional case can be generalized in a routine way to construct the unique isomorphism

$$Q : \text{Sym}^*(H^\vee) \rightarrow L^2(H, \mu_H)$$

with the property

$$Q(\text{Exp } h) = e^{h - \frac{1}{2}\|h\|^2}.$$

As in the one dimensional case, Q induces a unitary map

$$Q : \text{Sym}^*(H_{\mathbb{C}}^\vee) \rightarrow L^2(H, \mu_H, \mathbb{C})$$

where $H_{\mathbb{C}}$ is the complexification of H and $L^2(H, \mu_H, \mathbb{C})$ is the L^2 space of complex valued functions. This map Q can be thought of as a “backward heat” operator which maps polynomials to L^2 functions. Next, we attempt to justify this assertion.

Let us identify H with \mathbb{R}^n such that μ_H corresponds to the standard Gaussian measure μ_n on \mathbb{R}^n . Let ν_n denote the measure on \mathbb{C}^n with the density $\pi^{-n}e^{-\|z\|^2}$ with respect to the standard Lebesgue measure on $\mathbb{C}^n = \mathbb{R}^{2n}$. We use $\mathcal{H}(\mathbb{C}^n)$ to denote the space of holomorphic functions on \mathbb{C}^n and $\mathcal{HL}^2(\mathbb{C}^n, \nu_n)$ to denote the space of L^2 holomorphic functions on \mathbb{C}^n .

There is a well known unitary map [10, Theorem 2.3] $S : L^2(\mathbb{R}^n, d\mu, \mathbb{C}) \rightarrow \mathcal{HL}^2(\mathbb{C}^n, \mu_{\mathbb{C}})$ defined by

$$Sf(z) = e^{-\frac{1}{2}\langle z, z \rangle} \int_{\mathbb{R}^n} f(u) e^{\langle z, u \rangle} d\mu_n(u) = \int_{\mathbb{R}^n} B_n(z, u) f(u) d\mu_n(u)$$

where

$$B_n(z, u) = e^{-\frac{1}{2}\langle z, z \rangle + \langle z, u \rangle}.$$

Here \langle, \rangle is the complex bilinear extension of the standard inner product on \mathbb{R}^n . The map S takes Hermite polynomials to the z^α s. Note that

$$(4.2) \quad Sf(z) = (2\pi)^{-\frac{n}{2}} \int_{\mathbb{R}^n} f(z - u) e^{-\frac{1}{2}\langle u, u \rangle} du.$$

Identifying $\mathcal{HL}^2(\mathbb{C}^n, \mu_{\mathbb{C}})$ with $\text{Sym}^*((\mathbb{C}^n)^\vee)$ using the tautological unitary map T called the *Taylor* map

$$T : \mathcal{HL}^2(\mathbb{C}^n, \mu_{\mathbb{C}}) \rightarrow \text{Sym}^*((\mathbb{C}^n)^\vee)$$

and identifying S with $T \circ S$, it can be easily checked that Q and S are inverse of each other. In fact, Q from $\mathcal{HL}^2(\mathbb{C}^n, \nu_n)$ to $L^2(\mathbb{R}^n, \mu_n, \mathbb{C})$ is given

by

$$QF(x) = \int_{\mathbb{C}^n} B(\bar{z}, x)F(z) d\nu_n(z).$$

We also observe that the following diagram is commutative.

$$\begin{CD} L^2(\mathbb{R}^n, d\mu) @>S>> \text{Sym}^*((\mathbb{R}^n)^\vee) \\ @VVV @VVV \\ L^2(\mathbb{R}^n, d\mu, \mathbb{C}) @>S>> \text{Sym}^*((\mathbb{C}^n)^\vee) \end{CD}$$

Let $k(y, x) \in L^2(\mathbb{R}^m \times \mathbb{R}^n, \mu_m \otimes \mu_n)$. Then k induces a Hilbert Schmidt operator

$$L_k : L^2(\mathbb{R}^n, \mu_n) \rightarrow L^2(\mathbb{R}^m, \mu_m)$$

given by

$$L_k f(y) = \int_{\mathbb{R}^n} k(y, x)f(x) d\mu_n(x).$$

Furthermore L_k induces a Hilbert-Schmidt operator

$$Q^{-1} \circ L_k \circ Q : \mathcal{H}L^2(\mathbb{C}^n, \nu_n) \rightarrow \mathcal{H}L^2(\mathbb{C}^m, \nu_m)$$

where $Q : \mathcal{H}L^2(\mathbb{C}^n, \nu_n) \rightarrow L^2(\mathbb{R}^n, \mu_n)$ and $Q^{-1} : L^2(\mathbb{R}^m, \mu_m) \rightarrow \mathcal{H}L^2(\mathbb{C}^m, \nu_m)$. We are interested in computing the kernel of the operator $Q^{-1} \circ L_k \circ Q$ explicitly. Let $K(w, z)$ denote the image of $k(y, x)$ under the isomorphism $Q^{-1} : L^2(\mathbb{R}^m \times \mathbb{R}^n, \mu_m \otimes \mu_n) \rightarrow \mathcal{H}L^2(\mathbb{C}^{m+n}, \nu_{m+n})$.

Proposition 4.10. $Q^{-1} \circ L_k \circ QF(w) = \int_{\mathbb{C}^n} K(w, \bar{z})F(z)\nu_n(z)$.

Proof. We note

$$\begin{aligned} (4.3) \quad & Q^{-1} \circ L_k \circ QF(w) \\ &= \int_{\mathbb{R}^m} \int_{\mathbb{R}^n} \int_{\mathbb{C}^n} B_n(\bar{z}, x)B_m(w, y)k(y, x)F(z) d\nu_n(z) d\mu_n(x) d\mu_m(y). \end{aligned}$$

Interchanging the order of integration (See for example Proposition 1.81 [5]) and using the fact

$$B_{m+n}(w, z), (y, x) = B_m(w, y)B_n(z, x),$$

we observe

$$\begin{aligned} & Q^{-1} \circ L_k \circ QF(w) \\ &= \int_{\mathbb{C}^n} \int_{\mathbb{R}^m} \int_{\mathbb{R}^n} B_{m+n}(w, \bar{z}), (y, x) k(y, x) d\mu_n(x) d\mu_m(y) F(z) d\nu_n(z) \end{aligned}$$

when F is a polynomial. Now we can conclude

$$Q^{-1} \circ L_k \circ QF(w) = \int_{\mathbb{C}^n} K(w, \bar{z}) F(z) d\nu_n(z).$$

Since all the maps in the consideration are continuous it is sufficient to check when F is a polynomial. \square

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