

Conserved quantities on asymptotically hyperbolic initial data sets

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In this article, we consider the limit of quasi-local conserved quantities [9, 31] at the infinity of an asymptotically hyperbolic initial data set in general relativity. These give notions of total energy-momentum, angular momentum, and center of mass. Our assumption on the asymptotics is less stringent than any previous ones to validate a Bondi-type mass loss formula. The Lorentz group acts on the asymptotic infinity through the exchange of foliations by coordinate spheres. For foliations aligning with the total energy-momentum vector, we prove that the limits of quasi-local center of mass and angular momentum are finite, and evaluate the limits in terms of the expansion coefficients of the metric and the second fundamental form.

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The notion of energy, linear momentum, angular momentum, center of mass and their conservation laws are of fundamental importance for any physical theory. However, there have been great difficulties in understanding these notions for gravitation since Einstein's time, as there is no well-defined concept of energy density due to the equivalence principle. It is nevertheless possible to use asymptotic symmetries to define these notions for an isolated system. At spatial infinity, the Arnowitt-Deser-Misner (ADM) energy-momentum [2] is well-understood. ADM energy-momentum is fundamental in general relativity and has been proven to be natural and to satisfy the important positivity property by the work of Schoen-Yau [26] and Witten [34]. It is also shown that the ADM energy-momentum satisfies important invariant properties in the work of Bartnik [4] and Chruściel [13]. There are several existing definitions of total angular momentum and total center of mass such as the Arnowitt-Deser-Misner (ADM) angular momentum [2, 3, 23] and the center of mass proposed by Huisken-Yau, Regge-Teitelboim, Beig-Ó Murchadha, Christodoulou and Schoen [5, 11, 17, 18, 23]. In [9], the authors proposed a new definition of quasi-local angular momentum and center of mass and used their limits to define new total conserved quantities for asymptotically flat initial data sets. The new definitions are finite for asymptotically flat initial data of order 1, and satisfy important dynamical formulas for solutions of the Einstein equation. See Theorem 7.4 and Theorem 9.6 of [9].

When the system is viewed from null infinity, the situation is more complicated. The notion of mass at null infinity is first studied by Bondi [6] and Trautman [29]. While the ADM energy-momentum at spatial infinity is conserved for solutions of the Einstein equation, the Bondi energy at null infinity is decreasing, see [6], [25] and [29]. A similar mass loss formula at null infinity is derived in [12] by Christodoulou and Klainerman, as a consequence of their proof of global stability of the Minkowski space. Their mass loss formula plays a key role in the study of nonlinear memory effect in gravitational radiation [10]. In [24], Rizzi proposed a new definition of total angular momentum at null infinity by assuming the existence of a special

foliation. He also studied the change of this total angular momentum by choosing a particular lapse function.

Schoen and Yau modified their proof for the positivity of ADM energy to prove the positivity of Bondi mass [27]. Their main strategy is to study a spacelike hypersurface asymptotic to the null cone at infinity. Both the induced metric and the second fundamental form are asymptotic to the metric of the hyperbolic 3-space. On such an asymptotically hyperbolic hypersurface, they solve Jang's equation to obtain an asymptotically flat manifold whose ADM energy is a positive multiple of the Bondi mass of the null cone.

Motivated by the proof of the positivity of Bondi mass and the study of asymptotically Anti-de-Sitter space, it is natural to find a suitable notion of mass and other conserved quantities for general asymptotically hyperbolic manifolds. In [33], X. Wang proved the positivity of mass for umbilical and conformally compact asymptotically hyperbolic manifolds satisfying the dominant energy condition. There are many works along this direction, see for example [1], [14], [15], [28] and [35]. In [36], Zhang proved a positivity and rigidity theorem for the mass of asymptotically hyperbolic manifolds without the umbilical assumption. In the same article, a new definition of total angular momentum was also proposed.

In this article, we study the total energy, linear momentum, angular momentum and center of mass at null infinity of asymptotically flat spacetime using the quasi-local energy of [31] and the quasi-local angular momentum and center of mass of [9]. The null infinity is modeled on 3-manifolds asymptotic to a hyperboloid in the Minkowski spacetime where the induced metric is isometric to the hyperbolic 3-space and the second fundamental form is the same as the induced metric. Let \mathbb{H}^3 denote the hyperbolic 3-space with metric $\frac{1}{r^2+1}dr^2 + r^2\tilde{\sigma}_{ab}du^adu^b$ where $\tilde{\sigma}_{ab}$ is the standard metric on the unit 2-sphere S^2 . Specifically, here are the decay conditions:

Definition 1.1. A triple (M, g, k) of a complete 3-manifold M , a Riemannian metric g on M , and a symmetric 2-tensor (the second fundamental form) k is said to be an asymptotically hyperbolic initial data set if there exists a compact subset K of M such that $M \setminus K$ is diffeomorphic to a finite union of ends $\cup_i \mathbb{H}^3 \setminus B_i$ where each B_i is a geodesic ball in \mathbb{H}^3 . Under the diffeomorphism, we have

$$g = g_{rr}dr^2 + 2g_{ra}drdu^a + g_{ab}du^adu^b \quad \text{and} \quad k = g + p,$$

where

$$g_{rr} = \frac{1}{r^2} - \frac{1}{r^4} + \frac{g_{rr}^{(-5)}}{r^5} + \frac{g_{rr}^{(-6)}}{r^6} + O(r^{-7}), \quad g_{ra} = \frac{g_{ra}^{(-3)}}{r^3} + O(r^{-4}),$$

$$g_{ab} = r^2 \tilde{\sigma}_{ab} + g_{ab}^{(0)} + \frac{g_{ab}^{(-1)}}{r} + \frac{g_{ab}^{(-2)}}{r^2} + O(r^{-3}),$$

and

$$p_{rr} = \frac{p_{rr}^{(-4)}}{r^4} + O(r^{-5}) \quad p_{ra} = \frac{p_{ra}^{(-3)}}{r^3} + O(r^{-4}),$$

$$p_{ab} = p_{ab}^{(0)} + \frac{p_{ab}^{(-1)}}{r} + \frac{p_{ab}^{(-2)}}{r^2} + O(r^{-3}).$$

Here $\tilde{\sigma}_{ab}$ is the standard round metric on the unit 2-sphere S^2 . $g_{rr}^{(-5)}$, $g_{rr}^{(-6)}$, $p_{rr}^{(-4)}$ are considered to be functions on S^2 that do not depend on r , $g_{ra}^{(-3)}$, $p_{ra}^{(-3)}$ are considered to be one-forms on S^2 that do not depend on r , and $g_{ab}^{(0)}$, $g_{ab}^{(-1)}$, $g_{ab}^{(-2)}$, $p_{ab}^{(0)}$, $p_{ab}^{(-1)}$, $p_{ab}^{(-2)}$ are considered to be symmetric two-tensors on S^2 that do not depend on r . Furthermore, we assume that $p_{ab}^{(0)}$ and $g_{ab}^{(0)}$ are traceless with respect to $\tilde{\sigma}_{ab}$.

We assume that the triple (M, g, k) satisfies the following dominant energy condition:

Definition 1.2. (M, g, k) satisfies the dominant energy condition if

$$\mu = \frac{1}{2}(R(g) + (tr_g k)^2 - |k|_g^2)$$

$$\mathfrak{J}_i = D^j(k_{ij} - (tr_g k)g_{ij})$$

satisfies

$$\mu \geq |\mathfrak{J}|.$$

Here $R(g)$ is the scalar curvature of the metric g and D^j is the covariant derivative with respect to the metric g .

Definition 1.3. The mass aspect function m of an asymptotically hyperbolic initial data set is defined to be

$$(1.1) \quad m = \frac{3}{2} tr_{S^2} g_{ab}^{(-1)} + tr_{S^2} p_{ab}^{(-1)} + g_{rr}^{(-5)}.$$

The energy-momentum of an asymptotically hyperbolic initial data set is defined as follows:

Definition 1.4. Let (M, g, k) be an asymptotically hyperbolic initial data set. The energy-momentum of (M, g, k) is the four vector $(E_{AH}, P_{AH}^i), i = 1, 2, 3$ where

$$E_{AH}(M, g, k) = \frac{1}{8\pi} \int_{S^2} m dS^2$$

$$P_{AH}^i(M, g, k) = \frac{1}{8\pi} \int_{S^2} \tilde{X}^i m dS^2, i = 1, 2, 3$$

where $\tilde{X}^i, i = 1, 2, 3$ are the three standard coordinate functions on S^2 .

Remark 1.5. There are two types of asymptotically hyperbolic initial data sets. One is modeled on the hyperbola in the Minkowski spacetime, and thus $g = k$ (umbilical) is the same as the metric on \mathbb{H}^3 . The other is modeled on a static slice of the Anti-de-Sitter spacetime. Most studies only consider a Riemannian manifold that is asymptotically to a hyperbolic space and it is implicitly assumed that the second fundamental form is either zero or the same as the induced metric. However, in studying the mass of an initial data set, it is important to take into account of the second fundamental form. For example, there exists an asymptotically umbilical spacelike hypersurface in the Schwarzschild spacetime with exactly hyperbolic induced metric and an asymptotically totally geodesic spacelike hypersurface in the Anti-de-Sitter Schwarzschild spacetime with exactly hyperbolic induced metric. In both cases, the mass can only be read off from the second fundamental form.

Remark 1.6. In [33], X. Wang defined energy-momentum for asymptotically hyperbolic Riemannian manifolds with $g_{ra} = 0$ and $g_{rr} = \frac{1}{r^2+1}$. It is easy to see that his definition agrees with ours in this simplest case.

In [7] and [32], the authors proved that the limit of the Wang-Yau quasi-local energy of the coordinates spheres recovers the ADM energy momentum vector at spatial infinity and the Bondi-Sachs energy momentum vector at null infinity. In section 3, we prove

Theorem A (Theorem 3.2) *Given an asymptotically hyperbolic initial data set (M, g, k) , the limit of the quasi-local energy of Σ_r with reference embedding X_r into \mathbb{R}^3 is a linear function dual to the energy-momentum of (M, g, k) .*

In [19], Kwong and Tam evaluated the limit of a quasi-local energy-momentum defined by Wang and Yau in [30] at the infinity of an asymptotically hyperbolic manifold. The quasi-local energy-momentum of [30] again

only uses the hyperbolic space as reference and does not include the effect of the second fundamental form of the initial data.

In Section 4, we prove the following theorem regarding the causality of the energy-momentum vector of an asymptotically hyperbolic initial data set.

Theorem B (Theorem 4.2) *Suppose (M, g, k) is an asymptotically hyperbolic initial data set satisfying the dominant energy condition, then its energy-momentum 4-vector is future directed non-spacelike. Namely,*

$$E_{AH} \geq \sqrt{\sum_i (P_{AH}^i)^2}.$$

To show that the energy-momentum 4-vector is non-spacelike, we apply a Lorentz boost of the Minkowski spacetime to the coordinate spheres, obtain a new foliation, and compute the limit of quasi-local energy with respect to the new foliation. For any future directed unit timelike vector (a^0, a^i) in $\mathbb{R}^{3,1}$, we show that the limit of the Liu-Yau quasi-local energy of corresponding boost is $a^0 E_{AH} - \sum_i a^i P_{AH}^i$. The positivity of the Liu-Yau quasi-local energy [20] implies the energy momentum 4-vector is non-spacelike.

In addition to the positive mass theorem, we derive an energy loss formula for asymptotically hyperbolic initial data sets. Our energy loss formula is similar to the Bondi mass loss formula derived in [12]. In particular, the source of the energy loss is $p_{ab}^{(0)} + g_{ab}^{(0)}$, which plays the role of the leading order term of $\hat{\chi}$ in the mass loss formula derived in [12]. Theorem 3.2 enables us to use the limit of the Liu-Yau mass for the total energy of asymptotically hyperbolic initial data set and to compute the variation of the Liu-Yau mass along the Einstein equation. This approach is similar to that of [12], where the total energy of a null cone is defined to be the limit of the Hawking mass and the mass loss formula is derived by studying the variation of the Hawking mass. The energy loss formula we derive is the following:

Theorem C (Theorem 5.1) *Let (M, g, k) be an asymptotically hyperbolic initial data set satisfying the vacuum constraint equation. Let $(M, g(t), k(t))$ be the solution to the vacuum Einstein equation with $g(0) = g$ and $k(0) = k$, and with lapse $\sqrt{r^2 + 1}$ and shift vector $-re_3$ where e_3 is the unit normal vector of the coordinate spheres on $(M, g(t), k(t))$. Let $\Sigma_{r,t}$ be the coordinate spheres on $(M, g(t), k(t))$ and*

$$M(t) = \lim_{r \rightarrow \infty} m_{LY}(\Sigma_{r,t})$$

be the limit. Along the vacuum Einstein equation, we have

$$\partial_t M(t) = -\frac{1}{8\pi} \int_{S^2} |p_{ab}^{(0)} + g_{ab}^{(0)}|^2 dS^2.$$

Remark 1.7. The lapse function and shift vector considered here are natural for asymptotically hyperbolic initial data sets. The Minkowski metric can be written as

$$g = -dt^2 - \frac{2r}{\sqrt{r^2 + 1}} dr dt + \frac{dr^2}{r^2 + 1} + r^2(d\theta^2 + \sin^2 \theta d\phi^2).$$

Each level set of t is a standard hyperboloid. Let Σ_r be a level set of r on a hyperboloid, e_3 be its unit normal vector in the hyperboloid, and e_4 be the unit normal of the hyperboloid. Thus

$$\begin{aligned} e_3 &= \sqrt{r^2 + 1} \frac{\partial}{\partial r} \\ e_4 &= \frac{1}{\sqrt{r^2 + 1}} \frac{\partial}{\partial t} + r \frac{\partial}{\partial r} \end{aligned}$$

and the Killing vector field $\frac{\partial}{\partial t}$ can be written as

$$\frac{\partial}{\partial t} = \sqrt{r^2 + 1} e_4 - r e_3.$$

Hence, the lapse function is $\sqrt{r^2 + 1}$ and the shift vector is $-r e_3$.

For spacetimes with matter fields, we prove a similar mass loss formula with an additional term from the matter field assuming some natural decay conditions on the null component of the stress-energy density of the matter field (see Definition 5.2 and Theorem 5.3). In particular, for solutions of the Einstein-Maxwell equation, our result resembles the mass loss formula derived in [37] by Zipser.

While the ADM energy and linear momentum are well-understood for asymptotically flat initial data sets, defining total angular momentum and center of mass is much more complicated and there were several delicate issues concerning the finiteness, well-definedness, and physical validity. In [9], the authors defined new total conserved quantities for asymptotically flat initial data sets using the limit of quasi-local conserved quantities and proved the finiteness of the new total angular momentum and center of mass

for asymptotically flat initial data sets of order 1. In this article, we compute the limit of quasi-local angular momentum and center of mass at infinity of asymptotically hyperbolic initial data sets. Unlike the asymptotically flat case, we show that, for an asymptotically hyperbolic initial data set, the limit is finite only for the foliation with vanishing linear momentum. Assuming the energy-momentum vector is timelike for a given foliation, we show that there exists another foliation with vanishing linear momentum. For the new foliation, the total center of mass C^i and total angular momentum J^i (See Definition 6.1) can be explicitly computed in terms of the expansion coefficients of the metric and the second fundamental form.

Theorem D (Theorem 7.3) *Given an asymptotically hyperbolic initial data set (M, g, k) , for the foliation with vanishing linear momentum, the total center of mass C^i and total angular momentum J^i of (M, g, k) are*

$$C^i = \frac{1}{8\pi} \int_{S^2} \tilde{X}^i (2tr_{S^2} g_{ab}^{(-2)} + g_{rr}^{(-6)} + \tilde{\nabla}^a g_{ra}^{(-3)} + tr_{S^2} p_{ab}^{(-2)}) dS^2$$

$$J^i = \frac{1}{8\pi} \int_{S^2} \tilde{X}^i \tilde{\epsilon}^{ab} \tilde{\nabla}_b (g_{ra}^{(-3)} + p_{ra}^{(-3)}) dS^2.$$

2. Quasi-local energy-momentum, angular momentum and center of mass

We recall the definition of quasi-local energy-momentum defined in [31] and quasi-local angular momentum and center of mass defined in [9]. Let Σ be a closed embedded spacelike 2-surface in a spacetime with spacelike mean curvature vector H . The data used in the definition of quasi-local mass is the triple $(\sigma, |H|, \alpha_H)$ where σ is the induced metric, $|H|$ is the norm of the mean curvature vector and α_H is the connection one form of the normal bundle with respect to the mean curvature vector

$$\alpha_H(\cdot) = \left\langle \nabla_{(\cdot)}^N \frac{J}{|H|}, \frac{H}{|H|} \right\rangle$$

where J is the reflection of H through the incoming light cone in the normal bundle.

Given an isometric embedding $X : \Sigma \rightarrow \mathbb{R}^{3,1}$ and a constant future timelike unit vector $T_0 \in \mathbb{R}^{3,1}$, we consider the projected embedding \hat{X} into the orthogonal complement of T_0 . We denote the mean curvature of the image by \hat{H} .

In terms of $\tau = -\langle X, T_0 \rangle$, the quasi-local energy with respect to the pair (X, T_0) is

$$E(\Sigma, \tau) = \frac{1}{8\pi} \int_{\widehat{\Sigma}} \widehat{H} d\widehat{\Sigma} - \frac{1}{8\pi} \int_{\Sigma} \left[\sqrt{1 + |\nabla\tau|^2} \cosh\theta |H| - \nabla\tau \cdot \nabla\theta - \alpha_H(\nabla\tau) \right] d\Sigma,$$

where $\theta = \sinh^{-1} \left(\frac{-\Delta\tau}{|H|\sqrt{1+|\nabla\tau|^2}} \right)$ and ∇ and Δ are the gradient and Laplace operator, respectively, with respect to σ . In particular, $E(\Sigma, 0)$ is the same as the Liu-Yau quasi-local mass, $m_{LY}(\Sigma)$.

Assuming the spacetime satisfying the dominant energy condition, the quasi-local energy defined above is non-negative for any admissible pair of (X, T_0) . The quasi-local mass of the surface Σ is defined to be the infimum of the quasi-local energy with respect to all admissible pairs.

Given an isometric embedding X of Σ into $\mathbb{R}^{3,1}$, let H_0 and α_{H_0} be the mean curvature vector and connection one form of $X(\Sigma)$ in $\mathbb{R}^{3,1}$. We introduce ρ and a j_a as follows:

$$(2.1) \quad \rho = \frac{\sqrt{|H_0|^2 + \frac{(\Delta\tau)^2}{1+|\nabla\tau|^2}} - \sqrt{|H|^2 + \frac{(\Delta\tau)^2}{1+|\nabla\tau|^2}}}{\sqrt{1 + |\nabla\tau|^2}}$$

$$(2.2) \quad j_a = \rho \nabla_a \tau - \nabla_a \left[\sinh^{-1} \left(\frac{\rho \Delta\tau}{|H_0||H|} \right) \right] - (\alpha_{H_0})_a + (\alpha_H)_a.$$

In terms of these, the quasi-local energy is $\frac{1}{8\pi} \int_{\Sigma} (\rho + j_a \nabla^a \tau)$.

A critical point of the quasi-local energy satisfies the optimal isometric embedding equation. A pair of an embedding $X : \Sigma \hookrightarrow \mathbb{R}^{3,1}$ and an observer T_0 satisfies the optimal isometric embedding equation for $(\sigma_{ab}, |H|, (\alpha_H)_a)$ if X is an isometric embedding and

$$(2.3) \quad \text{div}_{\sigma} j = 0.$$

We define the quasi-local angular momentum and center of mass for each optimal isometric embedding. Let (x^0, x^1, x^2, x^3) be the standard coordinate of $\mathbb{R}^{3,1}$. We recall that K is a rotational Killing field if K is the image of $x^i \frac{\partial}{\partial x^j} - x^j \frac{\partial}{\partial x^i}$ under a Lorentz transformation. Similarly, K is a boost Killing field if K is the image of $x^0 \frac{\partial}{\partial x^i} + x^i \frac{\partial}{\partial x^0}$ under a Lorentz transformation.

Definition 2.1. The quasi-local conserved quantity of Σ with respect to an optimal isometric embedding (X, T_0) and a Killing field K is

$$(2.4) \quad E(\Sigma, X, T_0, K) = \frac{(-1)}{8\pi} \int_{\Sigma} \left[\langle K, T_0 \rangle \rho + j(K^\top) \right] d\Sigma,$$

where K^\top is the tangential part of K to $X(\Sigma)$.

Suppose $T_0 = A(\frac{\partial}{\partial x^0})$ for a Lorentz transformation A , then the quasi-local conserved quantities corresponding to $A(x^i \frac{\partial}{\partial x^j} - x^j \frac{\partial}{\partial x^i})$, $i < j$ are called the quasi-local angular momentum integral with respect to T_0 and the quasi-local conserved quantities corresponding to $A(x^i \frac{\partial}{\partial x^0} + x^0 \frac{\partial}{\partial x^i})$, $i = 1, 2, 3$ are called the quasi-local center of mass integral with respect to T_0 .

3. Limit of quasi-local energy of coordinate spheres.

We first prove the following lemma about the geometry of coordinate spheres in an asymptotically hyperbolic initial data set. Given an asymptotically hyperbolic initial data set (M, g, k) in a spacetime N , such that g and k are the induced metric and second fundamental form of M , respectively.

Lemma 3.1. *Let Σ_r be coordinate spheres in an asymptotically hyperbolic space. Let σ, H be the mean curvature vector of Σ_r and α_H be the induced metric, mean curvature vector and connection one-form of the normal bundle in mean curvature gauge of Σ_r in the spacetime N . We have*

$$(3.1) \quad \begin{aligned} \sigma_{ab} &= r^2 \tilde{\sigma}_{ab} + g_{ab}^{(0)} + \frac{g_{ab}^{(-1)}}{r} + O(r^{-2}), \\ |H| &= \frac{2}{r} - \frac{m}{r^2} + O(r^{-3}), \\ (\alpha_H)_a &= \frac{\partial_a(m)}{2r} + O(r^{-2}), \end{aligned}$$

where m is the mass aspect function defined in Definition 1.3.

Proof. The expansion for σ_{ab} follows directly from the definition of asymptotically hyperbolic initial data set. The inverse metric is

$$\sigma^{ab} = \frac{\tilde{\sigma}^{ab}}{r^2} - \frac{(g^{(0)})^{ab}}{r^4} - \frac{(g^{(-1)})^{ab}}{r^5} + O(r^{-6}).$$

Let e_3 be the unit normal vector of Σ_r in M and e_4 be the unit normal vector of M in N . We compute $\langle H, e_3 \rangle$ and $\langle H, e_4 \rangle$. e_3 is given by

$$e_3 = \frac{1}{\sqrt{g_{rr} - |V|_\sigma^2}} \left(\frac{\partial}{\partial r} + V^a \frac{\partial}{\partial u^a} \right)$$

where $V_a = -g_{ra}$, and the second fundamental form of Σ_r in the direction of e_3 is

$$\begin{aligned} (3.2) \quad \left\langle \nabla_{\frac{\partial}{\partial u^a}} e_3, \frac{\partial}{\partial u^b} \right\rangle &= \sqrt{r^2 + 1 - \frac{g_{rr}^{(-5)}}{r}} \left\langle \nabla_{\frac{\partial}{\partial u^a}} \frac{\partial}{\partial r}, \frac{\partial}{\partial u^b} \right\rangle + O(r^{-2}) \\ &= \sqrt{r^2 + 1 - \frac{g_{rr}^{(-5)}}{r}} \left(r \tilde{\sigma}_{ab} - \frac{g_{ab}^{(-1)}}{2r^2} \right) + O(r^{-2}). \end{aligned}$$

Hence,

$$\begin{aligned} -\langle H, e_3 \rangle &= \sigma^{ab} \sqrt{r^2 + 1 - \frac{g_{rr}^{(-5)}}{r}} \left(r \tilde{\sigma}_{ab} - \frac{g_{ab}^{(-1)}}{2r^2} \right) + O(r^{-4}) \\ &= r \left(1 + \frac{1}{2r^2} - \frac{g_{rr}^{(-5)}}{2r^3} \right) \left(r \tilde{\sigma}_{ab} - \frac{g_{ab}^{(-1)}}{2r^2} \right) \\ &\quad \times \left(\frac{\tilde{\sigma}^{ab}}{r^2} - \frac{(g^{(0)})^{ab}}{r^4} - \frac{(g^{(-1)})^{ab}}{r^5} \right) + O(r^{-4}) \\ &= 2 + \frac{1}{r^2} - \frac{g_{rr}^{(-5)} + \frac{3}{2} tr_{S^2} g_{ab}^{(-1)}}{r^3} + O(r^{-4}). \end{aligned}$$

For $\langle H, e_4 \rangle$, we have

$$\begin{aligned} -\langle H, e_4 \rangle &= \sigma^{ab} (\sigma_{ab} + p_{ab}) \\ &= 2 + \frac{tr_{S^2} p_{ab}^{(-1)}}{r^3} + O(r^{-4}). \end{aligned}$$

As a result,

$$\begin{aligned} |H|^2 &= \langle H, e_3 \rangle^2 - \langle H, e_4 \rangle^2 \\ &= \frac{4}{r^2} - \frac{4g_{rr}^{(-5)} + 4tr_{S^2} p_{ab}^{(-1)} + 6tr_{S^2} g_{ab}^{(-1)}}{r^3} + O(r^{-4}) \end{aligned}$$

and the expansion for $|H|$ follows.

Recall from [32], we have

$$(\alpha_H)_a = -k(e_3, \partial_a) + \nabla_a \theta$$

where

$$\sinh(\theta) = \frac{-\langle H, e_4 \rangle}{|H|}.$$

By the definition of asymptotically hyperbolic initial data set, we have $k(e_3, \partial_a) = O(r^{-2})$. Using the expansion for $\langle H, e_4 \rangle$ and $|H|$, we conclude that

$$e^\theta = 2r + (tr_{S^2} p_{ab}^{(-1)} + \frac{3}{2} tr_{S^2} g_{ab}^{(-1)} + g_{rr}^{(-5)}) + O(r^{-1})$$

which implies

$$\nabla_a \theta = \frac{\partial_a (tr_{S^2} p_{ab}^{(-1)} + \frac{3}{2} tr_{S^2} g_{ab}^{(-1)} + g_{rr}^{(-5)})}{2r} + O(r^{-2}). \quad \square$$

In [32], it is shown that for an asymptotically flat initial data set, the limit of the quasi-local energy of coordinate spheres with reference isometric embedding into \mathbb{R}^3 recovers the ADM energy-momentum 4-vector. We prove a similar result for asymptotically hyperbolic initial data sets. Let X_r be the isometric embedding of Σ_r into the totally geodesics \mathbb{R}^3 in $\mathbb{R}^{3,1}$.

Theorem 3.2. *Given an asymptotically hyperbolic initial data set (M, g, k) , the limit of the quasi-local energy of Σ_r with reference embedding X_r into \mathbb{R}^3 is a linear function dual to the energy-momentum of (M, g, k) .*

Proof. Let H_0 be the mean curvature of the image of X_r in \mathbb{R}^3 . From the expansion of the induced metric and the result in [7], it follows that $H_0 = \frac{2}{r} + O(r^{-3})$ and thus

$$\lim_{r \rightarrow \infty} \frac{|H|}{H_0} = 1.$$

By Corollary 2.1 of [32], the limit of

$$\lim_{r \rightarrow \infty} (\Sigma_r, X_r, T_0)$$

is a linear function dual to the four vector (e, p^i) where

$$e = \frac{1}{8\pi} \lim_{r \rightarrow \infty} \int_{\Sigma_r} (H_0 - |H|) d\Sigma_r$$

$$p^i = \frac{1}{8\pi} \lim_{r \rightarrow \infty} \int_{\Sigma_r} X^i \nabla^a (\alpha_H)_a d\Sigma_r.$$

By Equation (3.1),

$$\begin{aligned} & \frac{1}{8\pi} \int_{\Sigma_r} (H_0 - |H|) d\Sigma_r \\ &= \frac{1}{8\pi} \int_{S^2} \left(\frac{3}{2} tr_{S^2} \tilde{g}_{ab}^{(-1)} + tr_{S^2} p_{ab}^{(-1)} + g_{rr}^{(-5)} \right) dS^2 + O(r^{-1}) \end{aligned}$$

and

$$\begin{aligned} & \frac{1}{8\pi} \int_{\Sigma_r} X^i \nabla^a (\alpha_H)_a d\Sigma_r \\ &= -\frac{1}{8\pi} \int_{S^2} \tilde{X}^i \left(\frac{3}{2} tr_{S^2} g_{ab}^{(-1)} + tr_{S^2} p_{ab}^{(-1)} + g_{rr}^{(-5)} \right) dS^2 + O(r^{-1}). \end{aligned}$$

This finishes the proof of the theorem. □

Remark 3.3. The data $(\sigma, |H|, \alpha_H)$ has the same form of expansion as that of [7] and section 3 and 4 of [7] can be applied to the asymptotically hyperbolic case as well. Assuming the total energy momentum vector (E_{AH}, P_{AH}^i) is timelike, Let

$$M_{AH} = \sqrt{E_{AH}^2 - \sum_i (P_{AH}^i)^2}.$$

There is a local minimum $(X(r), T_0(r))$ of the quasi-local energy with

$$\begin{aligned} X^0(r) &= (X^0)^{(0)} + O(r^{-1}) \\ X^i(r) &= r \tilde{X}^i + \frac{(X^i)^{(-1)}}{r} + O(r^{-2}) \\ T_0(r) &= (a^0, a^1, a^2, a^3) + \frac{T_0^{(-1)}}{r} + O(r^{-2}). \end{aligned}$$

where (a^0, a^1, a^2, a^3) is the unit timelike vector aligned with the total energy-momentum 4-vector.

$$M_{AH}(a^0, a^1, a^2, a^3) = (E_{AH}, P_{AH}^i).$$

Moreover,

$$\lim_{r \rightarrow \infty} E(\Sigma_r, X(r), T_0(r)) = M_{AH}.$$

4. Spacetime positive mass theorem

Theorem 3.2 and the positivity of Liu-Yau quasi-local mass implies $E_{AH} \geq 0$ if (M, g, k) satisfies the dominant energy condition. We prove that the energy-momentum vector is non-spacelike by evaluating the limit of quasi-local energy on boosted coordinate spheres in M .

The hyperbolic space can be isometrically embedded into $\mathbb{R}^{3,1}$ with metric $-dt^2 + \sum_{i=1}^3(dx^i)^2$ as the set

$$\mathbb{H}^3 = \left\{ (t, x^i) \mid t^2 = 1 + \sum_i (x^i)^2, t > 0 \right\}.$$

The Lorentz group of $\mathbb{R}^{3,1}$ acts on \mathbb{H}^3 by isometry. In particular, we consider the isometries corresponding to boost elements of the Lorentz group. Let (M, g, k) be an asymptotically hyperbolic initial data set and (r, u^1, u^2) be the asymptotically hyperbolic coordinate system at infinity. For each future directed unit timelike vector (a^0, a^i) in $\mathbb{R}^{3,1}$, we consider the family of surfaces $\Sigma_R, R \gg 1$ on M defined by

$$\Sigma_R = \{(r, u^1, u^2) : r = R\tilde{F}(u^1, u^2)\}$$

where

$$\tilde{F} = \frac{1}{a^0 + \sum_i a^i \tilde{X}^i}$$

and $\tilde{X}^i(u^1, u^2), i = 1, 2, 3$ are the three standard coordinate functions on S^2 . They are all -2 eigenfunctions with respect to the standard round metric $\tilde{\sigma}_{ab}$ on S^2 . Notice that $\tilde{F}^2 \tilde{\sigma}_{ab}$ is isometric to $\tilde{\sigma}_{ab}$ and \tilde{F} satisfies the following constant Gauss curvature equation

$$\frac{1}{\tilde{F}^2} (1 - \tilde{\Delta} \ln \tilde{F}) = 1$$

where $\tilde{\Delta}$ is the Laplace operator with respect to the metric $\tilde{\sigma}_{ab}$.

Lemma 4.1. *Let Σ_R be the above family of surfaces. Let σ and H be the induced metric and the mean curvature vector of Σ_R , respectively. Then,*

$$(4.1) \quad \begin{aligned} \sigma_{ab} &= (\tilde{F}^2 R^2) \tilde{\sigma}_{ab} + g_{ab}^{(0)} + \frac{\tilde{F}_a \tilde{F}_b}{\tilde{F}^2} + \frac{g_{ab}^{(-1)}}{\tilde{F} R} + O(R^{-2}) \\ |H| &= \frac{2}{R} - \frac{g_{rr}^{(-5)} + tr_{S^2} p_{ab}^{(-1)} + \frac{3}{2} tr_{S^2} g_{ab}^{(-1)}}{\tilde{F}^3 R^2} + O(R^{-3}). \end{aligned}$$

Proof. The tangent space of Σ_R is spanned by $\frac{\partial}{\partial u^a} + r_a \frac{\partial}{\partial r}$, $a = 1, 2$. It follows immediately from the definition of asymptotically hyperbolic initial data sets that

$$\sigma_{ab} = (\tilde{F}^2 R^2) \tilde{\sigma}_{ab} + g_{ab}^{(0)} + \frac{\tilde{F}_a \tilde{F}_b}{\tilde{F}^2} + \frac{g_{ab}^{(-1)}}{\tilde{F} R} + O(R^{-2}).$$

As a result, we also have

$$\sigma^{ab} = \tilde{F}^{-2} R^{-2} \left(\tilde{\sigma}^{ab} - \frac{(g^{(0)})^{ab}}{\tilde{F}^2 R^2} - \frac{\tilde{F}^a \tilde{F}^b}{\tilde{F}^4 R^2} - \frac{(g^{(-1)})^{ab}}{\tilde{F}^3 R^3} \right) + O(R^{-6}).$$

Let e_3 be the unit normal of Σ_R in M and e_4 be the unit normal of M in N . We compute

$$\begin{aligned} -\langle H, e_4 \rangle &= \sigma^{ab} (\sigma_{ab} + p_{ab} + r_a r_b p_{rr} + r_a p_{rb} + r_b p_{ra}) \\ &= 2 + \frac{tr_{S^2} p_{ab}^{(-1)}}{\tilde{F}^3 R^3} + O(R^{-4}). \end{aligned}$$

To compute $\langle H, e_3 \rangle$, we start with

$$e_3 = \frac{1}{\sqrt{g_{rr} + V^a V^b g_{ab} - 2V^a g_{ar}}} \left(\frac{\partial}{\partial r} + V^a \frac{\partial}{\partial u^a} \right)$$

where

$$V_a = \frac{-r_a}{1 + r^2} + O(R^{-3}).$$

By definition,

$$\begin{aligned} -\langle H, e_3 \rangle &= \frac{\sigma^{ab}}{\sqrt{g_{rr} + V^d V^e g_{de} - 2V^d g_{dr}}} \\ &\times \left\langle \nabla_{\frac{\partial}{\partial u^a} + r_a \frac{\partial}{\partial r}} \left(\frac{\partial}{\partial r} - \frac{r^c}{1 + r^2} \frac{\partial}{\partial u^c} \right), \frac{\partial}{\partial u^b} + r_b \frac{\partial}{\partial r} \right\rangle + O(R^{-4}). \end{aligned}$$

Plugging in $r = R\tilde{F}(u^1, u^2)$, we deduce

$$\frac{1}{\sqrt{g_{rr} + V^a V^b g_{ab} - 2V^a g_{ar}}} = R\tilde{F} \left[1 + \frac{\tilde{F}^2 - |\tilde{\nabla} \tilde{F}|^2}{2R^2 \tilde{F}^4} - \frac{g_{rr}^{(-5)}}{2R^3 \tilde{F}^3} \right] + O(R^{-3}).$$

Continuing the calculation,

$$\begin{aligned}
 & \left\langle \nabla_{\frac{\partial}{\partial u^a} + r_a \frac{\partial}{\partial r}} \left(\frac{\partial}{\partial r} - \frac{r^c}{1+r^2} \frac{\partial}{\partial u^c} \right), \frac{\partial}{\partial u^b} + r_b \frac{\partial}{\partial r} \right\rangle \\
 &= \frac{1}{2} \partial_r g_{ab} - \left\langle \nabla_{\frac{\partial}{\partial u^a}} \frac{r^c}{1+r^2} \frac{\partial}{\partial u^c}, \frac{\partial}{\partial u^b} \right\rangle + r_a r_b \frac{1}{2} \partial_r \left(\frac{1}{1+r^2} \right) \\
 &\quad - \left\langle \nabla_{r_a \frac{\partial}{\partial r}} \frac{r^c}{1+r^2} \frac{\partial}{\partial u^c}, \frac{\partial}{\partial u^b} \right\rangle - \left\langle \nabla_{\frac{\partial}{\partial u^a}} \frac{r^c}{1+r^2} \frac{\partial}{\partial u^c}, + r_b \frac{\partial}{\partial r} \right\rangle + O(R^{-3}) \\
 &= r \sigma_{ab} - \frac{g_{ab}^{(-1)}}{2r^2} + \frac{r_a r_b}{r^3} + \nabla_a \nabla_b \frac{1}{r} + O(R^{-3}) \\
 &= R \tilde{F} \tilde{\sigma}_{ab} + \frac{1}{R} \left(\frac{\tilde{F}_a \tilde{F}_b}{\tilde{F}^3} + \nabla_a \nabla_b \frac{1}{\tilde{F}} \right) - \frac{g_{ab}^{(-1)}}{2R^2 \tilde{F}^2} + O(R^{-3}).
 \end{aligned}$$

As a result,

$$\begin{aligned}
 -\langle H, e_3 \rangle &= 2 + \frac{1}{R^2} \left(\frac{1}{\tilde{F}^2} - \frac{|\tilde{\nabla} \tilde{F}|^2}{\tilde{F}^4} + \frac{1}{\tilde{F}} \tilde{\Delta} \frac{1}{\tilde{F}} \right) \\
 &\quad - \frac{1}{R^3} \left(\frac{g_{rr}^{(-5)}}{\tilde{F}^3} + \frac{3tr S^2 g_{ab}^{(-1)}}{2\tilde{F}^3} \right) + O(R^{-4}) \\
 &= 2 + \frac{1}{R^2} - \frac{1}{R^3} \left(\frac{g_{rr}^{(-5)}}{\tilde{F}^3} + \frac{3tr S^2 g_{ab}^{(-1)}}{2\tilde{F}^3} \right) + O(R^{-4}).
 \end{aligned}$$

In the last equality, we used that

$$\frac{1}{\tilde{F}^2} - \frac{|\tilde{\nabla} \tilde{F}|^2}{\tilde{F}^4} + \frac{1}{\tilde{F}} \tilde{\Delta} \frac{1}{\tilde{F}} = \frac{1}{\tilde{F}^2} (1 - \tilde{\Delta} \ln \tilde{F}) = 1.$$

Therefore,

$$\begin{aligned}
 |H|^2 &= \langle H, e_3 \rangle^2 - \langle H, e_4 \rangle^2 \\
 &= \frac{4}{R^2} - \frac{4g_{rr}^{(-5)} + 4tr S^2 p_{ab}^{(-1)} + 6tr S^2 g_{ab}^{(-1)}}{R^3 \tilde{F}^3} + O(R^{-4}).
 \end{aligned}$$

□

We prove the positive mass theorem for asymptotically hyperbolic initial data set in the following.

Theorem 4.2. *Suppose (M, g, k) is an asymptotically hyperbolic initial data set satisfying the dominant energy condition, then its energy-momentum*

4-vector is future directed non-spacelike. Namely,

$$E_{AH} \geq \sqrt{\sum_i (P_{AH}^i)^2}.$$

Proof. For any future directed unit timelike vector (a^0, a^i) in $\mathbb{R}^{3,1}$, let $\tilde{F} = \frac{1}{a^0 + \sum_i a^i X^i}$ and Σ_R be the surfaces in M defined by

$$\Sigma_R = \{(r, u^1, u^2) : r = R\tilde{F}(u^1, u^2)\}$$

By (4.1), the induced metric σ_{ab} and mean curvature vector H of Σ_R satisfy

$$(4.2) \quad \begin{aligned} \sigma_{ab} &= (\tilde{F}^2 R^2) \tilde{\sigma}_{ab} + O(1) \\ |H| &= \frac{2}{R} - \frac{g_{rr}^{(-5)} + tr_{S^2} p_{ab}^{(-1)} + \frac{3}{2} tr_{S^2} g_{ab}^{(-1)}}{\tilde{F}^3 R^2} + O(R^{-3}). \end{aligned}$$

Let H_0 be the mean curvature of the image of isometric embedding of Σ_R in \mathbb{R}^3 . Recall that the Liu-Yau mass of the surface Σ_R is

$$m_{LY}(\Sigma_R) = \frac{1}{8\pi} \int_{\Sigma_R} (H_0 - |H|) d\Sigma_R.$$

We have

$$H_0 = \frac{2}{R} + O(R^{-3}).$$

since $(\tilde{F}^2 R^2) \tilde{\sigma}_{ab}$ is isometric to $R^2 \tilde{\sigma}_{ab}$. As a result,

$$(4.3) \quad \int_{\Sigma_R} (H_0 - |H|) d\Sigma_R = \int_{S^2} \frac{g_{rr}^{(-5)} + tr_{S^2} p_{ab}^{(-1)} + \frac{3}{2} tr_{S^2} g_{ab}^{(-1)}}{\tilde{F}} dS^2 + O(R^{-1}).$$

By the positivity of the Liu-Yau mass [20], we have

$$\int_{S^2} \frac{g_{rr}^{(-5)} + tr_{S^2} p_{ab}^{(-1)} + \frac{3}{2} tr_{S^2} g_{ab}^{(-1)}}{\tilde{F}} dS^2 \geq 0.$$

Moreover, by Theorem 3.2,

$$\frac{1}{8\pi} \int_{S^2} \frac{g_{rr}^{(-5)} + tr_{S^2} p_{ab}^{(-1)} + \frac{3}{2} tr_{S^2} g_{ab}^{(-1)}}{\tilde{F}} dS^2 = a^0 E_{AH} - \sum_i a^i P_{AH}^i.$$

Hence,

$$a^0 E_{AH} - \sum_i a^i P_{AH}^i \geq 0$$

for any future directed unit timelike vector (a^0, a^i) in $\mathbb{R}^{3,1}$. It follows that (E_{AH}, P_{AH}^i) is future-directed non-spacelike. \square

Corollary 4.3. *Assuming the energy-momentum vector (E_{AH}, P_{AH}^i) for a given foliation is timelike, then there exists a foliation with vanishing linear momentum.*

Proof. From Equation (4.3), it follows that the energy component of the energy momentum vector transformed equivariantly under boost. As a result, if the energy-momentum vector (E_{AH}, P_{AH}^i) for a given foliation is timelike, then we can align the unit vector (a^0, a^i) with the energy momentum and the new foliation corresponding to (a^0, a^i) have vanishing linear momentum. \square

5. Energy loss formula for asymptotically hyperbolic manifolds

Energy loss formula at null infinity for solutions of the vacuum Einstein equation was first studied in [6] and [29]. In [12], a similar energy loss formula at null infinity is proved as a consequence of the global nonlinear stability of the Minkowski space, for any global development of asymptotically flat vacuum initial data sufficiently close to that of the Minkowski space. The source of the energy loss is the traceless part of the inward null second fundamental form of the coordinate spheres. In this section, we study the mass loss formula for asymptotically hyperbolic initial data sets under the Einstein equation. By Theorem 3.2, we can evaluate the total energy at infinity of an asymptotically hyperbolic initial data set with respect to a given foliation of coordinate spheres using the limit of the Liu-Yau mass. We first prove the following energy loss formula for a vacuum spacetime.

Theorem 5.1. *Let (M, g, k) be an asymptotically hyperbolic initial data set satisfying the vacuum constraint equation. Let $(M, g(t), k(t))$ be the solution to the vacuum Einstein equation with $g(0) = g$ and $k(0) = k$, and with lapse $\sqrt{r^2 + 1}$ and shift vector $-re_3$ where e_3 is the unit normal vector of the coordinate spheres. Let $\Sigma_{r,t}$ be the coordinate spheres on $(M, g(t), k(t))$. Let*

$$M(t) = \lim_{r \rightarrow \infty} m_{LY}(\Sigma_{r,t}),$$

then

$$\partial_t M(t)|_{t=0} = -\frac{1}{8\pi} \int_{S^2} |p_{ab}^{(0)} + g_{ab}^{(0)}|^2 dS^2.$$

Proof. Consider the timelike hypersurface

$$\Omega_s = \cup_{0 \leq t \leq s} \Sigma_{r,t}.$$

Let V be the unit normal vector of $\Sigma_{r,t}$ in Ω_s . We have $V = \sqrt{r^2 + 1}e_4 - re_3$ on $\Sigma_{r,t}$ and

$$(5.1) \quad \langle H, V \rangle = O(r^{-2})$$

by the calculation in Lemma 3.1. Let γ_{ij} and Θ_{ij} be the metric and second fundamental form of Ω_s and $\pi_{ij} = \Theta_{ij} - (tr_\gamma \Theta)\gamma_{ij}$ be the conjugate momentum. Let ν be the outward unit normal of Ω_s . Let $H(\Sigma_{r,0})$ and $H(\Sigma_{r,s})$ be the mean curvature vector of $\Sigma_{r,0}$ and $\Sigma_{r,s}$, respectively, in the spacetime and let $H_0(\Sigma_{r,0})$ and $H_0(\Sigma_{r,s})$ be the mean curvature of the image of isometric embedding of $\Sigma_{r,0}$ and $\Sigma_{r,s}$, respectively, into \mathbb{R}^3 . Integrating by parts,

$$\int_{\Sigma_{r,s}} \pi(V, V) d\Sigma_{r,s} = \int_{\Sigma_{r,0}} \pi(V, V) d\Sigma_{r,s} + \int_{\Omega_s} \pi_{ij} V_{i,j}.$$

By the definition of the conjugate momentum,

$$\begin{aligned} \pi(V, V)|_{\Sigma_{r,0}} &= \langle H(\Sigma_{r,0}), \nu \rangle \\ \pi(V, V)|_{\Sigma_{r,s}} &= \langle H(\Sigma_{r,s}), \nu \rangle. \end{aligned}$$

From (5.1), we have

$$\begin{aligned} |H(\Sigma_{r,0})| &= \langle H(\Sigma_{r,0}), \nu \rangle + O(r^{-3}) \\ |H(\Sigma_{r,s})| &= \langle H(\Sigma_{r,s}), \nu \rangle + O(r^{-3}) \end{aligned}$$

and thus

$$\int_{\Sigma_{r,s}} |H(\Sigma_{r,s})| d\Sigma_{r,s} = \int_{\Sigma_{r,0}} |H(\Sigma_{r,0})| d\Sigma_{r,0} + \int_{\Omega_s} \pi_{ij} V_{i,j} + O(r^{-1}).$$

From (5.1), the area $A(r, s)$ of $\Sigma_{r,s}$ satisfies

$$A(r, s) = A(r, 0) + O(1).$$

From Lemma 2.1 of [16], we have

$$\int_{\Sigma_{r,s}} H_0(\Sigma_{r,s})d\Sigma_{r,s} = 4\pi r + \frac{A(r, s)}{r} + O(r^{-1}).$$

We conclude that

$$\int_{\Sigma_{r,s}} H_0(\Sigma_{r,s})d\Sigma_{r,s} = \int_{\Sigma_{r,0}} H_0(\Sigma_{r,0})d\Sigma_{r,0} + O(r^{-1}).$$

It follows that

$$m_{LY}(\Sigma_{r,s}) = m_{LY}(\Sigma_{r,0}) + \frac{1}{8\pi} \int_{\Omega_s} \pi_{ij}V_{i,j} + O(r^{-1})$$

and

$$\begin{aligned} \partial_t m_{LY}(\Sigma_{r,t})|_{t=0} &= \frac{1}{8\pi} \int_{\Sigma_{r,0}} \pi_{ij}V_{i,j}d\Sigma_{r,0} + O(r^{-1}) \\ &= \frac{1}{8\pi} \int_{\Sigma_{r,0}} \pi^{ab}V_{a,b}d\Sigma_{r,0} + O(r^{-1}). \end{aligned}$$

Let $h_{ab}^{(3)}$ and $h_{ab}^{(4)}$ be the second fundamental form of Σ_r in the direction of e_3 and e_4 . Furthermore, let $\hat{h}_{ab}^{(3)}$ and $\hat{h}_{ab}^{(4)}$ be their traceless part.

$$\pi^{ab}V_{a,b} = \pi^{ab}(\sqrt{r^2 + 1}h_{ab}^{(4)} - rh_{ab}^{(3)})$$

since π is symmetric and the symmetrization of $V_{a,b}$ is the second fundamental form of Σ_r in the direction of $\sqrt{r^2 + 1}e_4 - re_3$. Equation (5.1) implies that

$$\sigma^{ab}(\sqrt{r^2 + 1}h_{ab}^{(4)} - rh_{ab}^{(3)}) = O(r^{-2}).$$

From the definition of conjugate momentum,

$$\begin{aligned} &\pi^{ab}(\sqrt{r^2 + 1}h_{ab}^{(4)} - rh_{ab}^{(3)}) \\ &= [\sqrt{r^2 + 1}(h^{(3)})^{ab} - r(h^{(4)})^{ab} - (tr_\gamma \Theta)g^{ab}](\sqrt{r^2 + 1}h_{ab}^{(4)} - rh_{ab}^{(3)}) + O(r^{-3}) \\ &= (\sqrt{r^2 + 1}(\hat{h}^{(3)})^{ab} - r(\hat{h}^{(4)})^{ab})(\sqrt{r^2 + 1}h_{ab}^{(4)} - rh_{ab}^{(3)}) + O(r^{-3}) \\ &= (\sqrt{r^2 + 1}(\hat{h}^{(3)})^{ab} - r(\hat{h}^{(4)})^{ab})(\sqrt{r^2 + 1}\hat{h}_{ab}^{(4)} - r\hat{h}_{ab}^{(3)}) + O(r^{-3}). \end{aligned}$$

From Equation (3.2), we have

$$h_{ab}^{(3)} = \sqrt{r^2 + 1 - \frac{g_{rr}^{(-5)}}{r}} r\tilde{\sigma}_{ab} + O(r^{-1}).$$

By Definition 1.1, we have

$$\begin{aligned}\sigma_{ab} &= r^2 \tilde{\sigma}_{ab} + g_{ab}^{(0)} + O(r^{-1}) \\ h_{ab}^{(4)} &= \sigma_{ab} + p_{ab}^{(0)} + O(r^{-1})\end{aligned}$$

where $tr_{S^2} p_{ab}^{(0)} = tr_{S^2} g_{ab}^{(0)} = 0$. It follows that

$$\begin{aligned}\hat{h}_{ab}^{(3)} &= -g_{ab}^{(0)} + O(r^{-1}) \\ \hat{h}_{ab}^{(4)} &= p_{ab}^{(0)} + O(r^{-1}).\end{aligned}$$

As a result, we have

$$\begin{aligned}& (\sqrt{r^2 + 1}(\hat{h}^{(3)})^{ab} - r(\hat{h}^{(4)})^{ab})(\sqrt{r^2 + 1}\hat{h}_{ab}^{(4)} - r\hat{h}_{ab}^{(3)}) \\ &= \frac{-|p_{ab}^{(0)} + g_{ab}^{(0)}|^2}{r^2} + O(r^{-3})\end{aligned}$$

and

$$\partial_t m_{LY}(\Sigma_{r,t})|_{t=0} = -\frac{1}{8\pi} \int_{S^2} \frac{|p_{ab}^{(0)} + g_{ab}^{(0)}|^2}{r^2} dS^2 + O(r^{-1}).$$

The theorem follows from taking the limit as r approaches infinity. □

In the remaining part of this section, we generalize the above mass loss formula to spacetime with matter fields. Instead of considering specific matter fields, we require the matter fields to satisfies only the dominant energy condition and the following decay condition for the stress-energy density $T_{\mu\nu}$

Definition 5.2. Let (M, g, k) be an asymptotically hyperbolic initial data set in a spacetime N with some matter field. Let e_3 and e_4 be the unit normal of Σ_r in M and the unit normal of M in N . We say that the stress-energy density $T_{\mu\nu}$ satisfies the null decay condition if

$$\begin{aligned}T(e_3 + e_4, e_3 + e_4) &= O(r^{-3}), \\ T(e_3 - e_4, e_3 - e_4) &= O(r^{-4}).\end{aligned}$$

The null decay condition is natural for the global development of an asymptotically flat initial data set. For example, see [37] for the asymptotics at null infinity for solutions of Einstein-Maxwell equation. In particular, the

null decomposition of the the stress-energy density satisfies

$$\begin{aligned} T(L, L) &= O(r^{-5}), \\ T(\underline{L}, \underline{L}) &= O(r^{-2}), \end{aligned}$$

and the leading term of $T(\underline{L}, \underline{L})$ has an additional contribution to the mass loss formula [37]. By considering the standard hyperbolic space embedded in the Minkowski space (see Remark 1.7), it is natural to identify L with $\frac{e_3 + e_4}{r}$ and \underline{L} with $r(e_3 + e_4)$ and assume the null condition in Definition 5.2.

The dominant energy condition implies that both $T(e_3 + e_4, e_3 + e_4)$ and $T(e_3 - e_4, e_3 - e_4)$ are non-negative. Let

$$(5.2) \quad T(e_3 - e_4, e_3 - e_4) = \frac{F^2}{r^4} + O(r^{-5}).$$

We have

Theorem 5.3. *Let (M, g, k) be an asymptotically hyperbolic initial data set satisfying the dominant energy condition equation. Let $(M, g(t), k(t))$ be the solution to the Einstein equation with $g(0) = g$ and $k(0) = k$ with lapse $\sqrt{r^2 + 1}$ and shift vector $-re_3$ where e_3 is the unit normal vector of the coordinate spheres. Let $\Sigma_{r,t}$ be the coordinate spheres on $(M, g(t), k(t))$. Let*

$$M(t) = \lim_{r \rightarrow \infty} m_{LY}(\Sigma_{r,t}).$$

Suppose the stress energy density satisfies both the null condition and the dominant energy condition. Then

$$\partial_t M(t)|_{t=0} = -\frac{1}{8\pi} \int_{S^2} (|p_{ab}^{(0)} + g_{ab}^{(0)}|^2 + F^2) dS^2$$

where F is given by (5.2).

Proof. The proof is similar to that of the previous theorem. Due to the non-zero stress energy density, we have

$$\begin{aligned} \int_{\Sigma_{r,s}} \pi(V, V) d\Sigma_{r,s} &= \int_{\Sigma_{r,0}} \pi(V, V) d\Sigma_{r,s} + \int_{\Omega_s} \pi_{ij} V_{i,j} + \int_{\Omega_s} \nabla^i \pi_{ij} V_j \\ &= \int_{\Sigma_{r,0}} \pi(V, V) d\Sigma_{r,s} + \int_{\Omega_s} \pi_{ij} V_{i,j} + \int_{\Omega_s} T(V, \nu) \end{aligned}$$

where ν is the unit normal of the timelike hypersurface Ω_s . As a result,

$$\partial_t m_{LY}(\Sigma_{r,t})|_{t=0} = \frac{1}{8\pi} \int_{\Sigma_{r,0}} (\pi^{ab} V_{a,b} + T(V, \nu)) d\Sigma_{r,0} + O(r^{-1})$$

The first term can be treated as before. For the second term,

$$\begin{aligned} T(V, \nu) &= T\left(\sqrt{1+r^2}e_4 - re_3, \sqrt{1+r^2}e_3 - re_4\right) \\ &= \frac{1}{4} \left[-(\sqrt{1+r^2} + r)^2 T(e_4 - e_3, e_4 - e_3) \right. \\ &\quad \left. + (\sqrt{1+r^2} - r)^2 T(e_4 + e_3, e_4 + e_3) \right] \\ &= -\frac{F^2}{r^2} + O(r^{-3}). \end{aligned}$$

□

6. Finiteness of total angular momentum and center of mass

Let $(\sigma, |H|, \alpha_H)$ be the induced metric, norm of mean curvature vector and connection one-form in mean curvature gauge on the coordinate spheres Σ_r . The data admits the same expansion as data of the coordinate spheres at the infinity of an asymptotically flat initial data set of order 1 studied in [9]. Thus, we solve the optimal embedding equation and define the total angular momentum and center of mass of (M, g, k) as the limit of quasi-local angular momentum and center of mass on the coordinate spheres in the same manner as in Section 6 of [9].

Denote the expansion of the data $(\sigma, |H|, \alpha_H)$ on the coordinate spheres by

$$\begin{aligned} (6.1) \quad \sigma &= r^2 \tilde{\sigma}_{ab} + O(r^{-1}) \\ |H| &= \frac{2}{r} + \frac{h^{(-2)}}{r^2} + \frac{h^{(-3)}}{r^3} + O(r^{-4}) \\ (\alpha_H)_a &= \frac{(\alpha_H^{(-1)})_a}{r} + \frac{(\alpha_H^{(-2)})_a}{r^2} + O(r^{-3}). \end{aligned}$$

From [8] and [9] for the solution of the optimal embedding equation, we have the following expansion for the solution $(X(r), T_0(r))$ of the optimal

embedding equation

$$\begin{aligned}
 X(r)^0 &= (X^0)^{(0)} + O(r^{-1}) \\
 (6.2) \quad X(r)^i &= r\tilde{X}^i + (X^i)^{(0)} + \frac{(X^i)^{(-1)}}{r} + O(r^{-2}) \\
 T_0(r) &= (a^0, a^1, a^2, a^3) + \frac{T_0^{(-1)}}{r} + O(r^{-2}).
 \end{aligned}$$

We have corresponding expansions for the data of the image of the isometric embedding

$$\begin{aligned}
 (6.3) \quad |H_0| &= \frac{2}{r} + \frac{h_0^{(-2)}}{r^2} + \frac{h_0^{(-3)}}{r^3} + O(r^{-4}) \\
 (\alpha_{H_0})_a &= \frac{(\alpha_{H_0}^{(-1)})_a}{r} + \frac{(\alpha_{H_0}^{(-2)})_a}{r^2} + O(r^{-3}),
 \end{aligned}$$

and the expansion for ρ as defined in Equation (2.1)

$$\rho = \frac{\rho^{(-2)}}{r^2} + \frac{\rho^{(-3)}}{r^3} + O(r^{-4}).$$

The metric expansion implies that $h_0^{(-2)} = 0$, $(X^i)^{(0)} = 0$, and $\rho^{(-2)} = \frac{-h^{(-2)}}{a^0}$.

The leading order term of the optimal embedding equation (see for example equation (7.4) of [9]) implies that $(X^0)^{(0)}$ solves the following equation

$$\begin{aligned}
 (6.4) \quad & \frac{1}{2}\tilde{\Delta}(\tilde{\Delta} + 2)(X^0)^{(0)} \\
 &= \operatorname{div}_{S^2} \alpha_H^{(-1)} - \sum_{i=1}^3 a^i \left[\tilde{\Delta} \left(\frac{\rho^{(-2)} \tilde{X}^i}{2} \right) + \operatorname{div}_{S^2} \rho^{(-2)} (\tilde{\nabla} \tilde{X}^i) \right] \\
 &= -\frac{1}{2}\tilde{\Delta} h^{(-2)} + \sum_{i=1}^3 \frac{a^i}{a^0} \left[\operatorname{div}_{S^2} (h^{(-2)} \tilde{\nabla} \tilde{X}^i) + \tilde{\Delta} \left(\frac{h^{(-2)} \tilde{X}^i}{2} \right) \right].
 \end{aligned}$$

This is solvable if and only if (a^0, a^i) is the future directed unit timelike vector in the direction of the total energy momentum.

The total center of mass and angular momentum of (M, g, k) is defined as follows.

Definition 6.1. Suppose $T_0(r) = A(r)(\frac{\partial}{\partial x^0})$ for a family of Lorentz transformation $A(r)$. Define

$$(6.5) \quad C^i = \lim_{r \rightarrow \infty} E \left(\Sigma_r, X(r), T_0(r), A(r) \left(x^i \frac{\partial}{\partial x^0} + x^0 \frac{\partial}{\partial x^i} \right) \right)$$

to be the total center of mass integral and

$$(6.6) \quad J_i = \lim_{r \rightarrow \infty} \epsilon_{ijkl} E \left(\Sigma_r, X(r), T_0(r), A(r) \left(x^j \frac{\partial}{\partial x^k} - x^k \frac{\partial}{\partial x^j} \right) \right)$$

to be the total angular momentum, where Σ_r are the coordinate spheres and $(X(r), T_0(r))$ is the unique family of optimal isometric embeddings of Σ_r such that $X(r)$ converges to a round sphere of radius r in \mathbb{R}^3 when $r \rightarrow \infty$.

In [9], the authors studied the limit of quasi-local conserved quantities at the infinity of asymptotically flat initial data sets of order 1 and proved the following theorem concerning finiteness of total conserved quantities:

Proposition 6.2 (Proposition 7.1, [9]). *The total center of mass and total angular momentum are finite if*

$$(6.7) \quad \int_{S^2} \tilde{X}^i \rho^{(-2)} dS^2 = 0 \text{ and } \int_{S^2} \tilde{X}^i \left(\tilde{\epsilon}^{ab} \tilde{\nabla}_b (\alpha_H^{(-1)})_a \right) dS^2 = 0.$$

Proposition 7.1 of [9] can be applied to the limit of quasi-local conserved quantities at infinity of asymptotically hyperbolic initial data since $(\sigma, |H|, \alpha_H)$ has expansions of the same forms.

In addition, the authors proved in [9] that Equation (6.7) follows from the vacuum constraint equation for asymptotically flat initial data sets of order 1. We prove in this section that, for coordinate spheres in asymptotically hyperbolic initial data sets, Equation (6.7) does not follow from the vacuum constraint equation. Instead, it holds only for the foliation with vanishing linear momentum.

Proposition 6.3. *Let Σ_r be a family of coordinate spheres in an asymptotically hyperbolic initial data sets. Then*

$$(6.8) \quad \int_{S^2} \tilde{X}^i \rho^{(-2)} dS^2 = 0 \text{ and } \int_{S^2} \tilde{X}^i \left(\tilde{\epsilon}^{ab} \tilde{\nabla}_b (\alpha_H^{(-1)})_a \right) dS^2 = 0$$

holds if and only if the linear momentum vanishes.

Proof. Recall from Lemma 3.1, we have

$$(6.9) \quad h^{(-2)} = -m \text{ and } (\alpha_H^{(-1)})_a = \frac{1}{2} \partial_a m,$$

where m is the mass aspect function. By the expansion of the induced metric on Σ_r in Lemma 3.1,

$$h_0^{(-2)} = 0.$$

It follows immediately that

$$\int_{S^2} \tilde{X}^i \left(\tilde{\epsilon}^{ab} \tilde{\nabla}_b (\alpha_H^{(-1)})_a \right) dS^2 = 0$$

holds and

$$\int_{S^2} \tilde{X}^i \rho^{(-2)} dS^2 = \frac{P_{AH}^i}{E_{AH}}. \quad \square$$

Corollary 6.4. *For an asymptotically hyperbolic initial data set with time-like energy-momentum vector, the total angular momentum and center of mass associated with any foliation with vanishing linear momentum are finite.*

Proof. The corollary follows from combining Corollary 4.3, Proposition 6.2, and Proposition 6.3. □

7. Evaluating total center of mass and angular momentum

Let (M, g, k) be an asymptotically hyperbolic initial data set. Suppose a foliation Σ_r has vanishing linear momentum, we evaluate its total center of mass and angular momentum. We first evaluate the total center of mass and angular momentum in terms of the expansion of $(\sigma, |H|, \alpha_H)$ and $(|H_0|, \alpha_{H_0})$. Then we express them in terms of the expansion for g and k .

Let $(X(r), T_0(r))$ be the solution of the optimal embedding equation on Σ_r . The solution takes a simpler form since the linear momentum vanishes. In particular, we have

$$(7.1) \quad \begin{aligned} X^0(r) &= (X^0)^{(0)} + O(r^{-1}) \\ X^i(r) &= r \tilde{X}^i + \frac{(X^i)^{(-1)}}{r} + O(r^{-2}) \\ T_0(r) &= (1, 0, 0, 0) + \frac{T_0^{(-1)}}{r} + O(r^{-2}) \end{aligned}$$

where $(X^0)^{(0)}$ satisfies the equation

$$(7.2) \quad \frac{1}{2}\tilde{\Delta}(\tilde{\Delta} + 2)(X^0)^{(0)} = -\frac{1}{2}\tilde{\Delta}h^{(-2)}.$$

Let $T_0^{(-1)} = (0, b^1, b^2, b^3)$. The time function $\tau = -X(r) \cdot T_0(r)$ has the following expansion

$$(7.3) \quad \tau = (X^0)^{(0)} - \sum_i b^i \tilde{X}^i + O(r^{-1}).$$

Proposition 7.1. *The total center of mass and angular momentum of an asymptotically hyperbolic initial data set are*

$$(7.4) \quad \begin{aligned} C^i &= \frac{1}{8\pi} \int_{S^2} \tilde{X}^i (h_0^{(-3)} - h^{(-3)}) dS^2 \\ J^i &= \frac{-1}{8\pi} \int_{S^2} \tilde{X}^i \epsilon^{ab} \tilde{\nabla}_b (\alpha_H^{(-2)})_a dS^2. \end{aligned}$$

Proof. Since $\tau = O(1)$,

$$(7.5) \quad \rho = |H_0| - |H| + O(r^{-4}).$$

Recall that,

$$j_a = \rho \nabla_a \tau - \nabla_a \left[\sinh^{-1} \left(\frac{\rho \Delta \tau}{|H_0| |H|} \right) \right] - (\alpha_{H_0})_a + (\alpha_H)_a$$

and

$$(7.6) \quad \begin{aligned} j_a^{(-1)} &= -(\alpha_{H_0})_a^{(-1)} + (\alpha_H)_a^{(-1)} \\ &= \frac{1}{2} \left[-\tilde{\nabla}_a (\tilde{\Delta} + 2)(X^0)^{(0)} - \tilde{\nabla} h^{(-2)} \right] = 0, \end{aligned}$$

by the optimal isometric embedding equation. The Killing vector field associated with the center of mass is $K_{0i} = A(r)(x^i \frac{\partial}{\partial x^0} + x^0 \frac{\partial}{\partial x^i})$ with $A(r) = Id + O(r^{-1})$. This implies that $(K_{0i}^\top)_a = O(r)$. We also have

$$\langle K_{0i}, T_0(r) \rangle = -r \tilde{X}^i + O(r^{-1}).$$

It follows that

$$\begin{aligned}
 (7.7) \quad & E\left(\Sigma_r, X(r), T_0(r), A(r)\left(x^i \frac{\partial}{\partial x^0} + x^0 \frac{\partial}{\partial x^i}\right)\right) \\
 &= \frac{-1}{8\pi} \int_{\Sigma_r} \left[\langle K_{0i}, T_0(r) \rangle \rho + (K_{0i}^\top)_a j^a\right] d\Sigma_r \\
 &= \frac{1}{8\pi} \int_{S^2} (h_0^{(-3)} - h^{(-3)}) \tilde{X}^i dS^2 + O(r^{-1}).
 \end{aligned}$$

For the angular momentum, the corresponding Killing vector field is $K_{ij} = A(r)(x^i \frac{\partial}{\partial x^j} - x^j \frac{\partial}{\partial x^i})$. We have $\langle K_{ij}, T_0(r) \rangle = 0$. Hence

$$\begin{aligned}
 (7.8) \quad & E\left(\Sigma_r, X(r), T_0(r), A(r)\left(x^i \frac{\partial}{\partial x^j} - x^j \frac{\partial}{\partial x^i}\right)\right) \\
 &= \frac{(-1)}{8\pi} \int_{\Sigma_r} K_{ij}^\top \cdot \left(\frac{\sqrt{(1 + |\nabla\tau|^2)|H_0|^2 + (\Delta\tau)^2} \nabla\tau}{(1 + |\nabla\tau|^2)}\right. \\
 &\quad \left. - \nabla \left[\sinh^{-1} \left(\frac{\Delta\tau}{|H_0| \sqrt{1 + |\nabla\tau|^2}}\right)\right] - \alpha_{H_0}\right) d\Sigma_r \\
 &+ \frac{1}{8\pi} \int_{\Sigma_r} K_{ij}^\top \cdot \left(\frac{\sqrt{(1 + |\nabla\tau|^2)|H|^2 + (\Delta\tau)^2} \nabla\tau}{(1 + |\nabla\tau|^2)}\right. \\
 &\quad \left. - \nabla \left[\sinh^{-1} \left(\frac{\Delta\tau}{|H| \sqrt{1 + |\nabla\tau|^2}}\right)\right] - \alpha_H\right) d\Sigma_r.
 \end{aligned}$$

We prove that the first integral is $O(r^{-1})$ by applying the conservation law to the image of the isometric embedding $X(r)$ of Σ_r , similar to the proof of Equation (8.1) in [9]. Let \mathcal{C}_r be the timelike cylinder in $\mathbb{R}^{3,1}$ obtained by translating $X(r)(\Sigma_r)$ along $T_0(r)$. Let $\hat{\Sigma}_r$ be the projection of $X(r)(\Sigma_r)$ onto the orthogonal complement of $T_0(r)$ and Ω_r be the portion of \mathcal{C}_r between $X(r)(\Sigma_r)$ and $\hat{\Sigma}_r$. Let γ_{ij} and Θ_{ij} be the metric and second fundamental form of \mathcal{C}_r and $\pi_{ij} = \Theta_{ij} - tr_\gamma \Theta \gamma_{ij}$ be the conjugate momentum. Let ν be the outward unit normal of \mathcal{C}_r . From the expansion of $X(r)$, it is easy to see that

$$\langle K_{ij}, \nu \rangle = O(r^{-1}).$$

Let

$$K^{\mathcal{C}} = K_{ij} - \langle K_{ij}, \nu \rangle \nu$$

be the tangential part of K_{ij} to \mathcal{C}_r and consider the vector field

$$Z = \pi(K^{\mathcal{C}}, \cdot).$$

We claim that $div_{\mathcal{C}}Z = O(r^{-3})$. Indeed, the conjugate momentum is divergence free since $\mathbb{R}^{3,1}$ is flat. We also have K_{ij} is Killing in $\mathbb{R}^{3,1}$. As a result, we have

$$div_{\mathcal{C}}Z = \pi_{ij}\Theta_{ij}\langle K, \nu \rangle = O(r^{-3}).$$

Applying the divergence theorem to Z , we derive

$$\int_{\Sigma_r} \pi(K^{\mathcal{C}}, n) = \int_{\hat{\Sigma}_r} \pi(K^{\mathcal{C}}, \hat{n}) + \int_{\Omega_r} div_{\mathcal{C}}Z$$

where n and \hat{n} are the unit normal of $X(r)(\Sigma_r)$ and $\hat{\Sigma}_r$ in \mathcal{C}_r .

By definition, $\int_{\Sigma_r} \pi(K^{\mathcal{C}}, n)$ is the reference term in the angular momentum. Moreover, since $\hat{\Sigma}_r$ lies in a totally geodesics slice in $\mathbb{R}^{3,1}$ we have

$$\int_{\hat{\Sigma}_r} \pi(K^{\mathcal{C}}, \hat{n}) = 0.$$

Finally, since $div_{\mathcal{C}}Z = O(r^{-3})$, we have $\int_{\Omega_r} div_{\mathcal{C}}Z = O(r^{-1})$. It follows that the first integral in (7.8) is of the order $O(r^{-1})$ as desired.

We proceed to study the second integral in (7.8). We have

$$(K_{ij}^{\top})_a = r^2(\tilde{X}^i\tilde{\nabla}_a\tilde{X}^j - \tilde{X}^j\tilde{\nabla}_a\tilde{X}^i) + O(1)$$

and

$$\begin{aligned} & \frac{1}{8\pi} \int_{\Sigma_r} K_{ij}^{\top} \cdot \left(\frac{\sqrt{(1 + |\nabla\tau|^2)|H|^2 + (\Delta\tau)^2}\nabla\tau}{(1 + |\nabla\tau|^2)} \right. \\ & \quad \left. - \nabla \left[\sinh^{-1} \left(\frac{\Delta\tau}{|H|\sqrt{1 + |\nabla\tau|^2}} \right) \right] - \alpha_H \right) d\Sigma_r \\ &= \frac{1}{8\pi} \int_{S^2} (\tilde{X}^i\tilde{\nabla}^a\tilde{X}^j - \tilde{X}^j\tilde{\nabla}^a\tilde{X}^i) \\ & \quad \times \left(h^{(-2)}\tilde{\nabla}_a \left(X_0^{(0)} - \sum_i b^i\tilde{X}^i \right) - \frac{1}{2}\tilde{\nabla}_a(\tilde{\Delta}X_0^{(0)}) - (\alpha_H^{(-2)})_a \right) dS^2 \\ & + O(r^{-1}). \end{aligned}$$

Notice that

$$(7.9) \quad \tilde{X}^i\tilde{\nabla}^a\tilde{X}^j - \tilde{X}^j\tilde{\nabla}^a\tilde{X}^i = \tilde{\epsilon}^{ab}\epsilon_k{}^{ij}\tilde{\nabla}_b\tilde{X}^k$$

is divergence free and thus

$$\int_{S^2} (\tilde{X}^i \tilde{\nabla}^a \tilde{X}^j - \tilde{X}^j \tilde{\nabla}^a \tilde{X}^i) \tilde{\nabla}_a (\tilde{\Delta} X_0^{(0)}) dS^2 = 0.$$

We also claim that

$$(7.10) \quad \int_{S^2} (\tilde{X}^i \tilde{\nabla}^a \tilde{X}^j - \tilde{X}^j \tilde{\nabla}^a \tilde{X}^i) \left[h^{(-2)} \tilde{\nabla}_a \left(X_0^{(0)} - \sum_k b^k \tilde{X}^k \right) \right] dS^2 = 0.$$

Recall from (7.2) that $\tilde{\Delta}(\tilde{\Delta} + 2)(X^0)^{(0)} = -\tilde{\Delta}h^{(-2)}$, or $(\tilde{\Delta} + 2)(X^0)^{(0)}$ and $-h^{(-2)}$ differ by a constant.

In view of (7.9), (7.10) is implied by

$$\int_{S^2} \tilde{\epsilon}^{ab} \tilde{\nabla}_b \tilde{X}^l \tilde{\nabla}_a \left(X_0^{(0)} - \sum_k b^k \tilde{X}^k \right) (\tilde{\Delta} + 2) X_0^{(0)} dS^2 = 0$$

for $l = 1, 2, 3$.

Define the differential operator L^l by $L^l(f) = \tilde{\epsilon}^{ab} \tilde{\nabla}_b \tilde{X}^l \tilde{\nabla}_a f$ for any C^1 function on S^2 . It suffices to show that

$$\int_{S^2} L^l \left(X_0^{(0)} - \sum_k b^k \tilde{X}^k \right) (\tilde{\Delta} + 2) X_0^{(0)} dS^2 = 0.$$

L^l corresponds to a rotational Killing field on S^2 and it is easy to check that the following properties hold for L^l :

$$L^l(fg) = L^l(f)g + fL^l(g), \quad [\tilde{\Delta}, L^l]f = 0, \quad \text{and} \quad \int_{S^2} L^l(f) dS^2 = 0,$$

for any C^1 functions f, g on S^2 .

Integrating by parts and applying the above relations, we derive

$$\begin{aligned} & \int_{S^2} L^l \left(X_0^{(0)} - \sum_k b^k \tilde{X}^k \right) (\tilde{\Delta} + 2) X_0^{(0)} dS^2 \\ &= \int_{S^2} L^l \left((\tilde{\Delta} + 2) \left(X_0^{(0)} - \sum_k b^k \tilde{X}^k \right) \right) X_0^{(0)} dS^2 \\ &= \int_{S^2} L^l((\tilde{\Delta} + 2)(X_0^{(0)})) X_0^{(0)} dS^2 \\ &= \int_{S^2} L^l(\tilde{\Delta} X_0^{(0)}) X_0^{(0)} dS^2. \end{aligned}$$

On the other hand,

$$\begin{aligned} \int_{S^2} L^l(\tilde{\Delta}X_0^{(0)})X_0^{(0)}dS^2 &= \int_{S^2} \tilde{\Delta}(L^l(X_0^{(0)}))X_0^{(0)}dS^2 \\ &= \int_{S^2} (L^l(X_0^{(0)}))\tilde{\Delta}X_0^{(0)}dS^2. \end{aligned}$$

Therefore,

$$\begin{aligned} &\int_{S^2} L^l(\tilde{\Delta}X_0^{(0)})X_0^{(0)}dS^2 \\ &= \frac{1}{2} \int_{S^2} L^l(\tilde{\Delta}X_0^{(0)})X_0^{(0)}dS^2 + \frac{1}{2} \int_{S^2} (L^l(X_0^{(0)}))\tilde{\Delta}X_0^{(0)}dS^2 \\ &= \frac{1}{2} \int_{S^2} L^l(X_0^{(0)})\tilde{\Delta}X_0^{(0)}dS^2 = 0. \end{aligned}$$

This finishes the proof of the proposition. □

We need one more lemma before we evaluate the limit in terms of the data (M, g, k) .

Lemma 7.2.

$$\int_{S^2} \tilde{X}^i(h_0^{(-3)} - h^{(-3)})dS^2 = - \int_{S^2} \tilde{X}^i \left[h^{(-3)} + \frac{(h^{(-2)})^2}{4} \right] dS^2.$$

Proof. It suffices to prove

$$\int_{S^2} \tilde{X}^i \left[h_0^{(-3)} + \frac{(h^{(-2)})^2}{4} \right] dS^2 = 0.$$

Recall that

$$|H_0|^2 = \sum_i (\Delta X^i)^2 - (\Delta X^0)^2$$

where

$$\begin{aligned} (7.11) \quad \Delta X^0 &= \frac{\tilde{\Delta}X_0^{(0)}}{r^2} + O(r^{-3}) \\ \Delta X^i &= r\Delta\tilde{X}^i + \frac{\tilde{\Delta}(X^i)^{(-1)}}{r^3} + O(r^{-4}) \end{aligned}$$

Furthermore, let Γ_{ab}^c be the Christoffel symbols of the metric σ_{ab} . We have

$$\Gamma_{ab}^c = \tilde{\Gamma}_{ab}^c + \frac{(\Gamma^{(-2)})_{ab}^c}{r^2} + O(r^{-3}).$$

This gives the following expansion for $\Delta\tilde{X}^i$

$$\begin{aligned} \Delta\tilde{X}^i &= \frac{-2\tilde{X}^i}{r^2} + \frac{\tilde{\sigma}^{ab}g_{ab}^{(0)}\tilde{X}^i + \tilde{\sigma}^{ab}(\Gamma^{(-2)})^c_{ab}\tilde{X}^i_c}{r^4} + O(r^{-5}) \\ &= \frac{-2\tilde{X}^i}{r^2} + \frac{\tilde{\sigma}^{ab}(\Gamma^{(-2)})^c_{ab}\tilde{X}^i_c}{r^4} + O(r^{-5}). \end{aligned}$$

As a result, we have

$$h_0^{(-3)} = -\sum_k \tilde{X}^k \tilde{\Delta}(X^k)^{(-1)} - \frac{1}{4}(\tilde{\Delta}X_0^{(0)})^2.$$

Recall that $(X^i)^{(-1)}$ is the solution of the linearized isometric embedding equation

$$\begin{aligned} (7.12) \quad &\sum_{i=1}^3 \tilde{\nabla}_a \tilde{X}^i \tilde{\nabla}_b (X^i)^{(-1)} + \tilde{\nabla}_b \tilde{X}^i \tilde{\nabla}_a (X^i)^{(-1)} \\ &= \tilde{\nabla}_a (X^0)^{(0)} \tilde{\nabla}_b (X^0)^{(0)} + g_{ab}^{(0)}. \end{aligned}$$

Let $T_{ab} = \tilde{\nabla}_a (X^0)^{(0)} \tilde{\nabla}_b (X^0)^{(0)} + g_{ab}^{(0)}$. We use the ansatz in Nirenberg’s paper [22] for solving the isometric embedding of surface with positive Gauss curvature into \mathbb{R}^3 . Write

$$\tilde{\nabla}_a (X^j)^{(-1)} = P_a \tilde{X}^j + \left(\frac{T_{ab}}{2} + F\epsilon_{ab} \right) \tilde{\sigma}^{bc} \tilde{X}^j_c$$

where P_a is a vector field, F is a function, and ϵ_{ab} is the area form on S^2 . P_a and F as considered to be the new unknowns that satisfy the compatibility condition. We have

$$\begin{aligned} (7.13) \quad \tilde{\nabla}_d \tilde{\nabla}_a X^{(-1)} &= (\tilde{\nabla}_d P_a) \tilde{X} + P_a \tilde{X}_d \\ &\quad + \left(\frac{\tilde{\nabla}_d T_{ab}}{2} + \tilde{\nabla}_d F \epsilon_{ab} \right) \tilde{\sigma}^{bc} \tilde{X}_c - \left(\frac{T_{ad}}{2} + F\epsilon_{ad} \right) \tilde{X}. \end{aligned}$$

Taking the trace of Equation (7.13) and multiplying by $-\tilde{X}$,

$$\sum_k -\tilde{X}^k \tilde{\Delta}(X^k)^{(-1)} = -\tilde{\nabla}^a P_a + \frac{1}{2} \tilde{\sigma}^{ab} T_{ab} = -\tilde{\nabla}^a P_a + \frac{1}{2} |\tilde{\nabla}(X^0)^{(0)}|^2.$$

Multiplying by \tilde{X}^i and integrating, we obtain

$$(7.14) \quad - \int \tilde{X}^i \sum_k \tilde{X}^k \tilde{\Delta}(X^k)^{(-1)} dS^2 = \frac{1}{2} \int \tilde{X}^i |\tilde{\nabla}(X^0)^{(0)}|^2 - \int \tilde{X}^i \tilde{\nabla}^a P_a dS^2.$$

We use the compatibility condition, $\epsilon^{ad} \tilde{\nabla}_d \tilde{\nabla}_a X^{(-1)} = 0$ and (7.13), to obtain equations for P_a and F :

$$\begin{aligned} \epsilon^{ad} \left(\tilde{\nabla}_d P_a - \frac{T_{ad}}{2} - F \epsilon_{ad} \right) &= 0, \\ \epsilon^{ad} \left(P_a \tilde{\sigma}_{cd} + \frac{1}{2} \tilde{\nabla}_d T_{ac} + \tilde{\nabla}_d F \epsilon_{ac} \right) &= 0. \end{aligned}$$

P_a can be solved from the second equation in terms of F and T_{ab} . The equation can be written as

$$(7.15) \quad \epsilon_{ac} P_b \tilde{\sigma}^{ab} + \frac{1}{2} \epsilon^{ad} \tilde{\nabla}_d T_{ac} + \tilde{\nabla}_c F = 0.$$

Multiplying the equation by ϵ^{ce} , we have

$$P^e = \frac{1}{2} \epsilon^{ce} \epsilon^{ad} \tilde{\nabla}_d T_{ac} + \epsilon^{ce} \tilde{\nabla}_c F.$$

Taking the divergence and integrating against \tilde{X}^i ,

$$\begin{aligned} \int_{S^2} \tilde{X}^i \tilde{\nabla}^e P_e dS^2 &= \int_{S^2} \tilde{X}_i \tilde{\nabla}_e \left(\frac{1}{2} \epsilon^{ce} \epsilon^{ad} \tilde{\nabla}_d T_{ac} + \epsilon^{ce} \tilde{\nabla}_c F \right) dS^2 \\ &= \frac{1}{2} \int_{S^2} \tilde{\nabla}_e \tilde{\nabla}_d \tilde{X}^i \epsilon^{ce} \epsilon^{ad} T_{ac} dS^2 \\ &= -\frac{1}{2} \int_{S^2} \tilde{X}^i \tilde{\sigma}_{de} \epsilon^{ce} \epsilon^{ad} T_{ac} dS^2 \\ &= -\frac{1}{2} \int_{S^2} \tilde{X}^i \tilde{\sigma}^{ab} T_{ab} dS^2. \end{aligned}$$

As a result, we have

$$- \int_{S^2} \tilde{X}^i \sum_k \tilde{X}^k \tilde{\Delta}(X^k)^{(-1)} = \int_{S^2} \tilde{X}^i |\tilde{\nabla}(X^0)^{(0)}|^2$$

and

$$\begin{aligned}
 (7.16) \quad & \int_{S^2} \tilde{X}^i \left(h_0^{(-3)} + \frac{(h^{(-2)})^2}{4} \right) dS^2 \\
 &= \int_{S^2} \tilde{X}^i \left(|\tilde{\nabla}(X^0)^{(0)}|^2 - \frac{1}{4}(\tilde{\Delta}X_0^{(0)})^2 + \frac{(h^{(-2)})^2}{4} \right) dS^2 \\
 &= \int_{S^2} \tilde{X}^i \left(|\tilde{\nabla}(X^0)^{(0)}|^2 - \frac{1}{4}(\tilde{\Delta}(X^0)^{(0)})^2 + \frac{[(\tilde{\Delta} + 2)(X^0)^{(0)}]^2}{4} \right) dS^2 \\
 &= \int_{S^2} \tilde{X}^i \left(|\tilde{\nabla}(X^0)^{(0)}|^2 + ((X^0)^{(0)})^2 + X_0^{(0)}\tilde{\Delta}(X^0)^{(0)} \right) dS^2 \\
 &= \int_{S^2} \tilde{X}^i \left(((X^0)^{(0)})^2 + \frac{\tilde{\Delta}}{2}((X^0)^{(0)})^2 \right) dS^2 = 0.
 \end{aligned}$$

□

We are now ready to express the total angular momentum and center of mass in terms of the expansion of the metric g and second fundamental form k . The calculation is straight-forward and is similar to Lemma 3.1. We simply need to expand one more order.

Theorem 7.3. *Given an asymptotically hyperbolic initial data set (M, g, k) , for the foliation with vanishing linear momentum, the center of mass C^i and angular momentum J^i of (M, g, k) are*

$$\begin{aligned}
 (7.17) \quad C^i &= \frac{1}{8\pi} \int_{S^2} \tilde{X}^i (2tr_{S^2}g_{ab}^{(-2)} + g_{rr}^{(-6)} + \tilde{\nabla}^a g_{ra}^{(-3)} + tr_{S^2}p_{ab}^{(-2)}) dS^2 \\
 J^i &= \frac{1}{8\pi} \int_{S^2} \tilde{X}^i \epsilon^{ab} \tilde{\nabla}_b (g_{ra}^{(-3)} + p_{ra}^{(-3)}) dS^2.
 \end{aligned}$$

Proof. We start with the angular momentum. Recall that

$$J^i = \frac{-1}{8\pi} \int_{S^2} \tilde{X}^i \epsilon^{ab} \tilde{\nabla}_b (\alpha_H^{(-2)})_a dS^2.$$

where

$$\begin{aligned}
 (\alpha_H)_a &= -k(e_3, \partial_a) + \nabla_a \theta \\
 k(e_3, \partial_a) &= \frac{g_{ra}^{(-3)} + p_{ra}^{(-3)}}{r^2} + O(r^{-3}).
 \end{aligned}$$

Hence,

$$\begin{aligned} J^i &= \frac{1}{8\pi} \int_{S^2} \tilde{X}^i \tilde{\epsilon}^{ab} \left[\tilde{\nabla}_b (g_{ra}^{(-3)} + p_{ra}^{(-3)}) + \tilde{\nabla}_b \tilde{\nabla}_a \theta \right] dS^2. \\ &= \frac{1}{8\pi} \int_{S^2} \tilde{X}^i \tilde{\epsilon}^{ab} \tilde{\nabla}_b (g_{ra}^{(-3)} + p_{ra}^{(-3)}) dS^2. \end{aligned}$$

For the center of mass, we start with $\langle H, e_4 \rangle$.

$$\begin{aligned} -\langle H, e_4 \rangle &= \sigma^{ab} (g_{ab} + p_{ab}) \\ &= 2 + \frac{tr_{S^2} p_{ab}^{(-1)}}{r^3} + \frac{tr_{S^2} p_{ab}^{(-2)}}{r^4} + O(r^{-5}). \end{aligned}$$

On the other hand, we compute

$$e_3 = \frac{1}{\sqrt{g_{rr} - |V|_\sigma^2}} \left(\frac{\partial}{\partial r} + V^a \frac{\partial}{\partial u^a} \right)$$

where $V_a = -g_{ra}$. The second fundamental form of Σ_r in the direction of e_3 is

$$\begin{aligned} \left\langle \nabla_{\frac{\partial}{\partial u^a}} e_3, \frac{\partial}{\partial u^b} \right\rangle &= \sqrt{r^2 + 1 - \frac{g_{rr}^{(-5)}}{r} - \frac{g_{rr}^{(-6)}}{r^2}} \\ &\quad \times \left\langle \nabla_{\frac{\partial}{\partial u^a}} \frac{\partial}{\partial r} + V^c \frac{\partial}{\partial u^c}, \frac{\partial}{\partial u^b} \right\rangle + O(r^{-3}) \\ &= \sqrt{r^2 + 1 - \frac{g_{rr}^{(-5)}}{r} - \frac{g_{rr}^{(-6)}}{r^2}} \\ &\quad \times \left[r \tilde{\sigma}_{ab} - \frac{g_{ab}^{(-1)}}{2r^2} - \frac{(g_{ab}^{(-2)} + \tilde{\nabla}_a g_{br}^{(-3)})}{r^3} \right] + O(r^{-3}). \end{aligned}$$

Hence,

$$\begin{aligned} -\langle H, e_3 \rangle &= \sigma^{ab} \sqrt{r^2 + 1 - \frac{g_{rr}^{(-5)}}{r} - \frac{g_{rr}^{(-6)}}{r^2}} \\ &\quad \times \left[r \tilde{\sigma}_{ab} - \frac{g_{ab}^{(-1)}}{2r^2} - \frac{(g_{ab}^{(-2)} + \tilde{\nabla}_a g_{br}^{(-3)})}{r^3} \right] + O(r^{-5}) \end{aligned}$$

$$\begin{aligned}
 &= r \left(1 + \frac{1}{2r^2} - \frac{g_{rr}^{(-5)}}{2r^3} - \frac{g_{rr}^{(-6)} + \frac{1}{4}}{2r^3} \right) \left(r\tilde{\sigma}_{ab} - \frac{g_{ab}^{(-1)}}{2r^2} - \frac{g_{ab}^{(-2)} + \tilde{\nabla}_a g_{br}^{(-3)}}{r^3} \right) \\
 &\quad \times \left(\frac{\tilde{\sigma}^{ab}}{r^2} - \frac{(g^{(-1)})^{ab}}{r^5} - \frac{(g^{(-2)})^{ab}}{r^6} \right) + O(r^{-5}) \\
 &= 2 + \frac{1}{r^2} - \frac{g_{rr}^{(-5)} + \frac{3}{2}tr_{S^2}g_{ab}^{(-1)}}{r^3} \\
 &\quad - \frac{g_{rr}^{(-6)} + \frac{1}{4} + 2tr_{S^2}g_{ab}^{(-1)} + \tilde{\nabla}^a g_{ar}^{(-3)}}{r^4} + O(r^{-5}).
 \end{aligned}$$

We have

$$\begin{aligned}
 |H|^2 &= \langle H, e_3 \rangle^2 - \langle H, e_4 \rangle^2 \\
 &= \frac{4}{r^2} + \frac{4h^{(-2)}}{r^3} \\
 &\quad - \frac{4(g_{rr}^{(-6)} + 2tr_{S^2}g_{ab}^{(-1)} + \tilde{\nabla}^a g_{ar}^{(-3)} + tr_{S^2}p_{ab}^{(-2)})}{r^4} + O(r^{-5}).
 \end{aligned}$$

Recall

$$|H| = \frac{2}{r} + \frac{h^{(-2)}}{r^2} + \frac{h^{(-3)}}{r^3} + O(r^{-4}).$$

Matching the coefficients,

$$h^{(-3)} + \frac{(h^{(-2)})^2}{4} = -(g_{rr}^{(-6)} + 2tr_{S^2}g_{ab}^{(-1)} + \tilde{\nabla}^a g_{ar}^{(-3)} + tr_{S^2}p_{ab}^{(-2)}).$$

The formula for the center of mass follows from Lemma 7.2. □

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